

# *Prospects for dark matter detection with inelastic transitions of xenon*



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*preliminary results  
—work in progress—*

# An old idea...

- The original direct detection paper:

PHYSICAL REVIEW D

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## Detectability of certain dark-matter candidates

Mark W. Goodman and Edward Witten

*Joseph Henry Laboratories, Princeton University, Princeton, New Jersey 08544*

(Received 7 January 1985)

Aside from the detector proposed in Ref. 5, an interesting possibility is to detect dark-matter particles via inelastic rather than elastic scattering from nuclei.

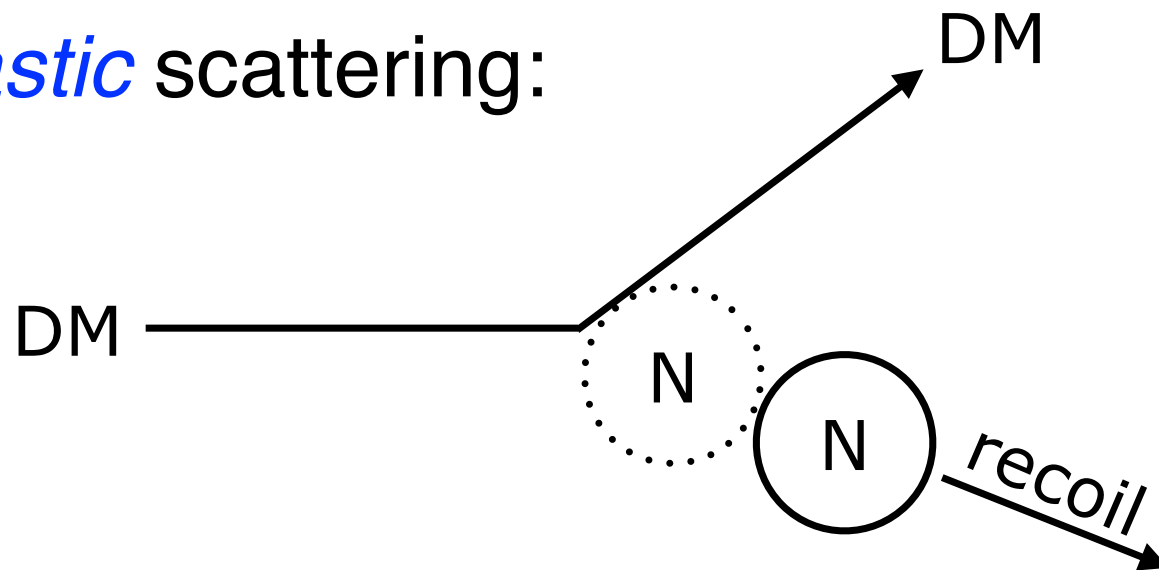
# An old idea... Inelastic scattering

- What is it?
- Why is it interesting?
- Why consider it now?

Can it ever be detected?

# What is it?

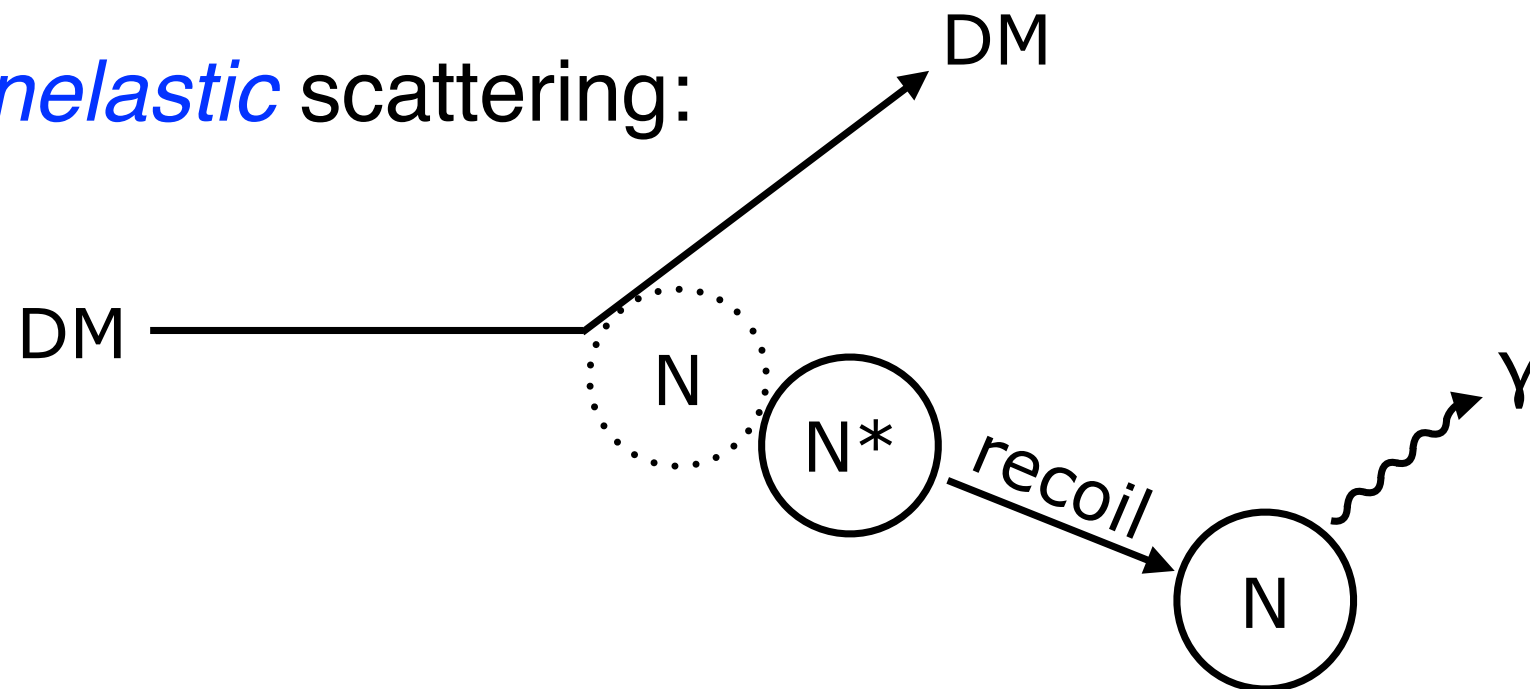
*elastic* scattering:



measure:

*N's recoil energy*

*inelastic* scattering:




measure:

*N's recoil energy*  
+ *photon energy*



# What is a good target?



**XENON**

# Why Xenon?

Inelastic scattering *is not*  $A^2$  enhanced

- ★ Only accessible for spin-dependent interactions
  - ➔ Elastic and inelastic scattering rates comparable

[Vietze et al arXiv:1412.6091](#)

- ★ Ideal target should have
  - good spin-dependent sensitivity
  - a low lying excitation ( $\lesssim E_{\text{DM-kinetic}} \approx 100 \text{ keV}$ )

# Why Xenon?

- 47.6% of xenon sensitive to spin-dependent interactions:

## $^{129}\text{Xe}$

Natural abundance: 26.4%

Lowest excitation: 39.6 keV

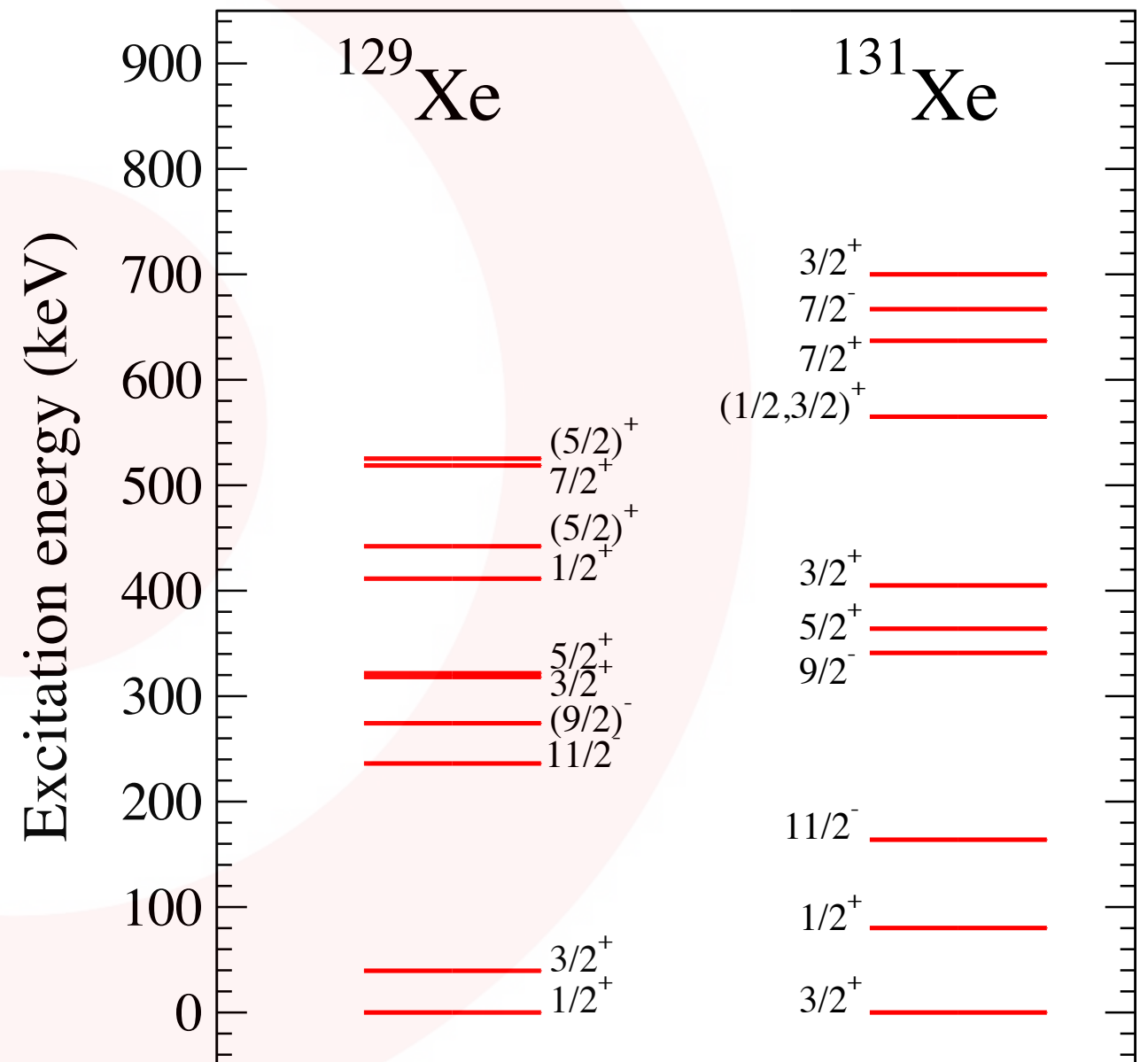
Lifetime: 0.97 ns

## $^{131}\text{Xe}$

Natural abundance: 21.2%

Lowest excitation: 80.2 keV

Lifetime: 0.48 ns



# Previous studies

## Limits on WIMP- $^{129}\text{Xe}$ inelastic scattering

P. Belli<sup>a</sup>, R. Bernabei<sup>a</sup>, V. Landoni<sup>a</sup>, F. Montecchia<sup>a</sup>, W. Di Nicolantonio<sup>b</sup>, A. Incicchitti<sup>b</sup>,  
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## Search for inelastic WIMP nucleus scattering on $^{129}\text{Xe}$ in data from the XMASS-I experiment

