

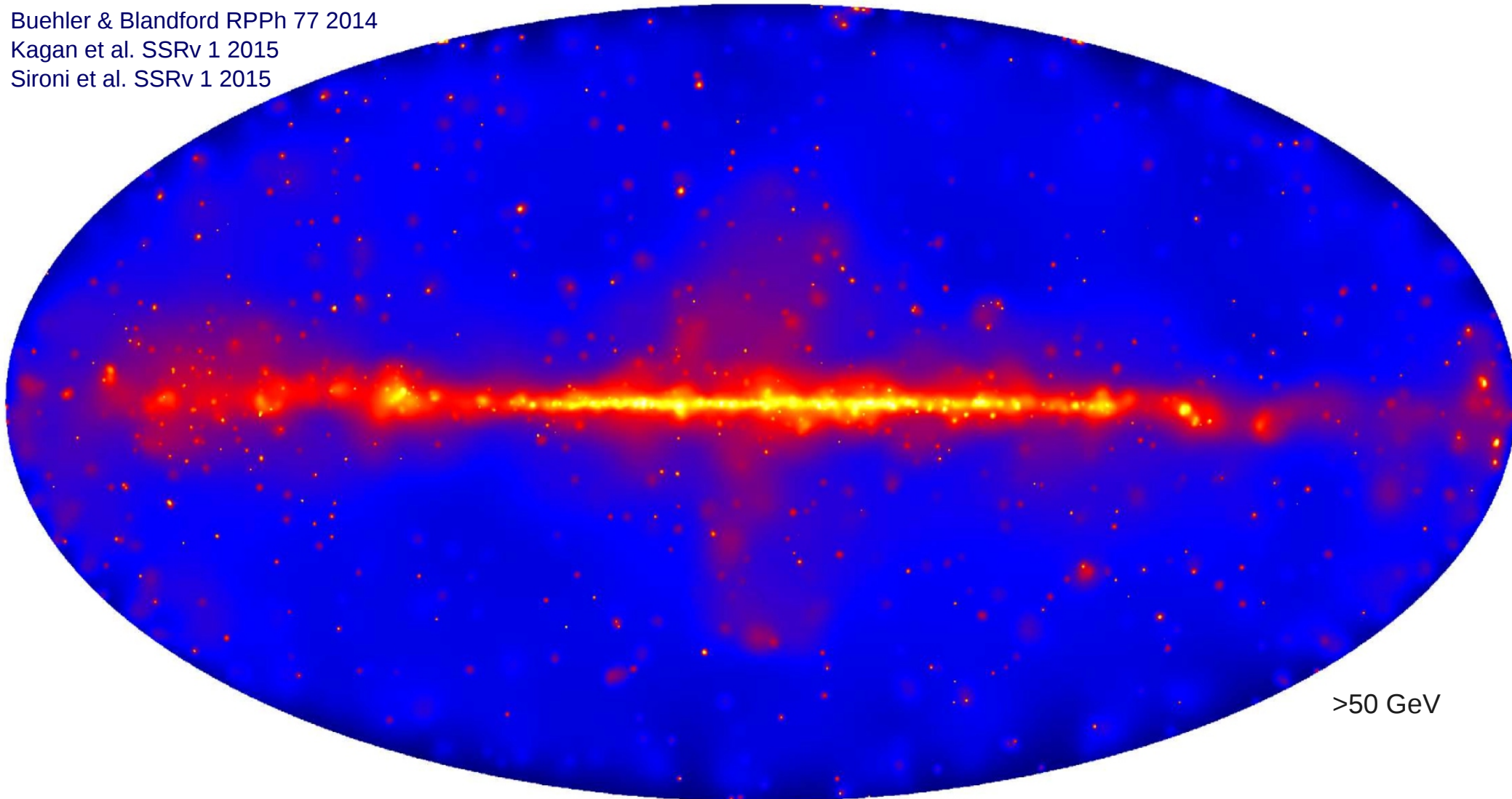
Particle acceleration observed in gamma rays

Based on

Buehler & Blandford RPPh 77 2014

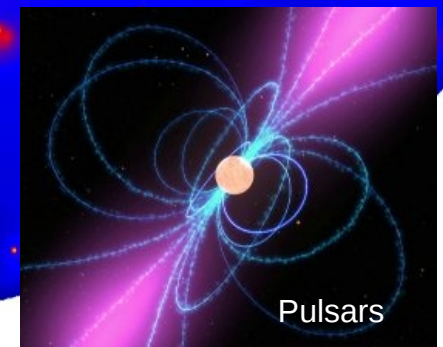
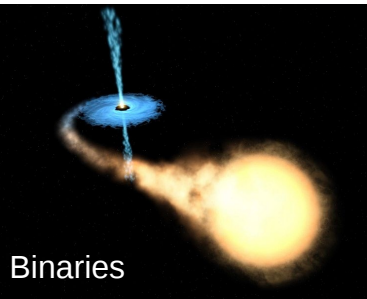
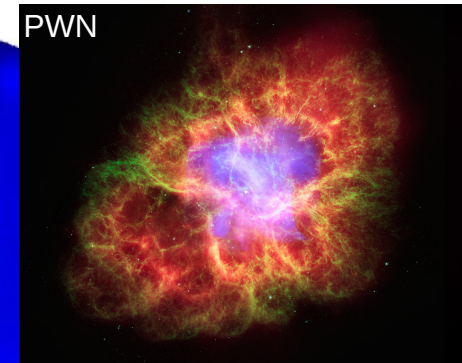
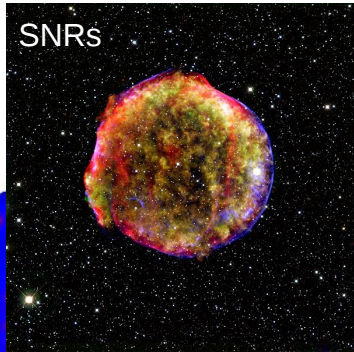
Kagan et al. SSRv 1 2015

Sironi et al. SSRv 1 2015

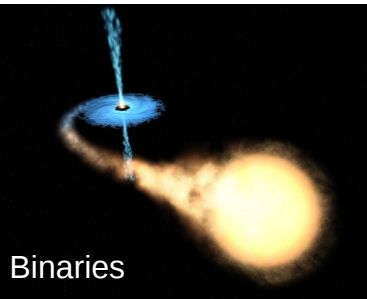
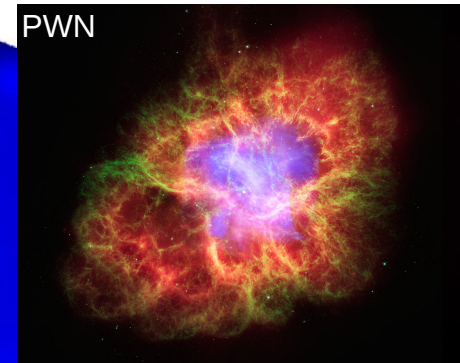
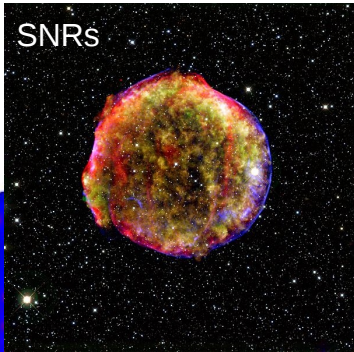


>50 GeV

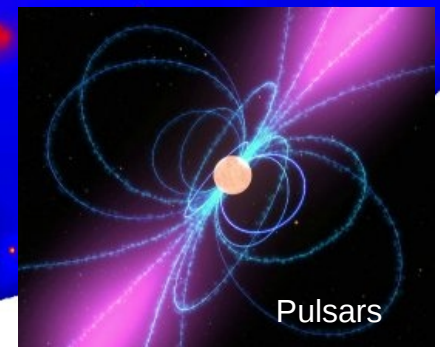
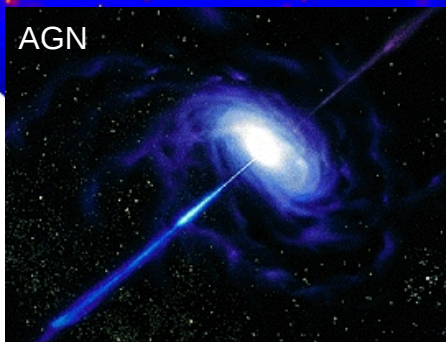
Cosmic accelerators



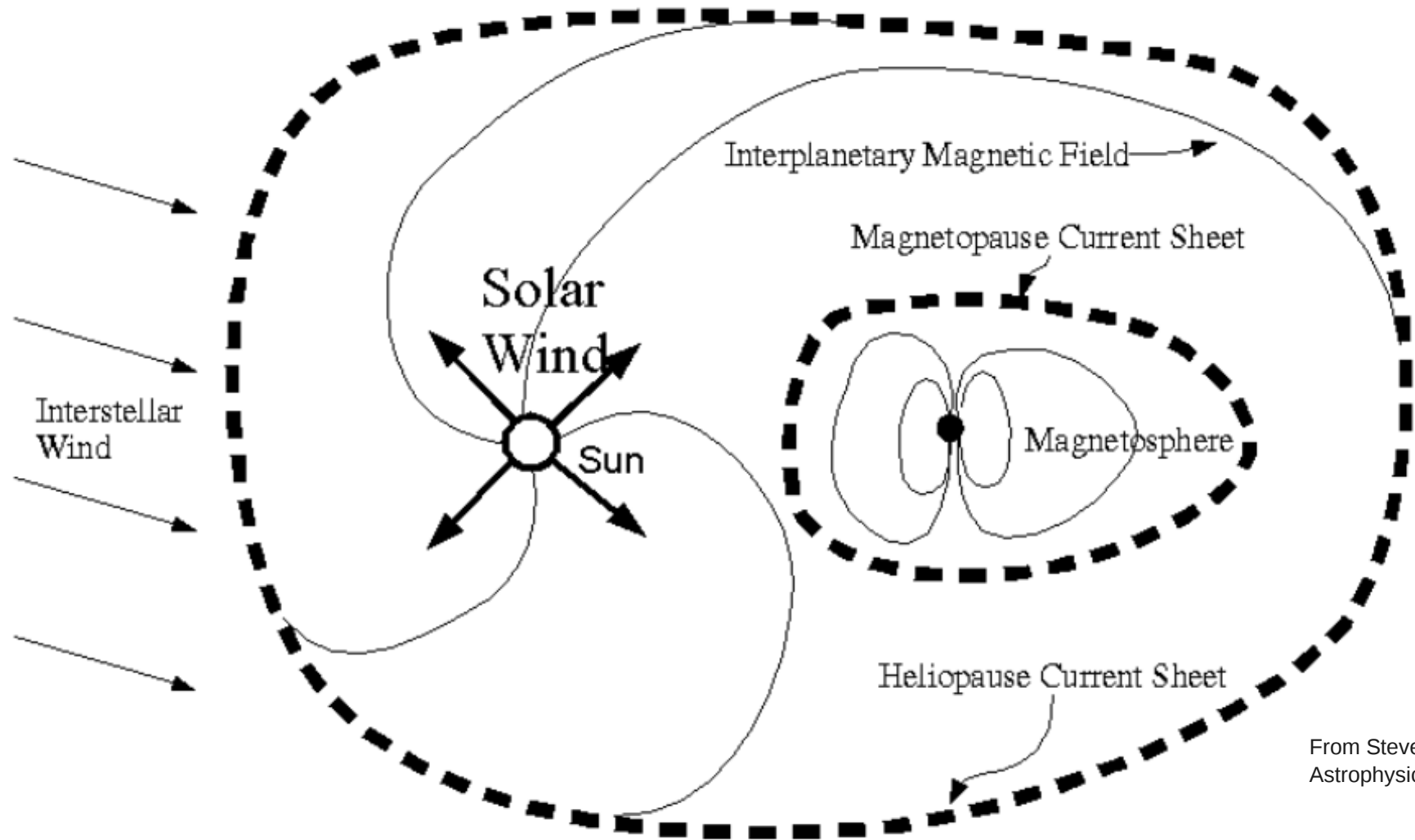
Cosmic accelerators



Why is acceleration beyond TeV energies so common?