- Previous searches with single phase-detectors
- *No limits or studies for two-phase detectors (LUX, XENON)*

# Why is it interesting?

*Inferring properties of dark matter is difficult!*  
*We should search for all signals that provide information*

- A detection should:
  - give independent evidence for dark matter scattering
  - point strongly to a spin-dependent interaction
  - help with mass reconstruction (because of different kinematics)

# Why now?

*We can accurately quantify the signal and background*

- Structure functions known (needed for cross-section)
- Backgrounds are more-or-less known
- Future detector properties are more-or-less known

# An old idea... Inelastic scattering

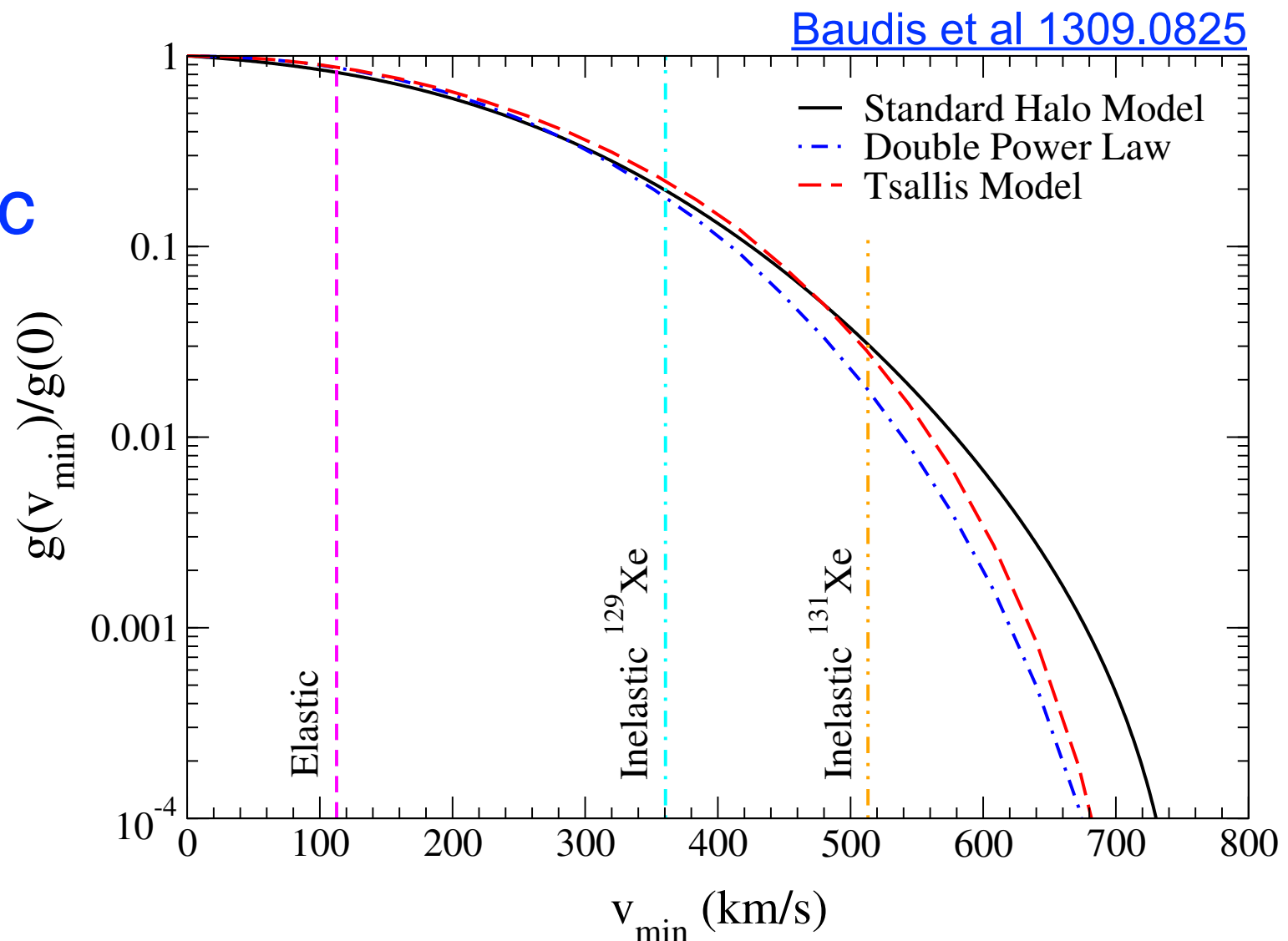
Can it ever be detected?

# Scattering rate

- Rate depends on the DM velocity distribution:

$$\frac{dR}{dE_R} \propto g(v_{\min}) = \int_{v_{\min}} d^3v \frac{f(v)}{v}$$

- $v_{\min}$  is higher for inelastic  
(DM kinetic energy must also excite the nucleus)
- This suppresses  
the inelastic rate  
by factor  $\sim 10$



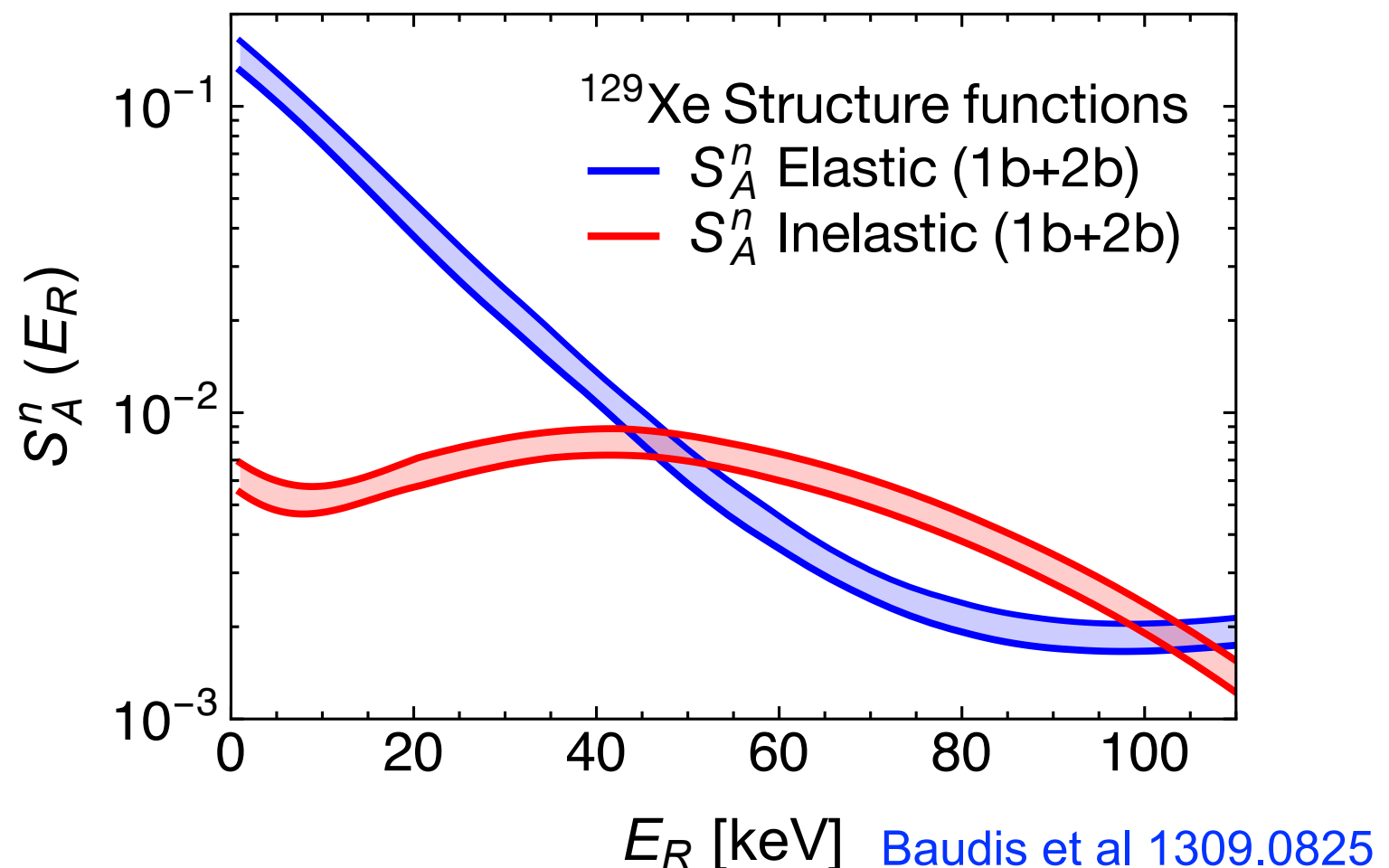


# Structure functions

- Known for axial-vector interaction:  $\mathcal{L} \propto -\bar{\chi}\gamma^\mu\gamma^5\chi \cdot \sum_q A_q \bar{\psi}_q\gamma_\mu\gamma^5\psi_q$
- Rate depends on the structure functions