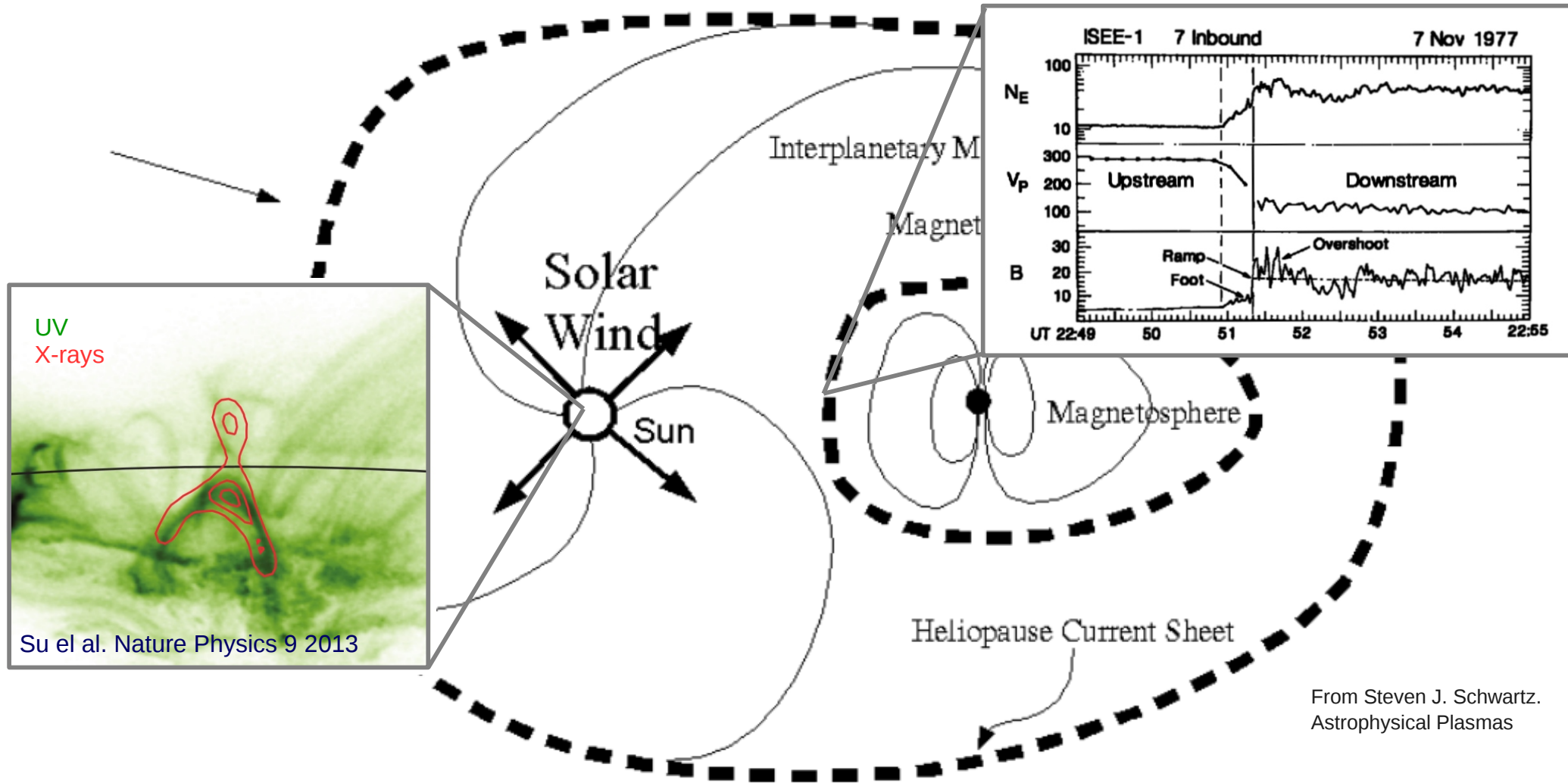


Universe of plasma bubbles



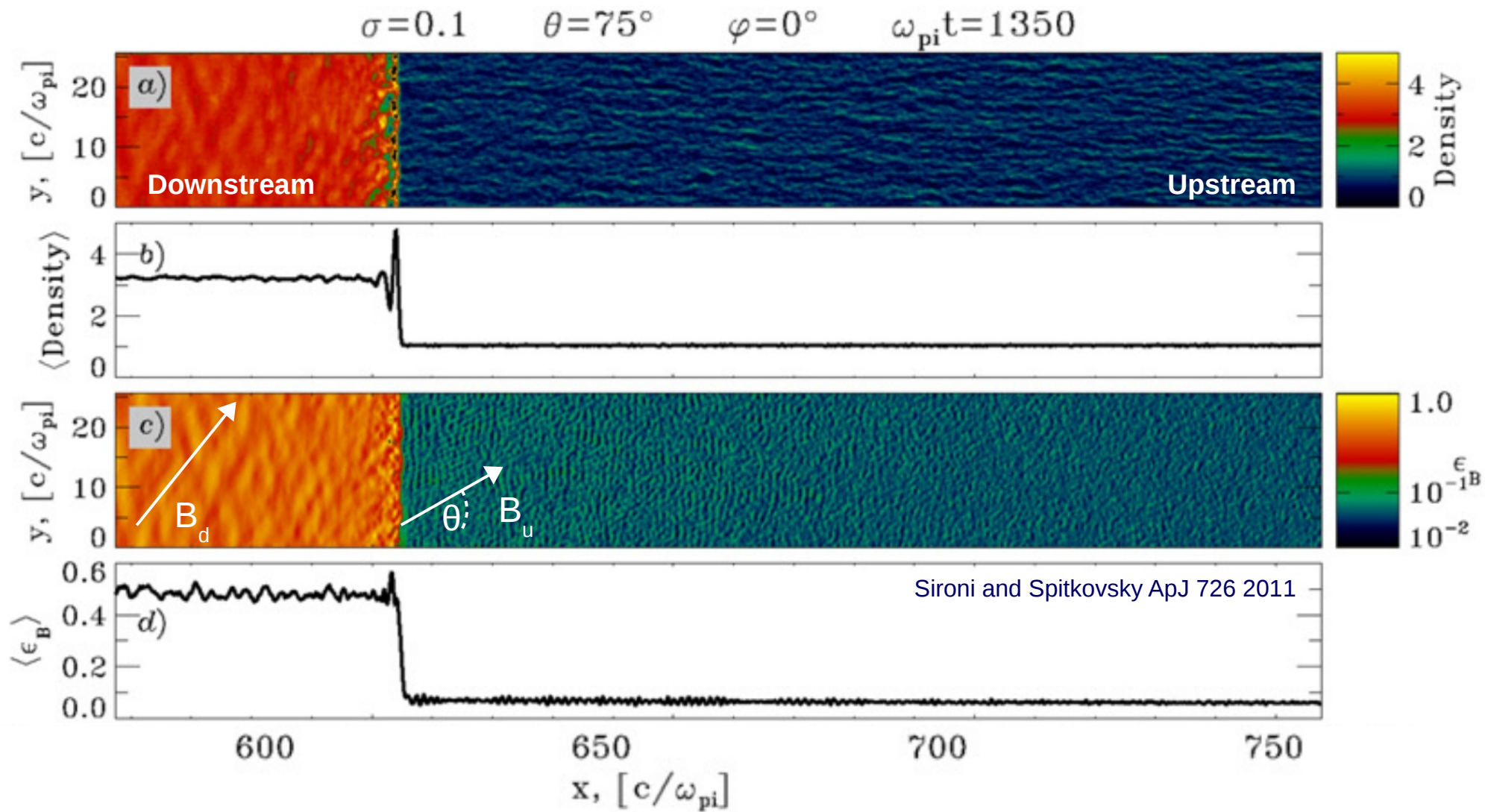
Contact of plasma flows leads to non-thermal particle acceleration,
magnetic and bulk flow energies are transferred to particles

Universe of plasma bubbles



The most studied acceleration mechanisms are shock acceleration
and magnetic reconnection

Shock acceleration



Theoretically studies since the 70's, recently with
2D/3D Particle In Cell simulations

Shock acceleration

The acceleration depends on the shock Lorentz factor and the upstream magnetization

$$\sigma = \frac{B^2}{4\pi \cdot (\gamma - 1) n m c^2}$$

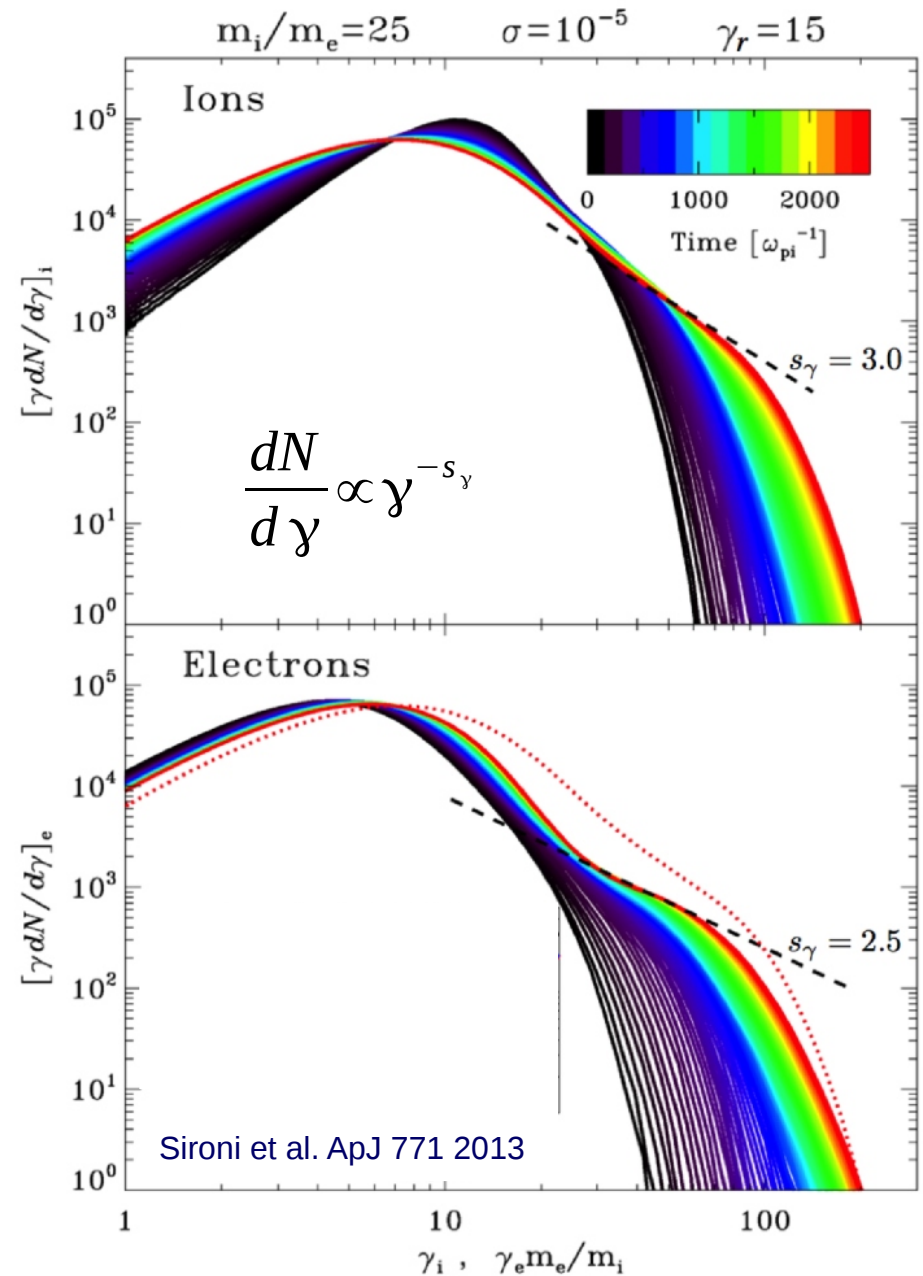
(in the shock frame)

For magnetized relativistic shocks ($\sigma > 10^{-3}$ and $\gamma > 1$) there is no acceleration if