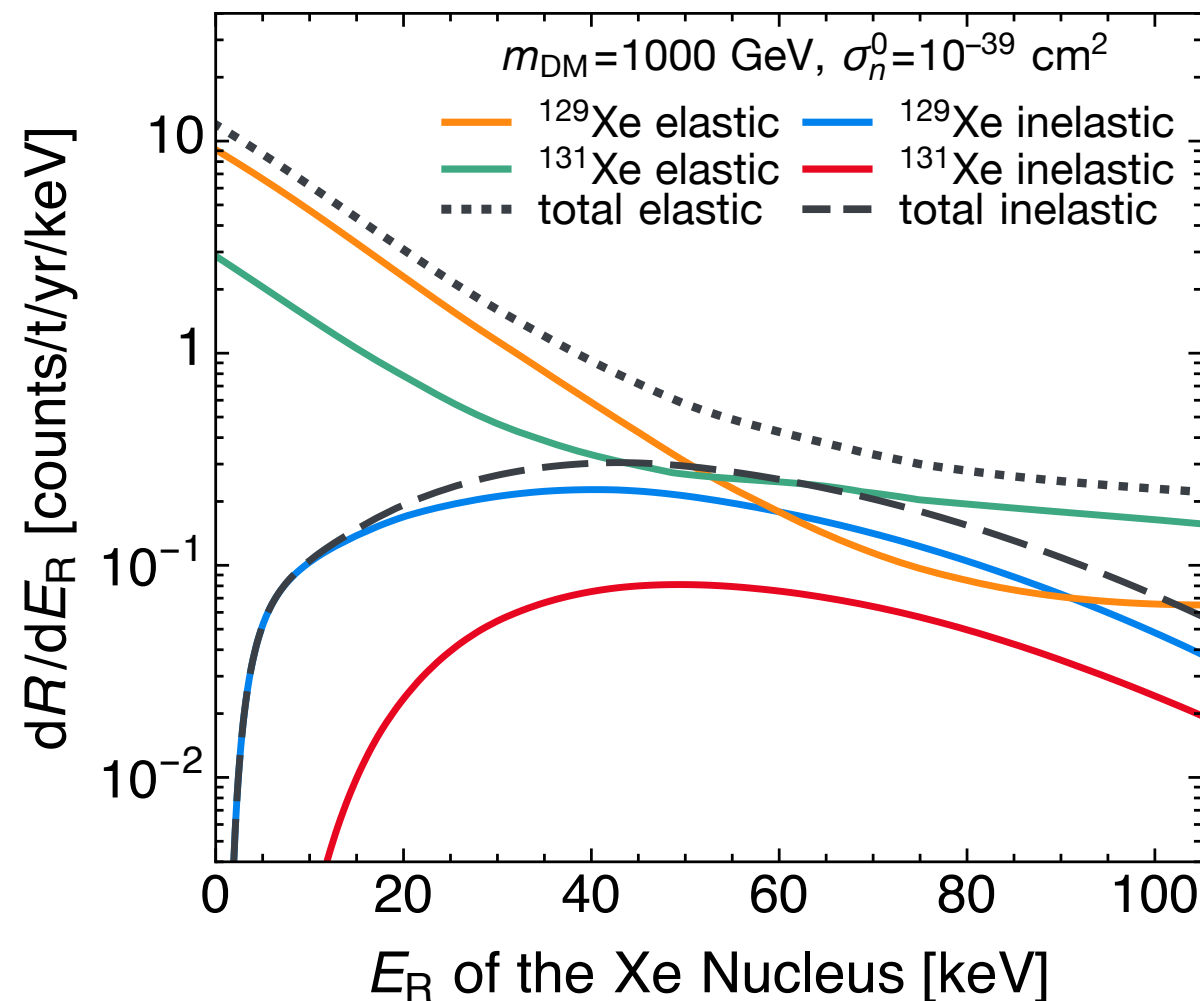
$$\frac{dR}{dE_R} \propto \frac{d\sigma}{dE_R} \propto S_A^n = \left| \langle \text{Xe}^* | \bar{\psi}_q \gamma_\mu \gamma^5 \psi_q | \text{Xe} \rangle \right|^2$$

- **Smaller** for inelastic  
(Small  $E_R$  most relevant)
- This **suppresses**  
the inelastic rate  
by factor  **$\sim 10$**



# The rate

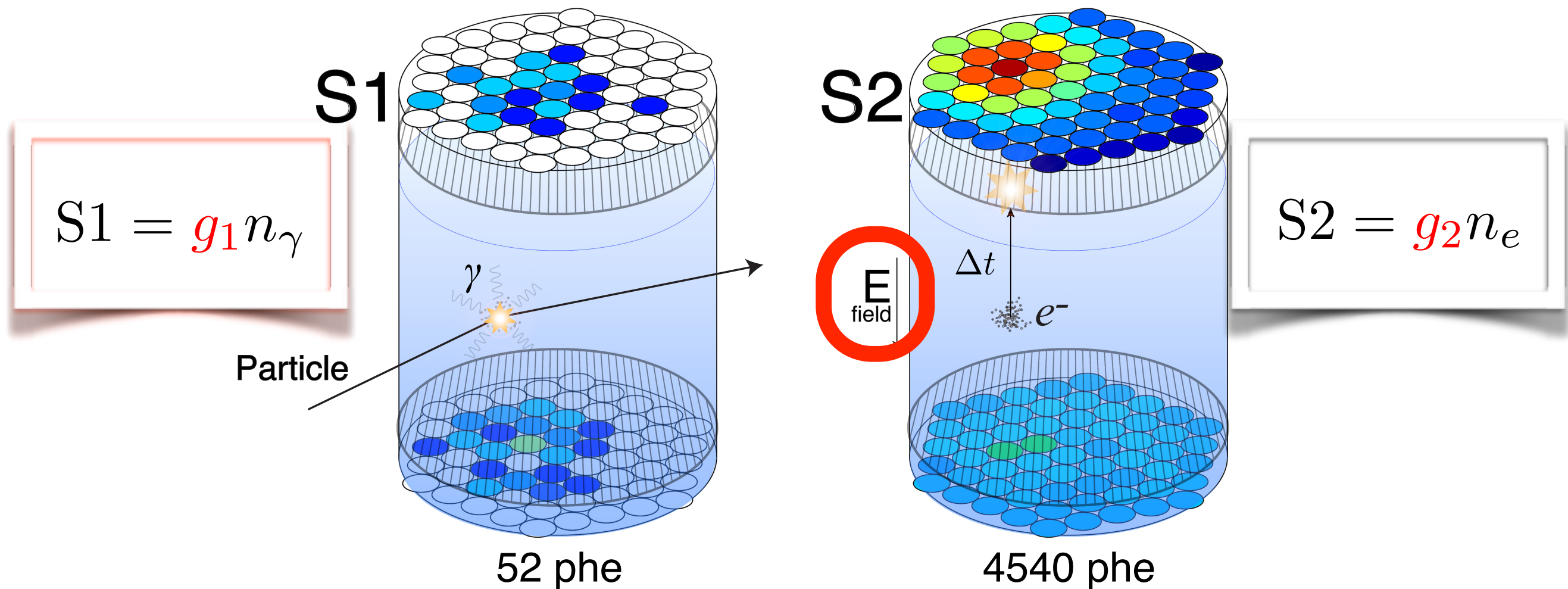
- Rate as a function recoil energy (not directly measured)



- Inelastic rate **smaller** by factor  **$\sim 100$**   
→ **Always see an elastic signal first**

# Two-phase xenon detectors

- Express the signal in terms of measured quantities:



$g_1$ ,  $g_2$  and drift field are the crucial parameters

# Mock detectors

- I'll consider two benchmark scenarios:

*XenonA200*

$g_1 = 0.07$  PE/ $\gamma$

$g_2 = 12.5$  PE/e  
(50% extraction efficiency)

drift field = 200 V/cm

*XenonB1000*

$g_1 = 0.12$  PE/ $\gamma$

$g_2 = 50$  PE/e  
(100% extraction efficiency)

drift field = 1000 V/cm

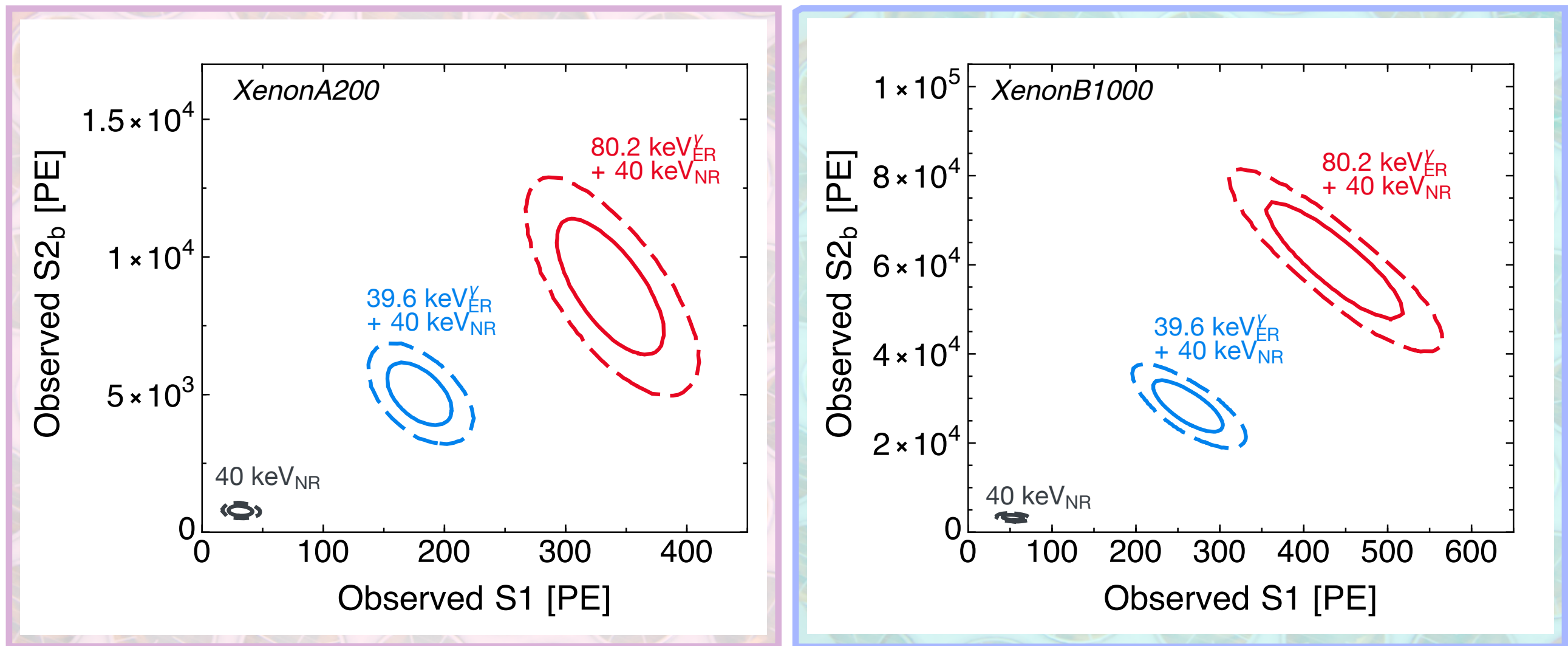
- Number of photons & electrons modelled with NEST