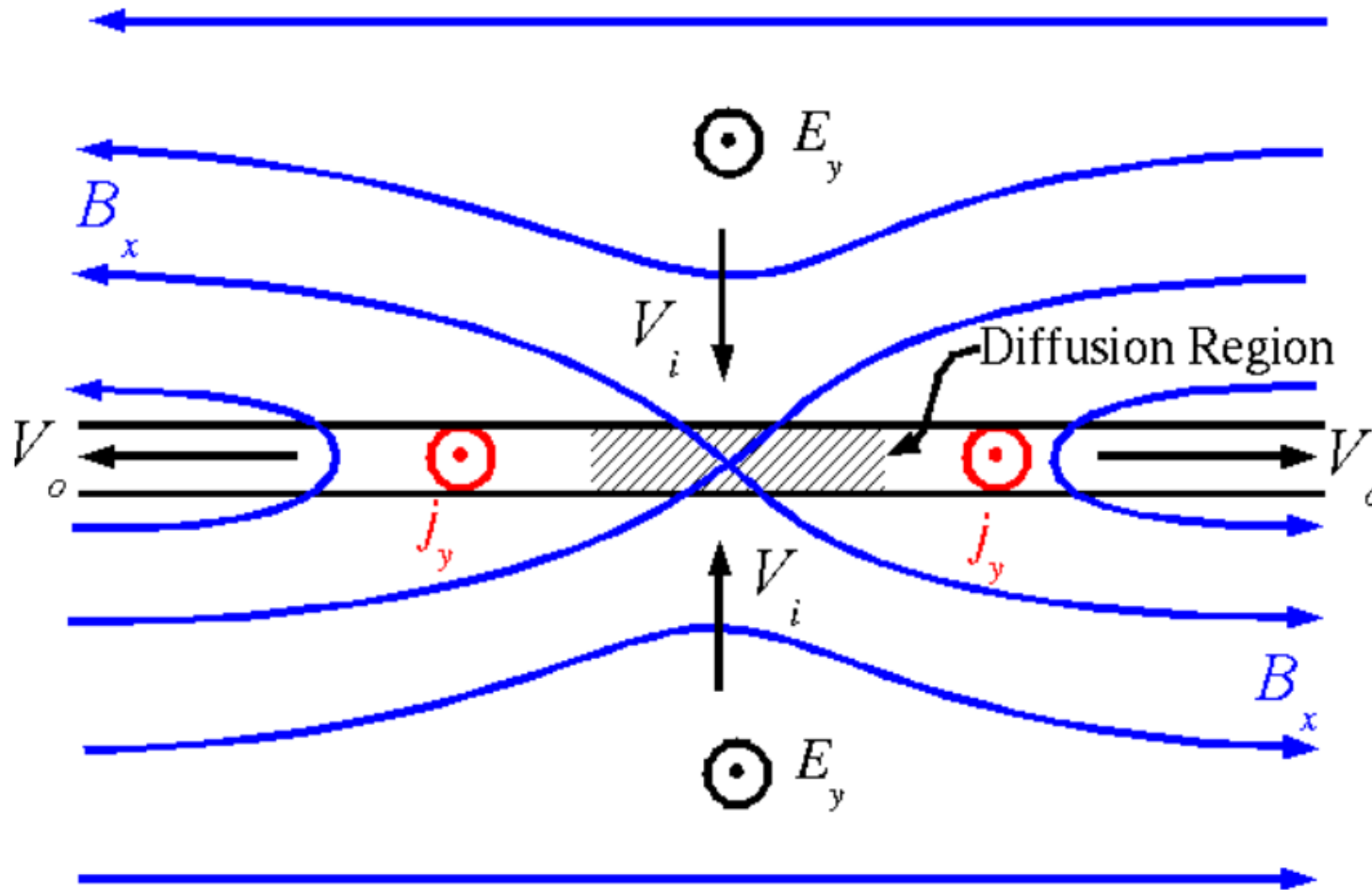
$$\theta > 34^\circ / \gamma$$

(in the upstream frame)

as particles can not come back upstream



Magnetic reconnection



$$\nabla \times \vec{B} \approx \frac{4\pi}{c} \vec{j}$$

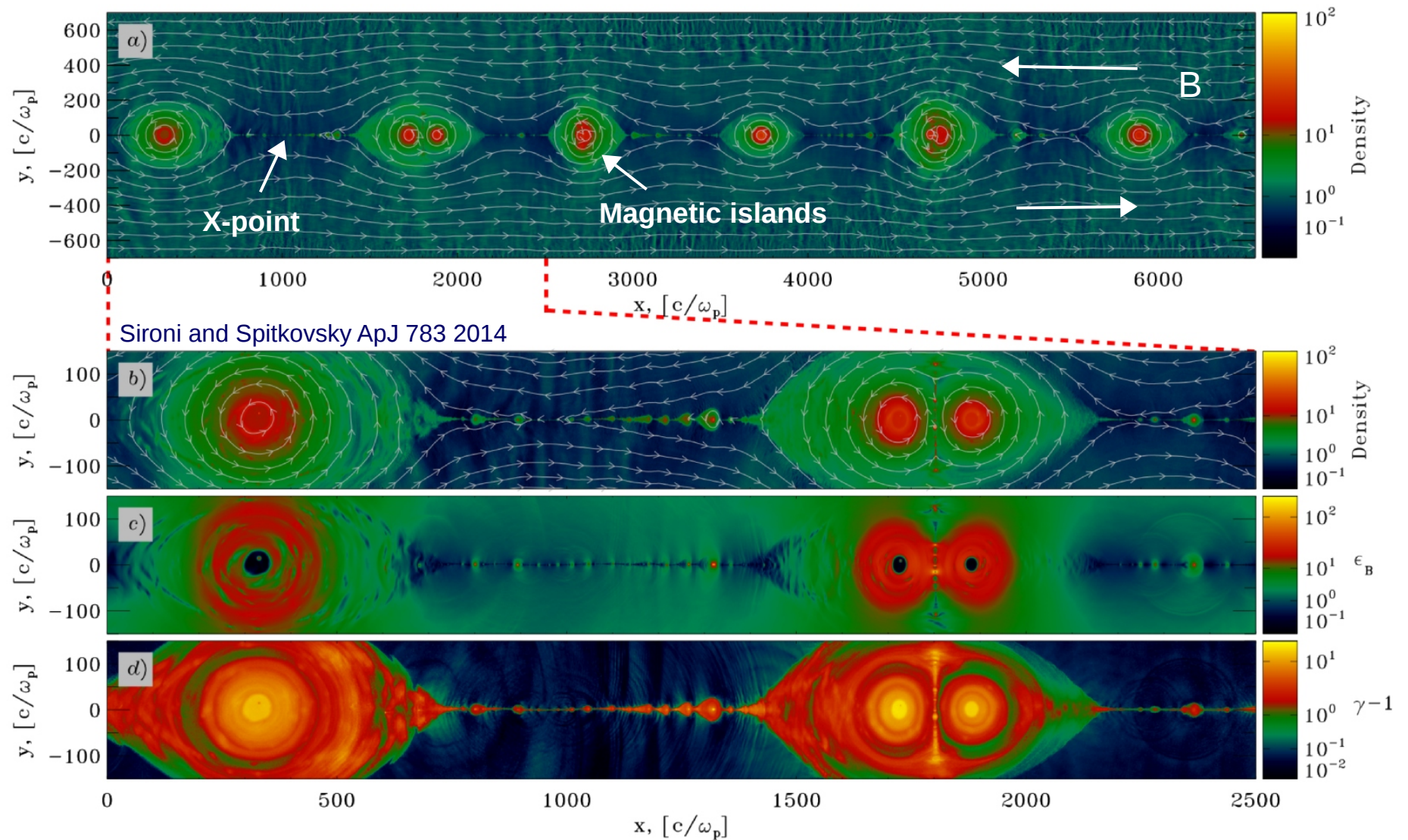
$$\vec{E} = -\frac{\vec{v}}{c} \times \vec{B}$$

$$\nabla \times \vec{E} = -\frac{1}{c} \frac{\partial \vec{B}}{\partial t} \approx 0$$