[Szydagis et al 1106.1613](#)



# Mock signals

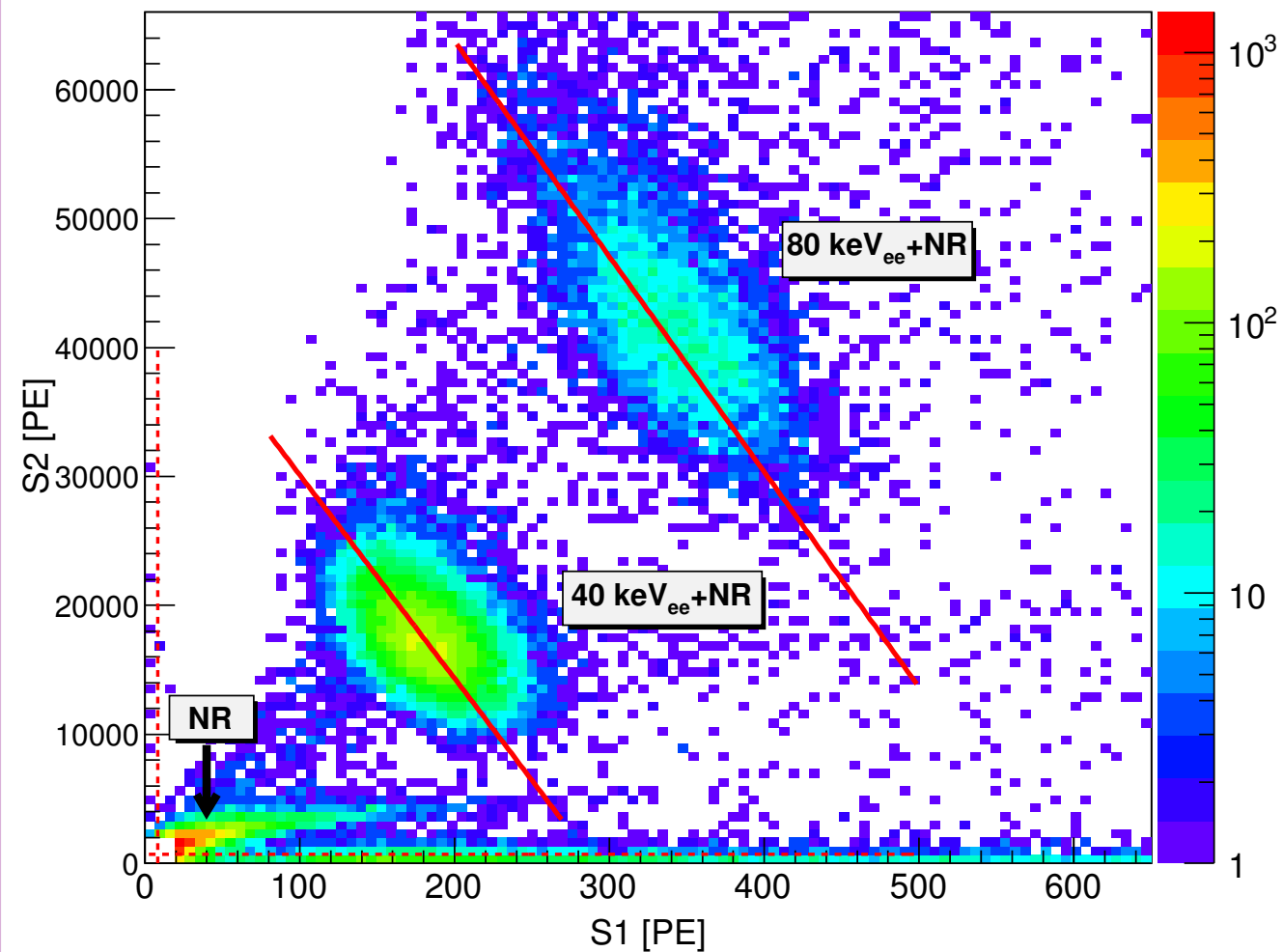
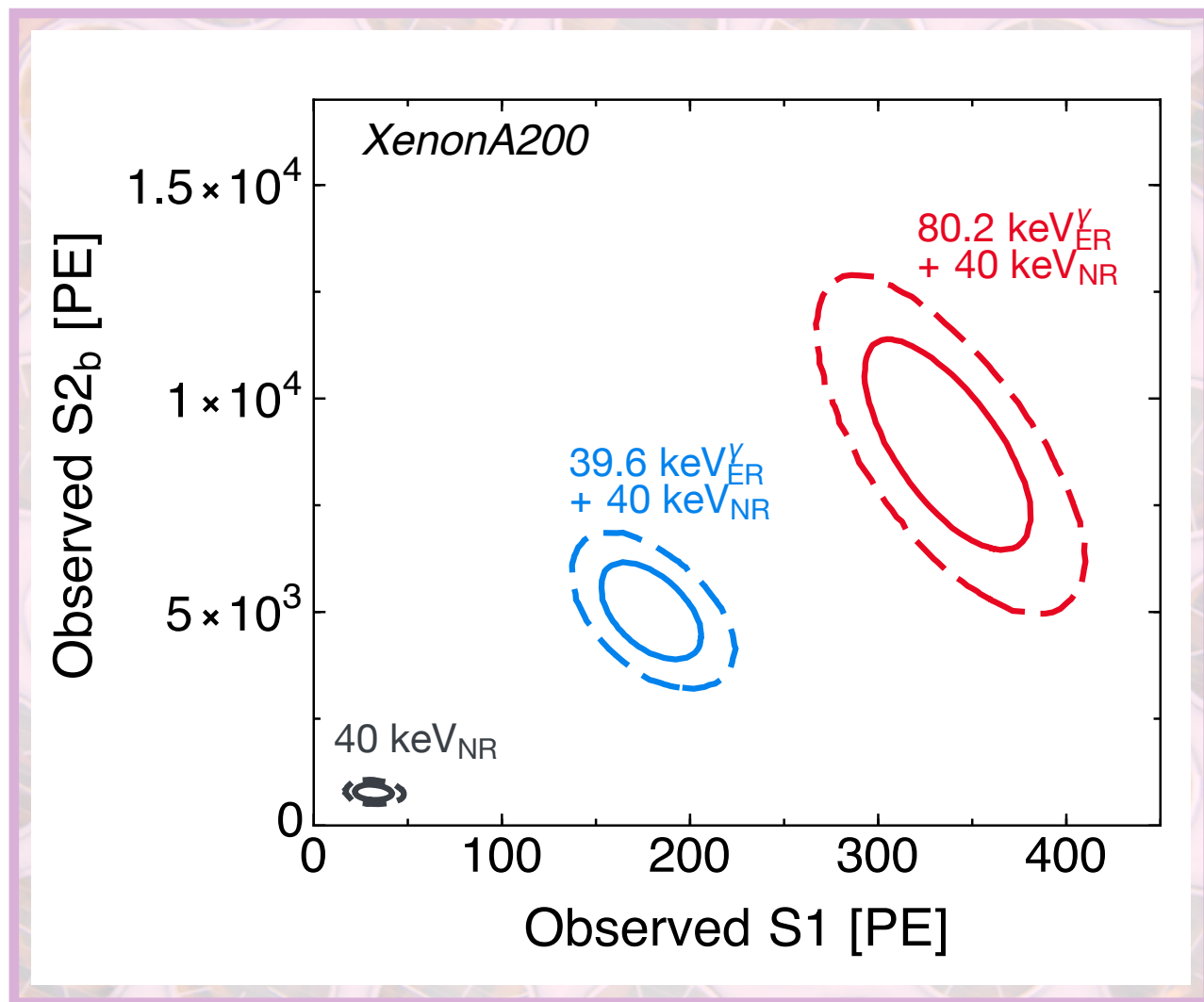
- Include **detector** and **recombination** fluctuations



- For same energy, electronic recoils produce a *much larger*  $S1$  and  $S2$

# Mock signals

- Looks like real data... 😊

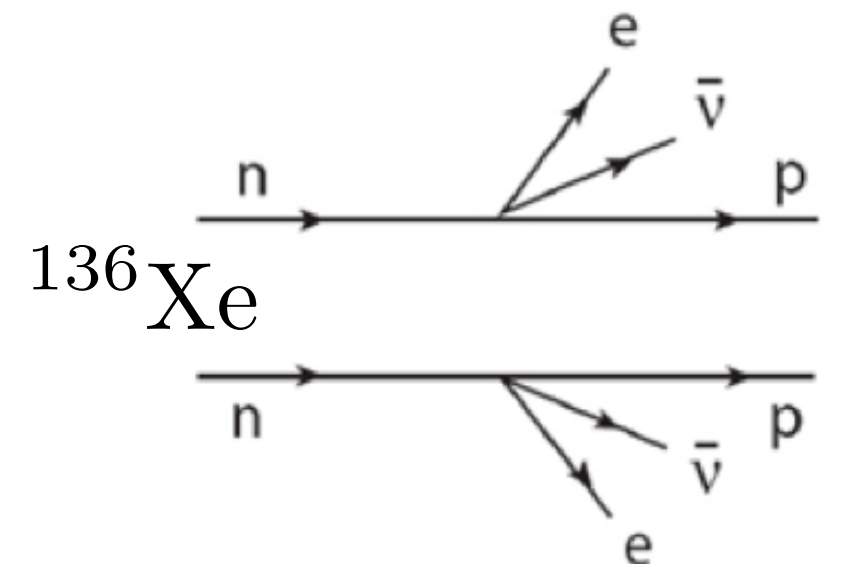
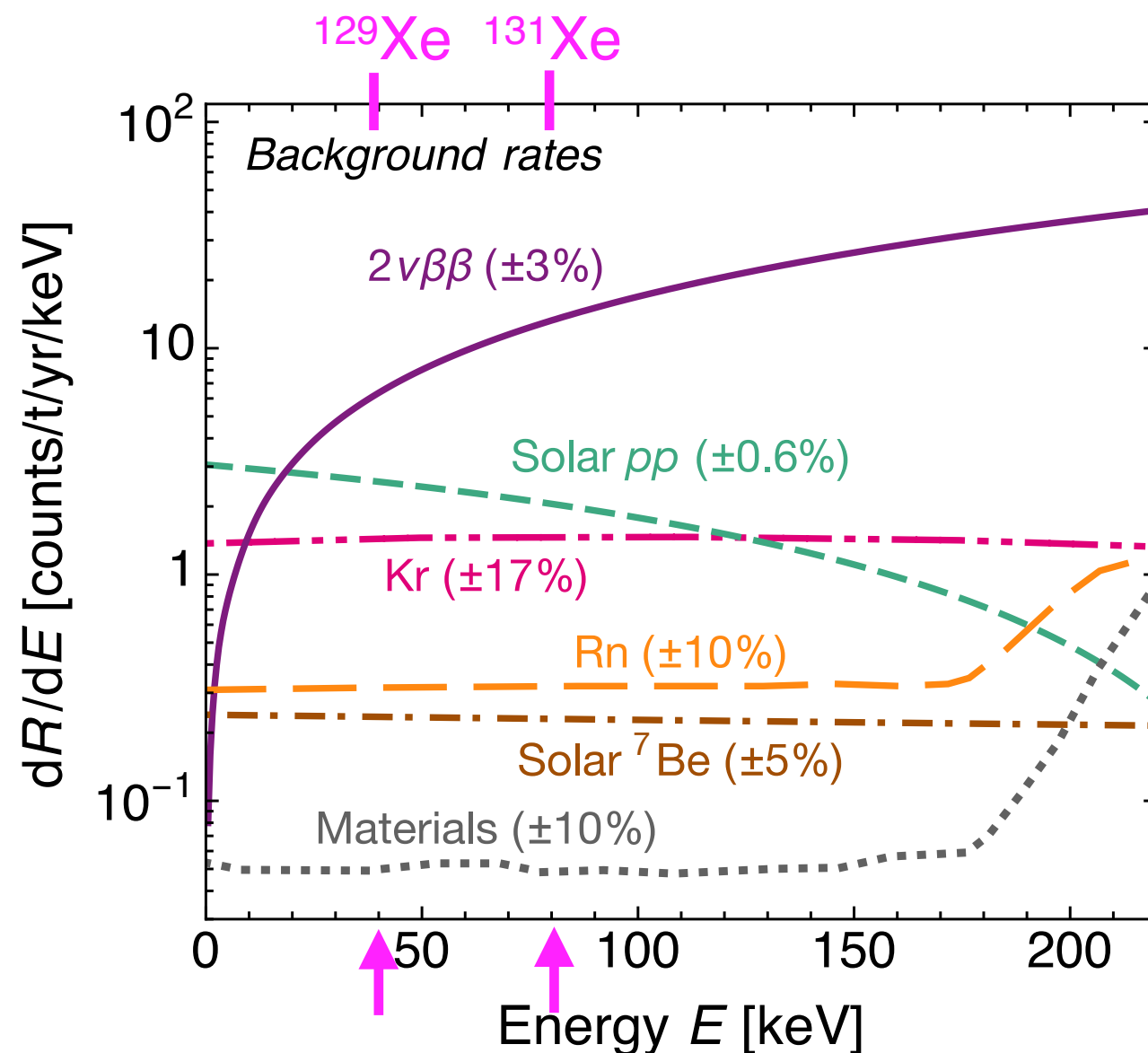


Data from PandaX-I [arXiv:1505.00771](https://arxiv.org/abs/1505.00771)

# Background

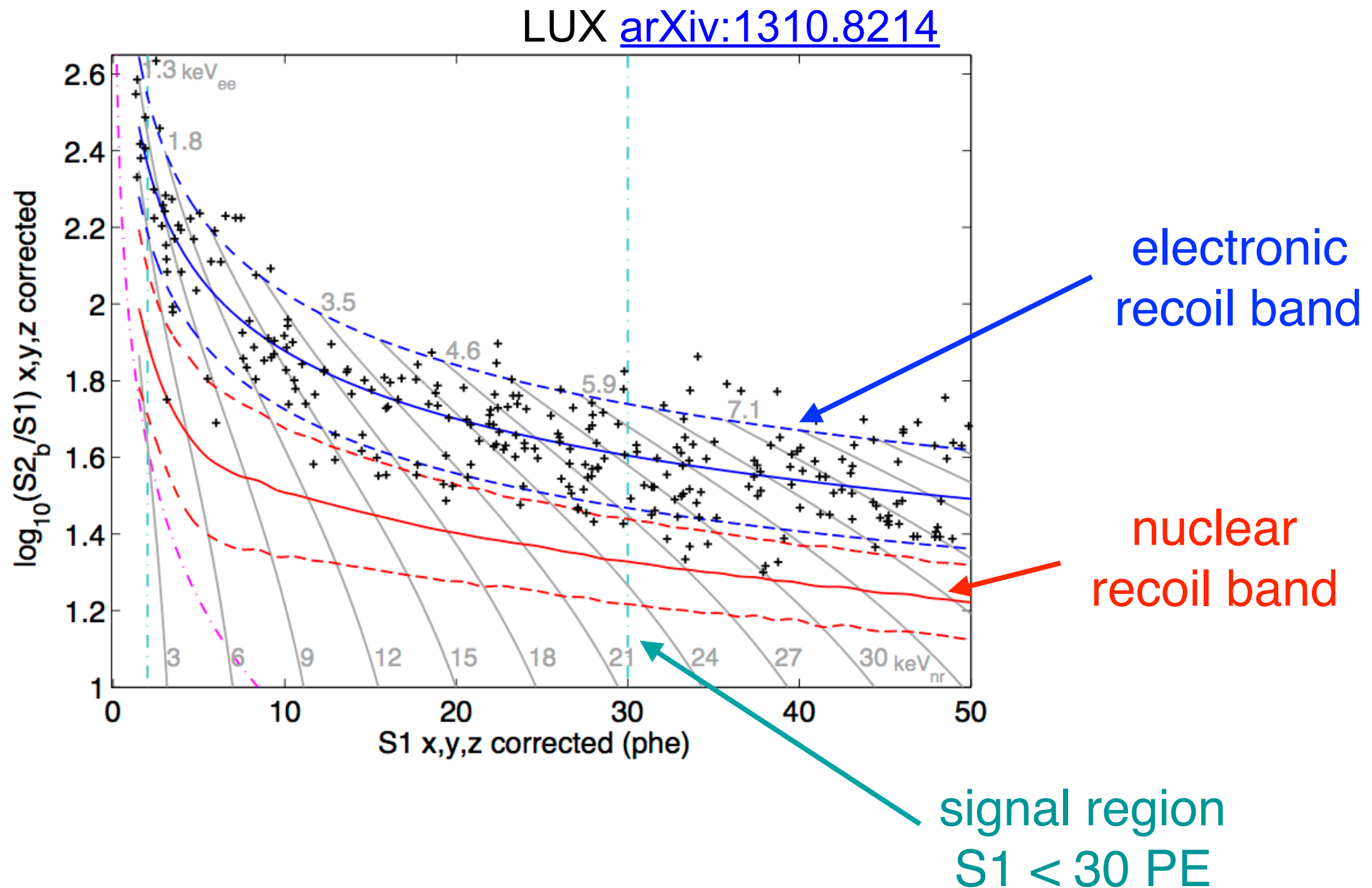
- Background spectra expected in LZ/XENONnT:

[LZ Design: 1509.02910](#)