$$\nabla \cdot \vec{E} = 4\pi \rho \approx 0$$

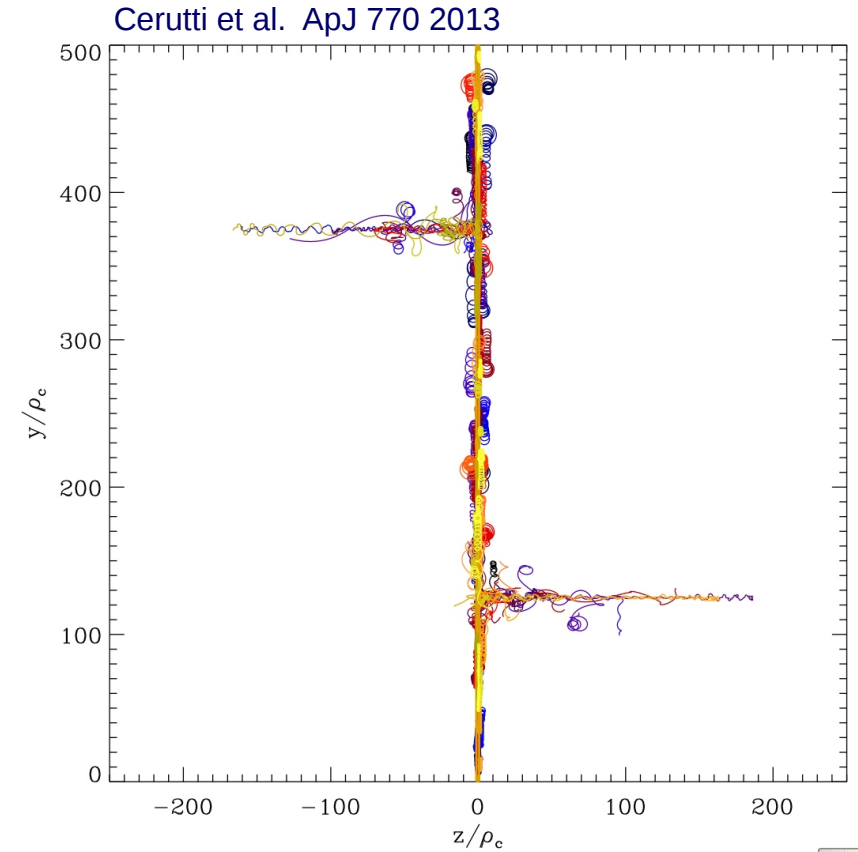
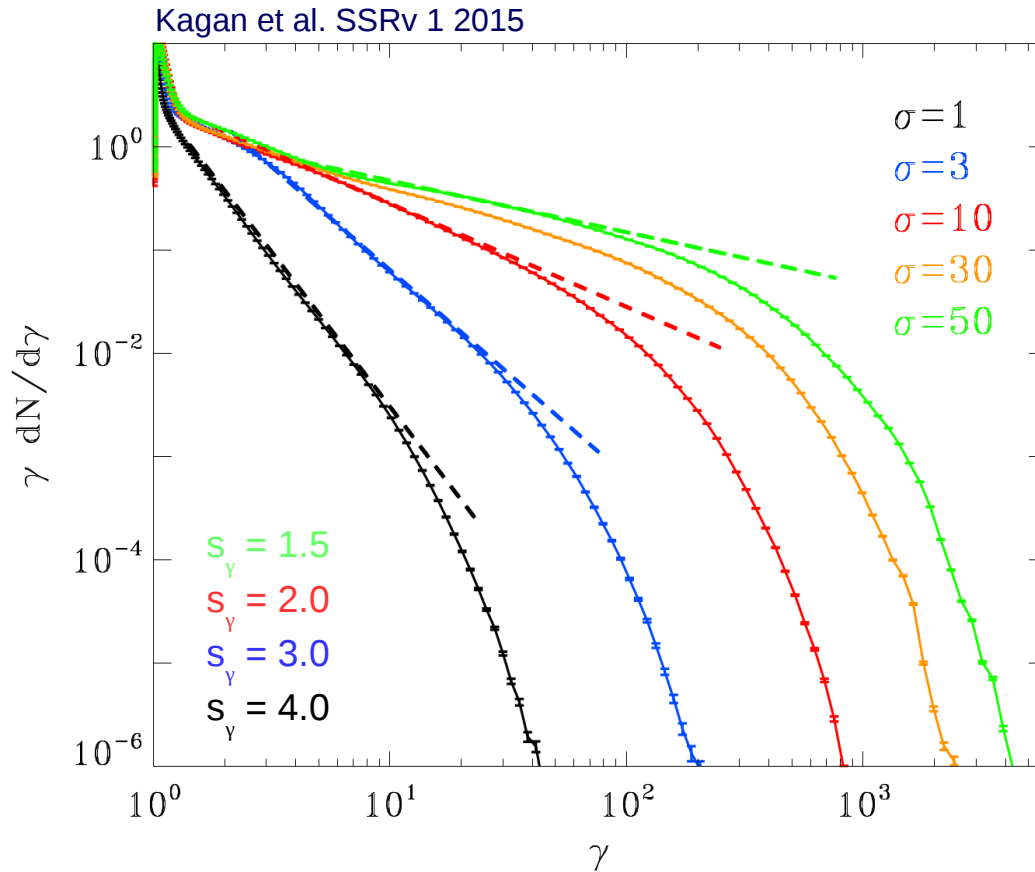
Particles are accelerated predominantly at X-points and concentrated into magnetic islands

Magnetic reconnection



Particles are accelerated predominantly at X-points and concentrated into magnetic islands

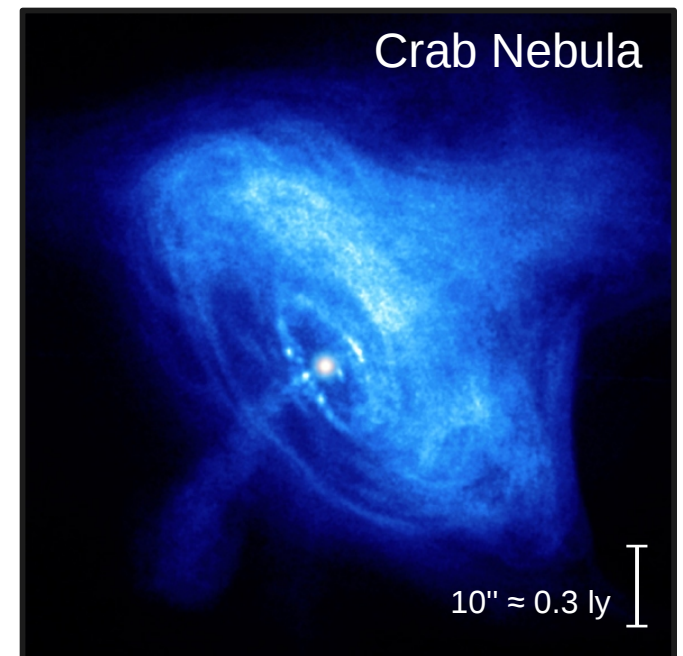
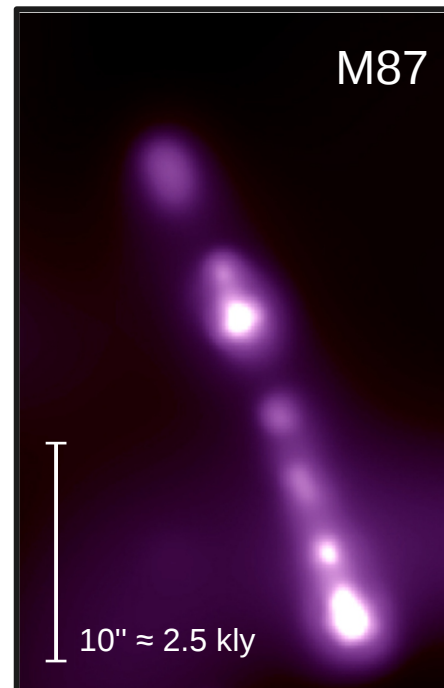
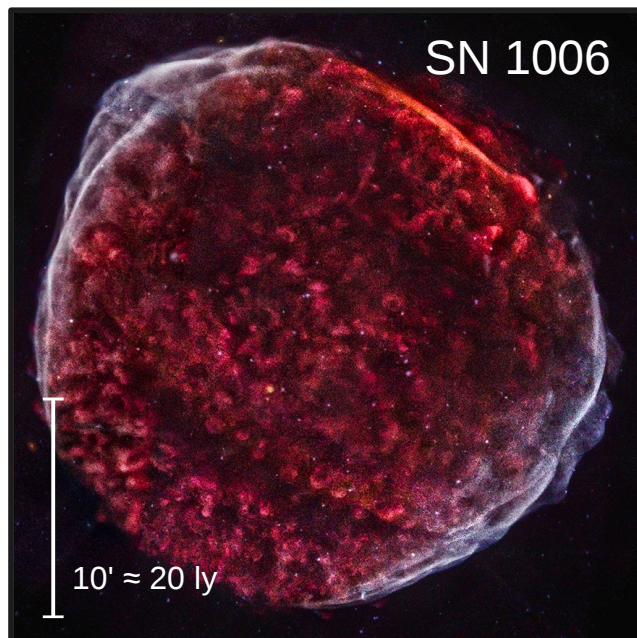
Magnetic reconnection