- 2-neutrino — 2-beta decay of  ${}^{136}\text{Xe}$  dominates above 20 keV*

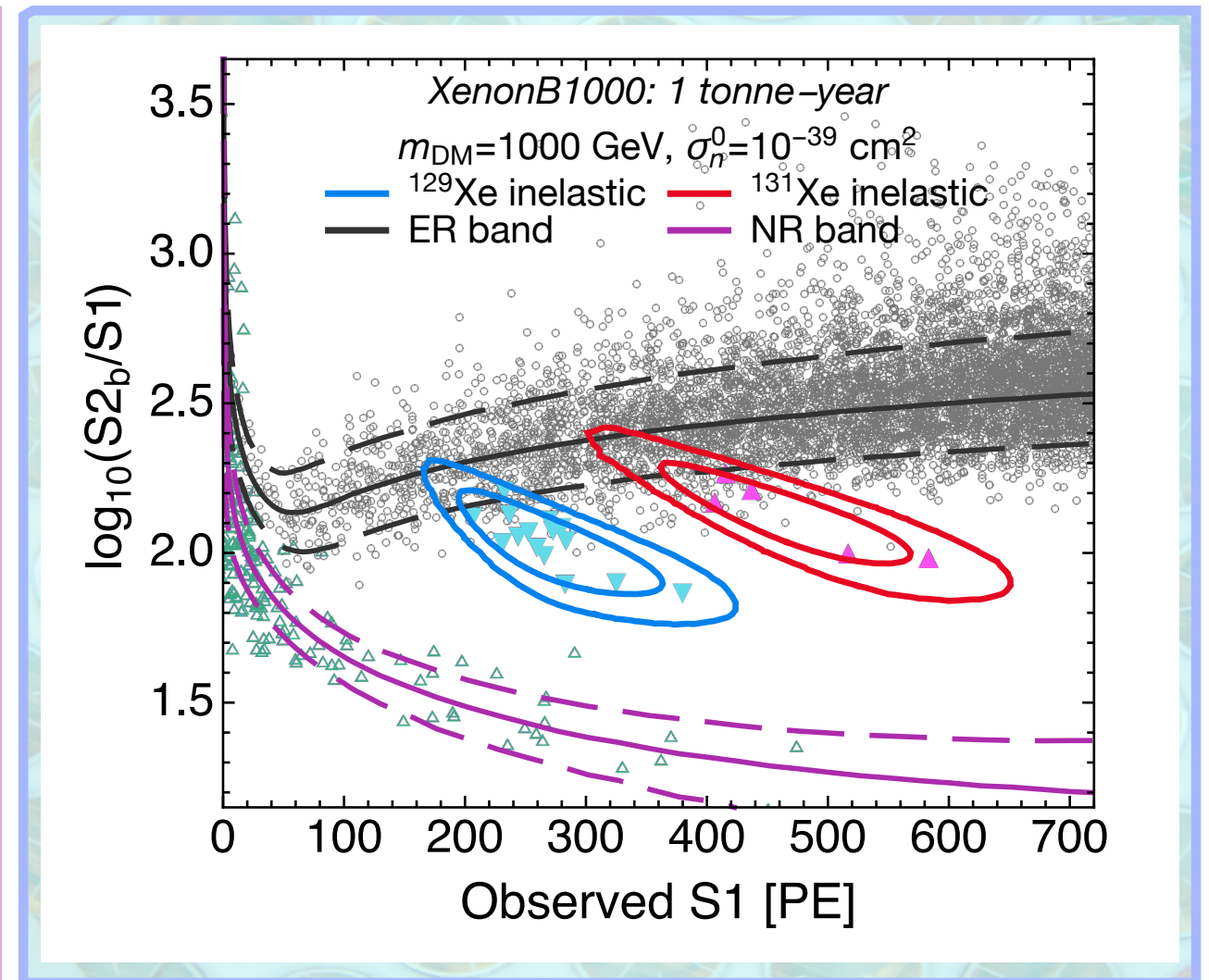
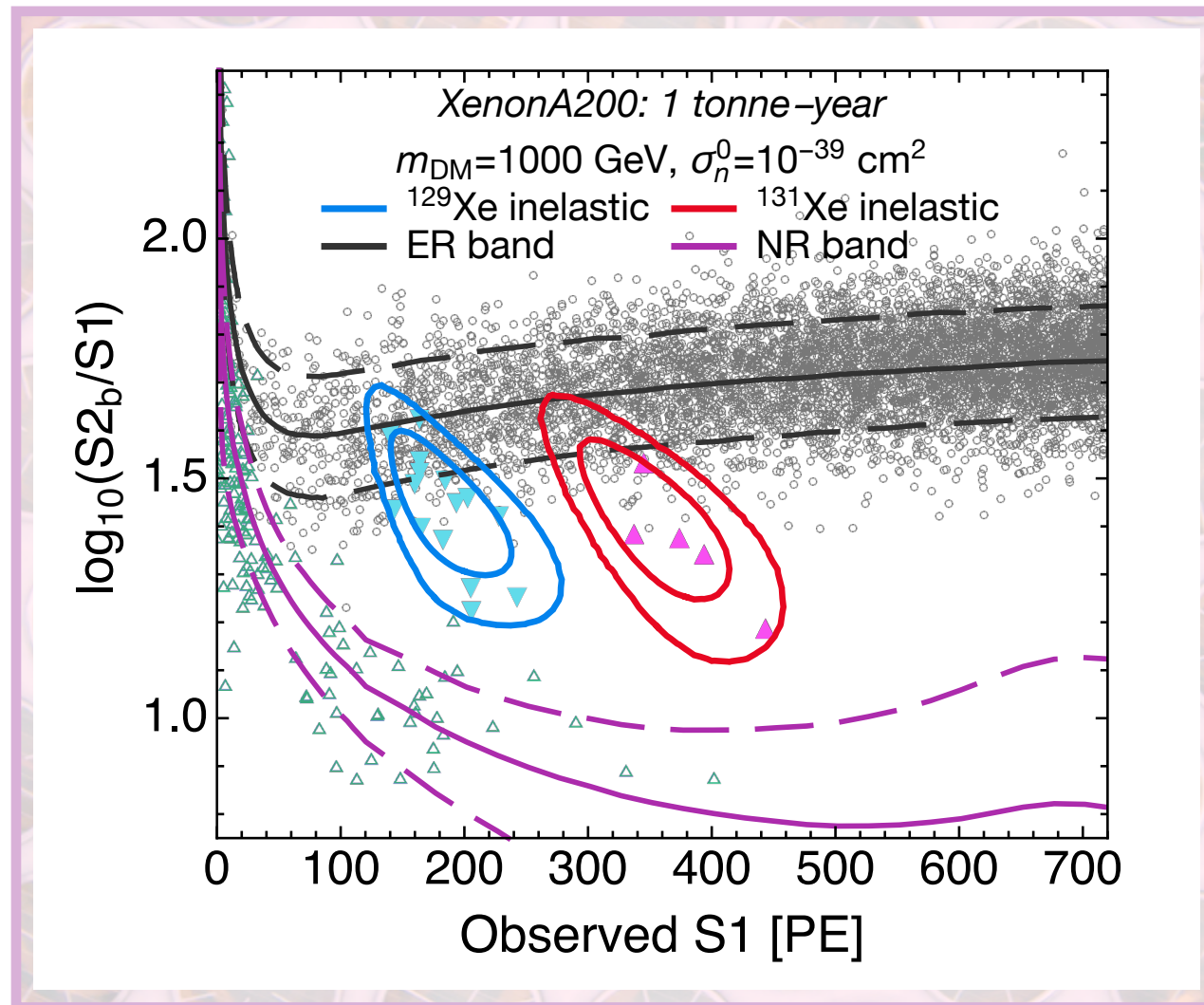
# Reminder: Usual signal plane





# Background versus signal

- Signal region at *higher values* of S1



- Large backgrounds...but some signal-to-background discrimination
- Better discrimination for higher drift fields

# Discovery limit

- Quantify the sensitivity of future experiments with a ‘discovery limit’ [Billard et al 1110.6079](#)

*The smallest cross-section at which 90% of experiments can make a  $3\sigma$  detection of the signal*

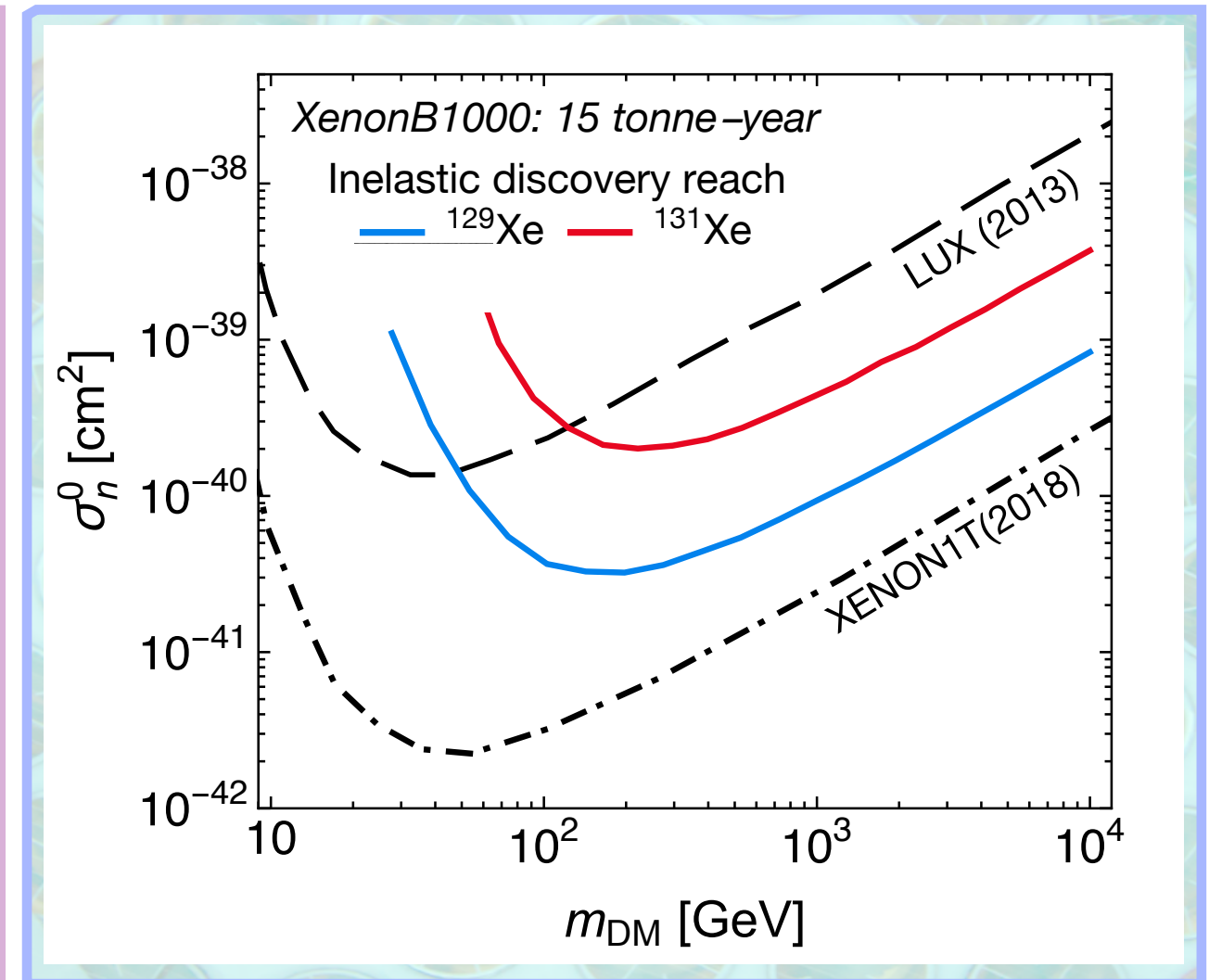
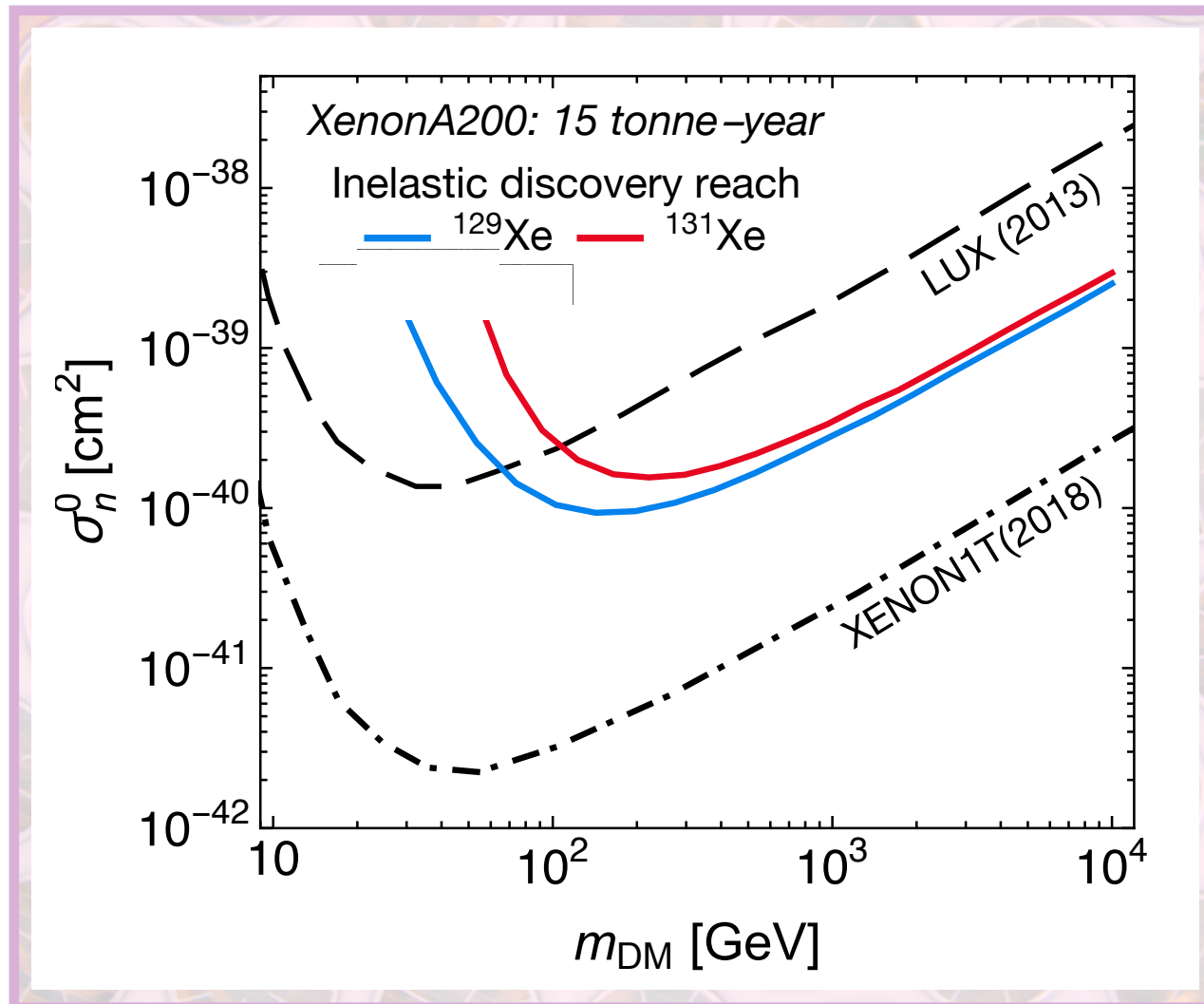
- Profile likelihood ratio:

$$\lambda(0) = \frac{L(\sigma_n^0 = 0, \hat{\vec{A}}_{\text{BG}})}{L(\hat{\sigma}_n^0, \hat{\vec{A}}_{\text{BG}})}$$

- Include background uncertainties

# Discovery limit

- Compare discovery limit with current/future (elastic) constraints



- Detectable if XENON1T make discovery in next run

# Summary

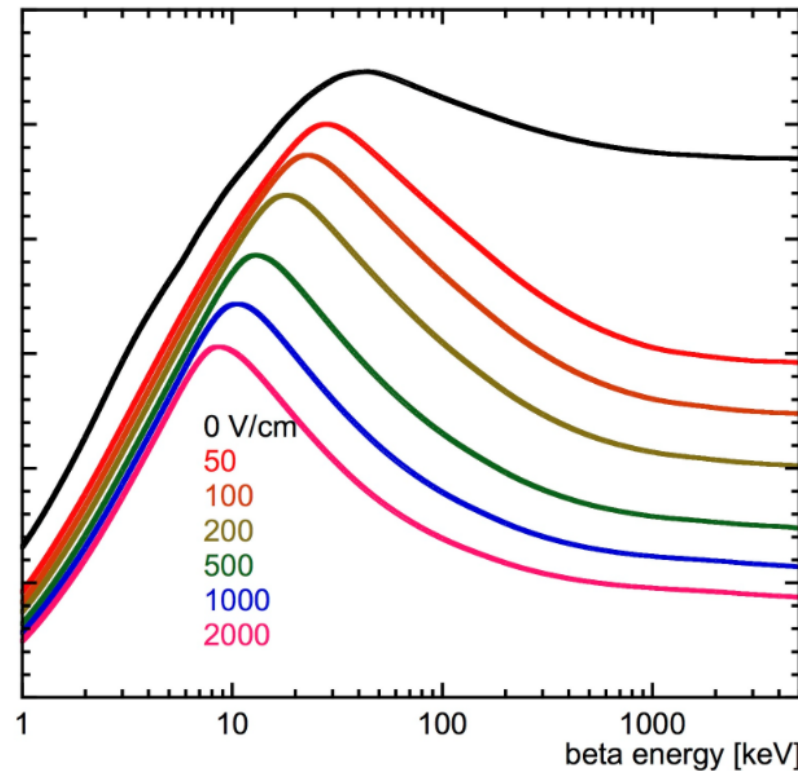
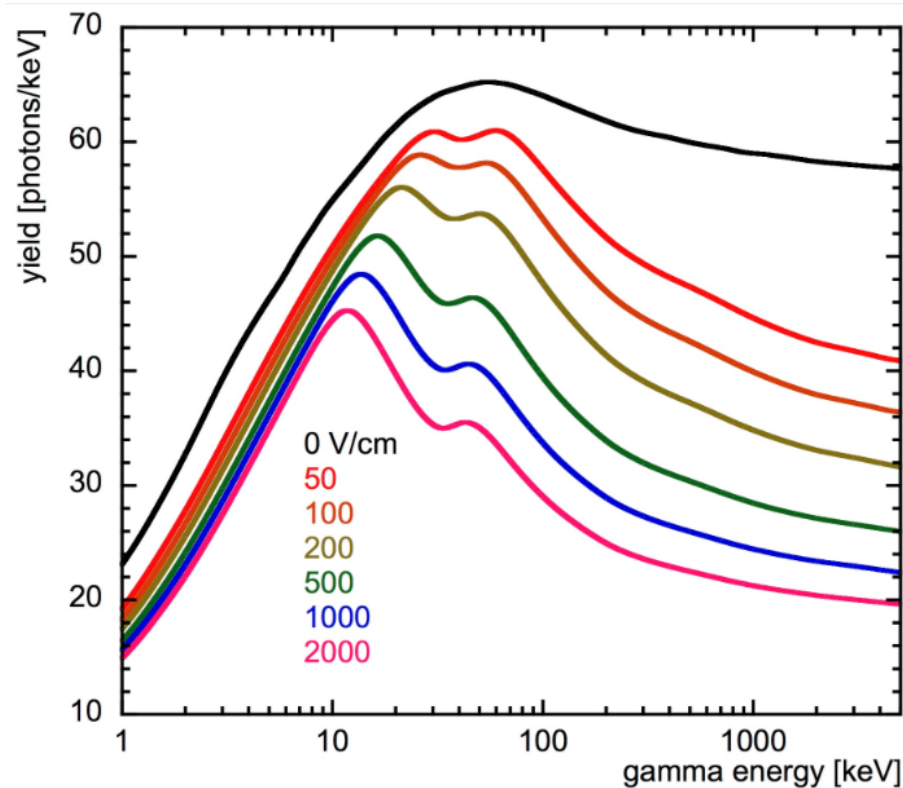
- Dark matter can excite the  $^{129}\text{Xe}$  and  $^{131}\text{Xe}$  isotopes
  - ➔ signal is nuclear recoil + photon
- Signal is always smaller than elastic rate
  - ➔ Can it be detected?

Yes!

...need an (elastic) discovery signal  
in the next run of XENON1T

Thank you

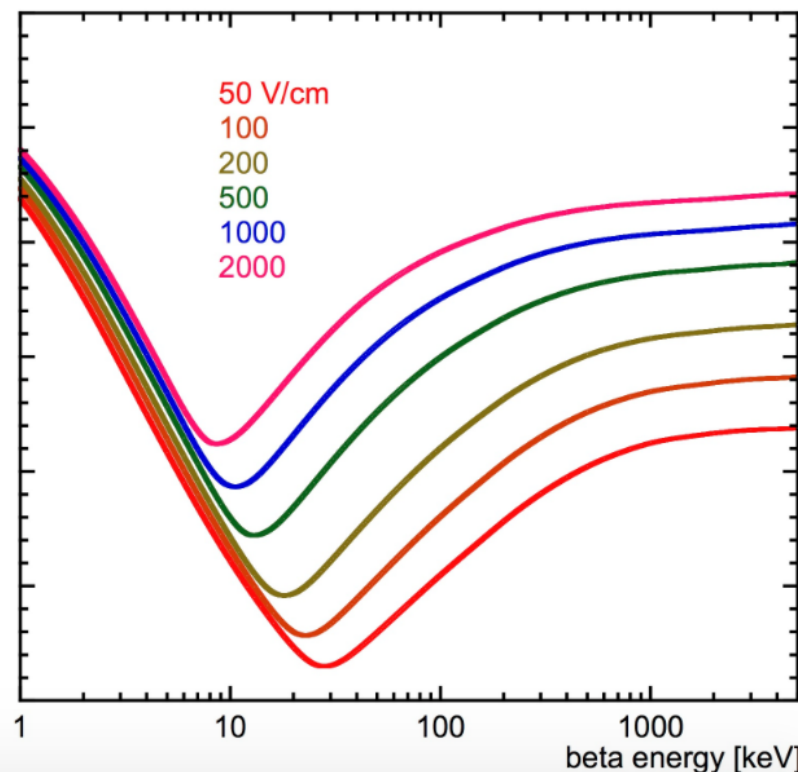
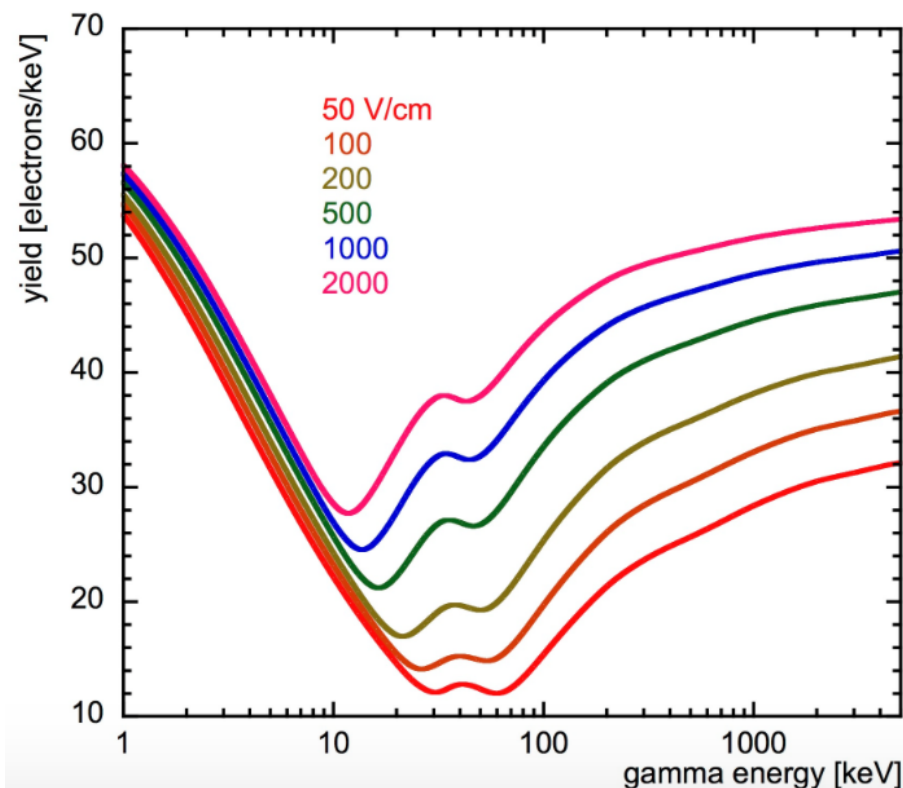
# Backup