Large magnetization ($\sigma > 1$) are required. Hard spectra ($s_\gamma < 2$) are possible. Highest energy particles are beamed into z -direction

Accelerator property estimates

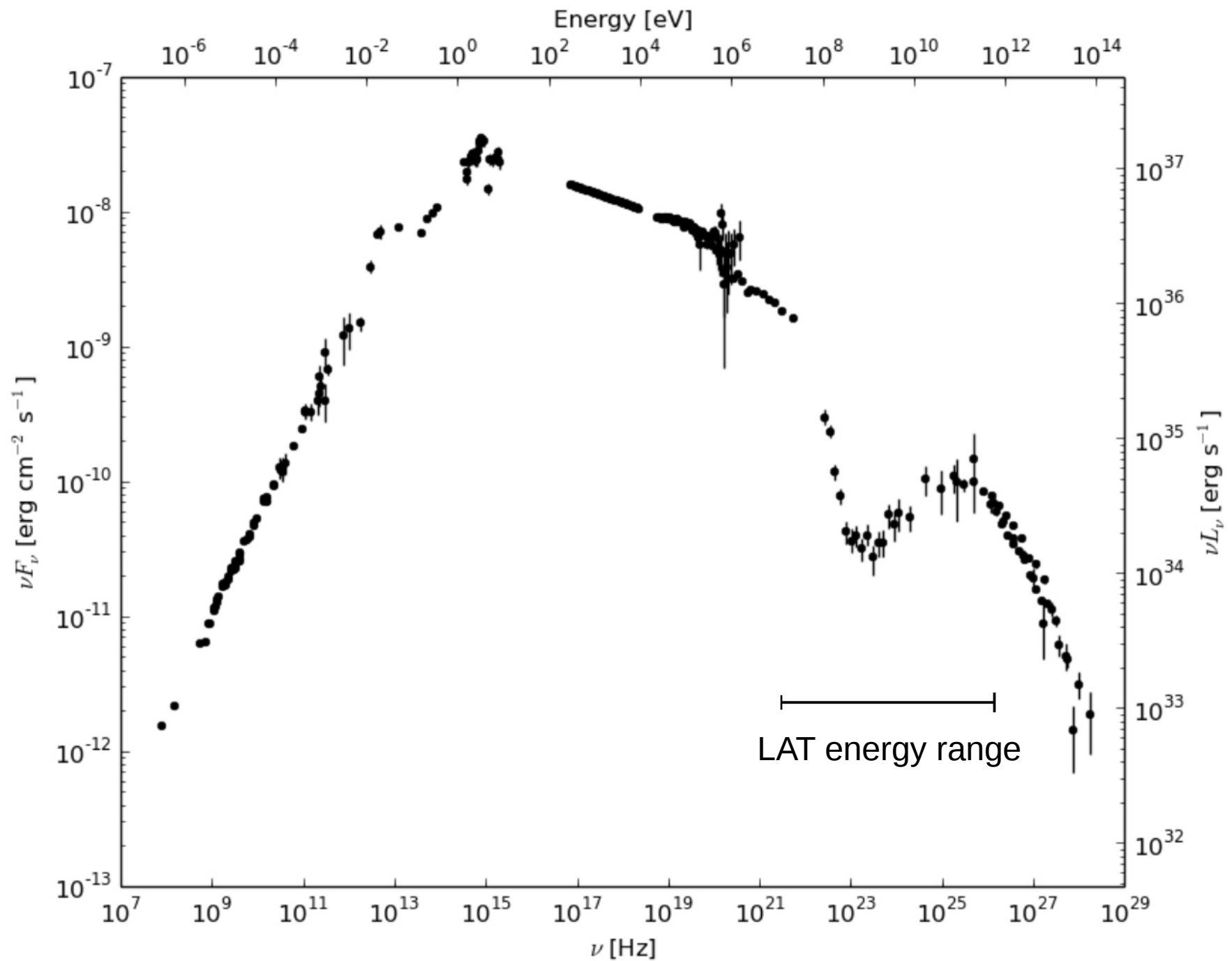
	Supernova Remnant shock	Active Galactic Nuclei jet	Pulsar Wind reverse shock
γ -1	$10^{-6} - 10^{-3}$	~ 10	$10^4 - 10^6$
σ	$< 10^{-3}$	$> 10^{-3}$	$0.01 - > 1$
Reconnection	no	(yes)	yes
Shock	yes	(no)	no