Gammas have shorter tracks,  
more recombination ( $r$  bigger)  
so  $n_e$  smaller,  $n_{\text{gamma}}$  bigger

$$n_e = n_i - r n_i$$

$$n_{\gamma} = n_{\text{ex}} + r n_i$$



$$\lambda(0) = \frac{L(\sigma_n^0 = 0, \hat{\vec{A}}_{\text{BG}})}{L(\hat{\sigma}_n^0, \hat{\vec{A}}_{\text{BG}})}$$

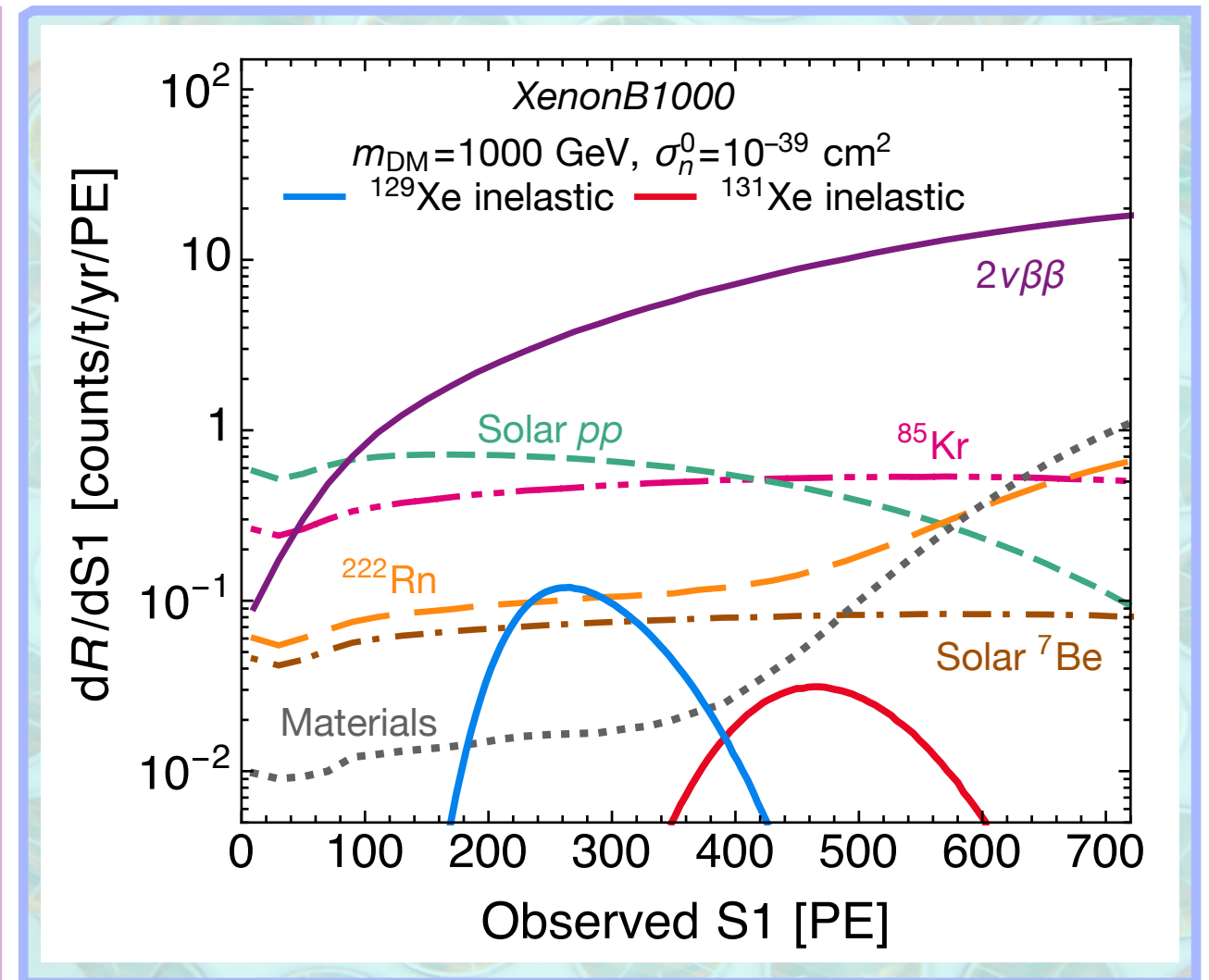
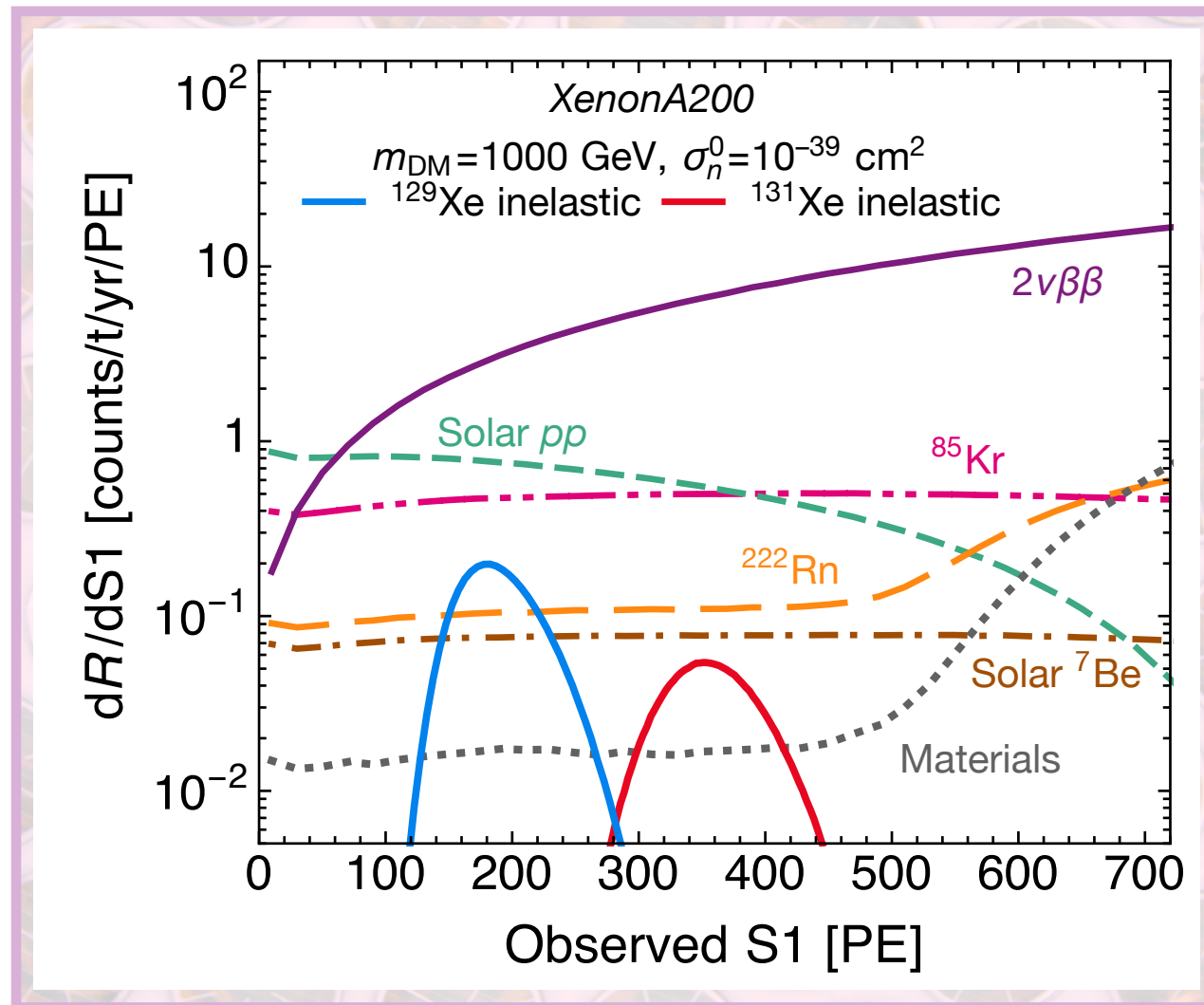
$$L(\sigma_n^0, \vec{A}_{\text{BG}}) = \frac{\left(\mu_{\text{DM}} + \sum_{j=1}^6 \mu_{\text{BG}j}\right)^N}{N!} \exp\left(-\mu_{\text{DM}} + \sum_{j=1}^6 \mu_{\text{BG}j}\right) \cdot \prod_{m=1}^6 L_m(A_{\text{BG}m})$$

$$\cdot \prod_{i=1}^N \left[ \frac{\mu_{\text{DM}}}{\mu_{\text{DM}} + \sum_{k=1}^6 \mu_{\text{BG}k}} f_{\text{DM}}(\text{S1}_i, \log_{10}(\text{S2}_b/\text{S1})_i) \right. \\ \left. + \sum_{j=1}^6 \frac{\mu_{\text{BG}j}}{\mu_{\text{DM}} + \sum_{k=1}^6 \mu_{\text{BG}k}} f_{\text{BG}j}(\text{S1}_i, \log_{10}(\text{S2}_b/\text{S1})_i) \right],$$



# Single-phase experiments

- Detecting this signal could be difficult...



...impossible for single phase (S1-only)?

# Improvements?

- Could have a larger exposure
  - ➔ background dominated so only scales with the square root
- Could reduce backgrounds
  - Largest: 2-beta—2-neutrino decay of  $^{136}\text{Xe}$ 
    - ➔ Remove the  $^{136}\text{Xe}$  isotope
  - Try to search for displaced the S2 signal from the recoil and photon?