Radio VLA
Optical HST
X-rays Chandra

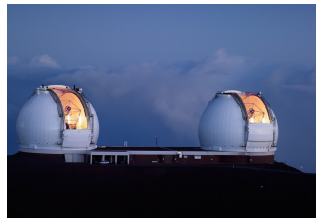
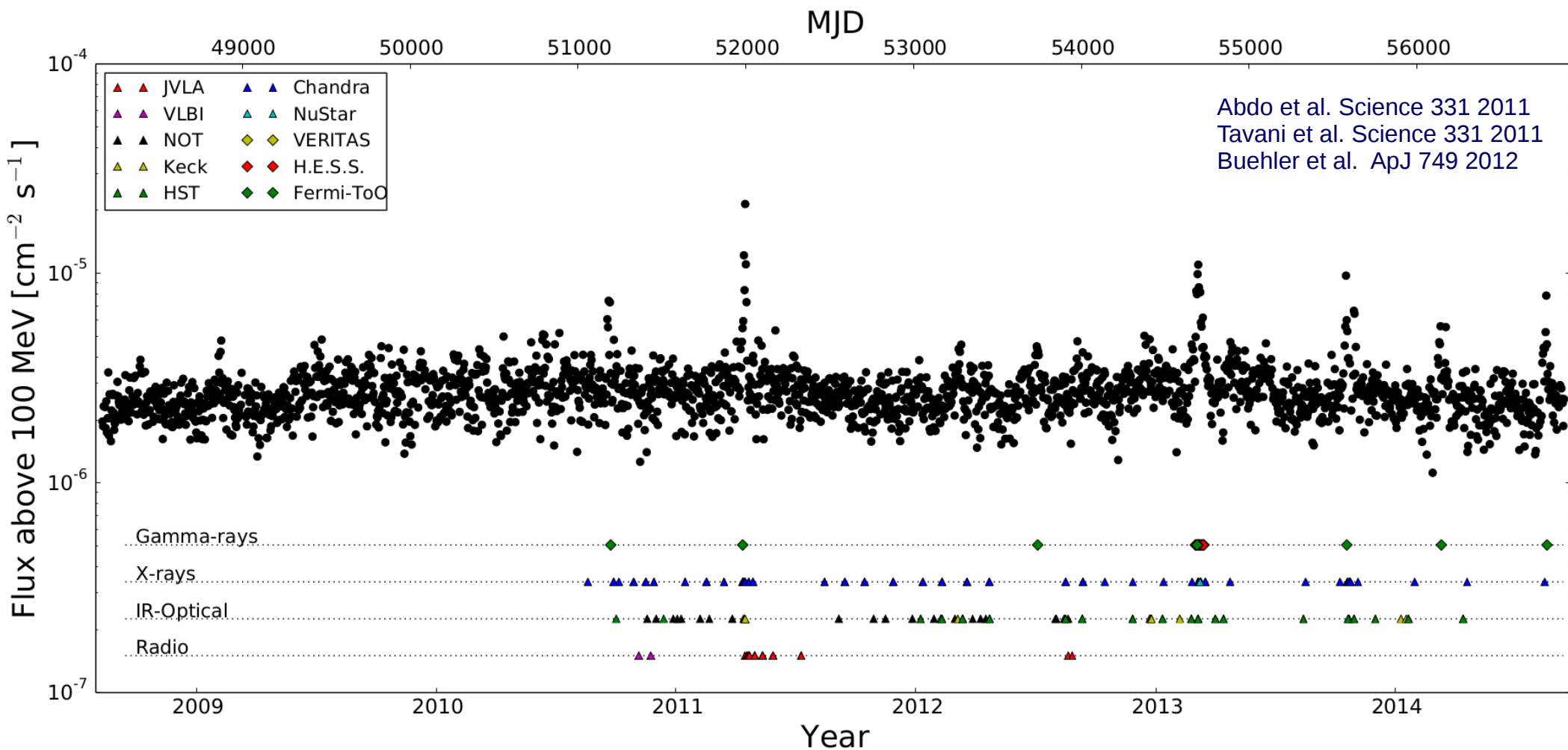
10" $\approx$ 0.3 ly

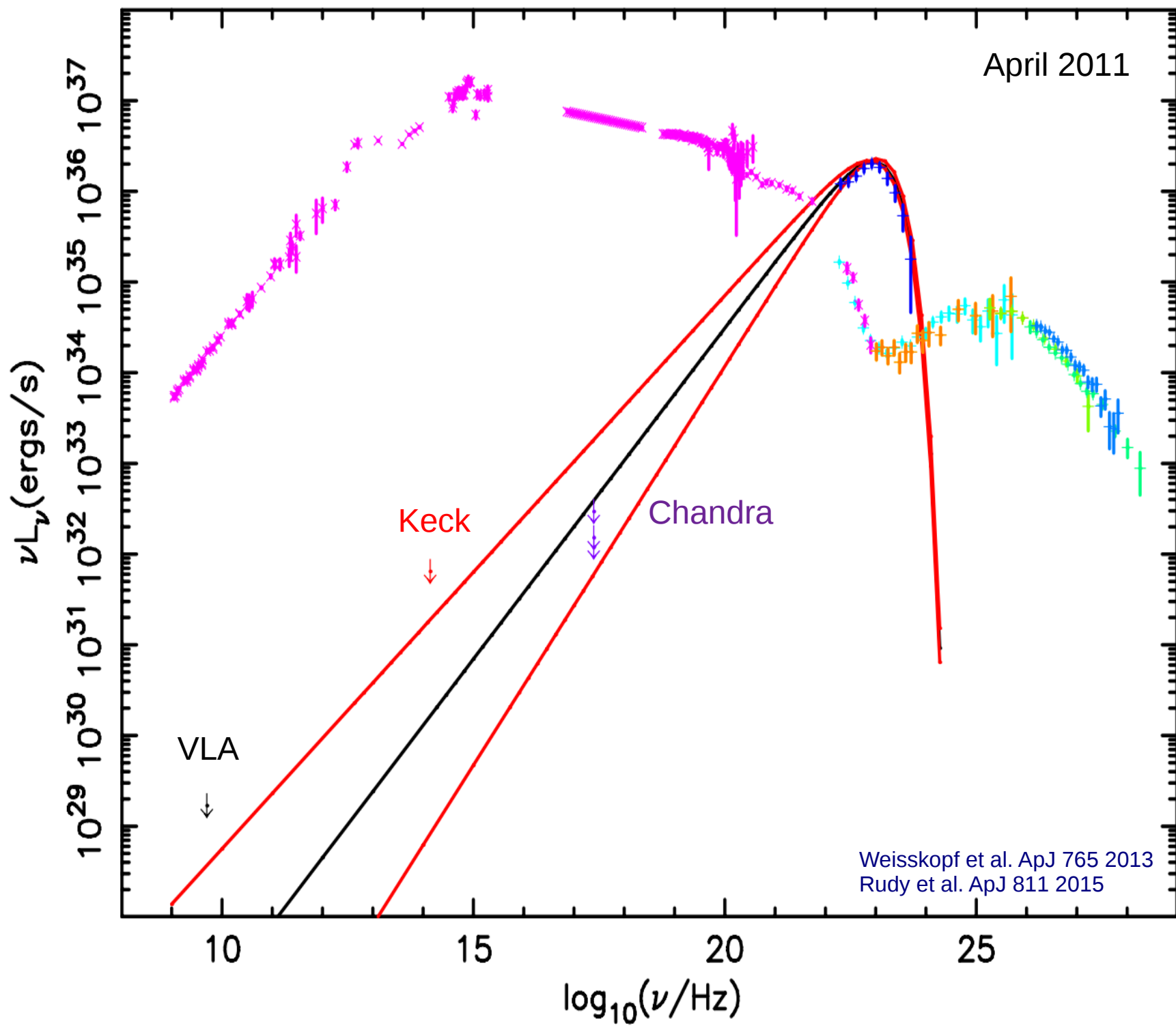


Radio
Optical
X-rays Chandra

$10 \approx 0.3 \text{ ly}$

The gamma-ray flares





Large energy release
of $\sim 1\%$ of the pulsar
spin-down power
suggests beaming

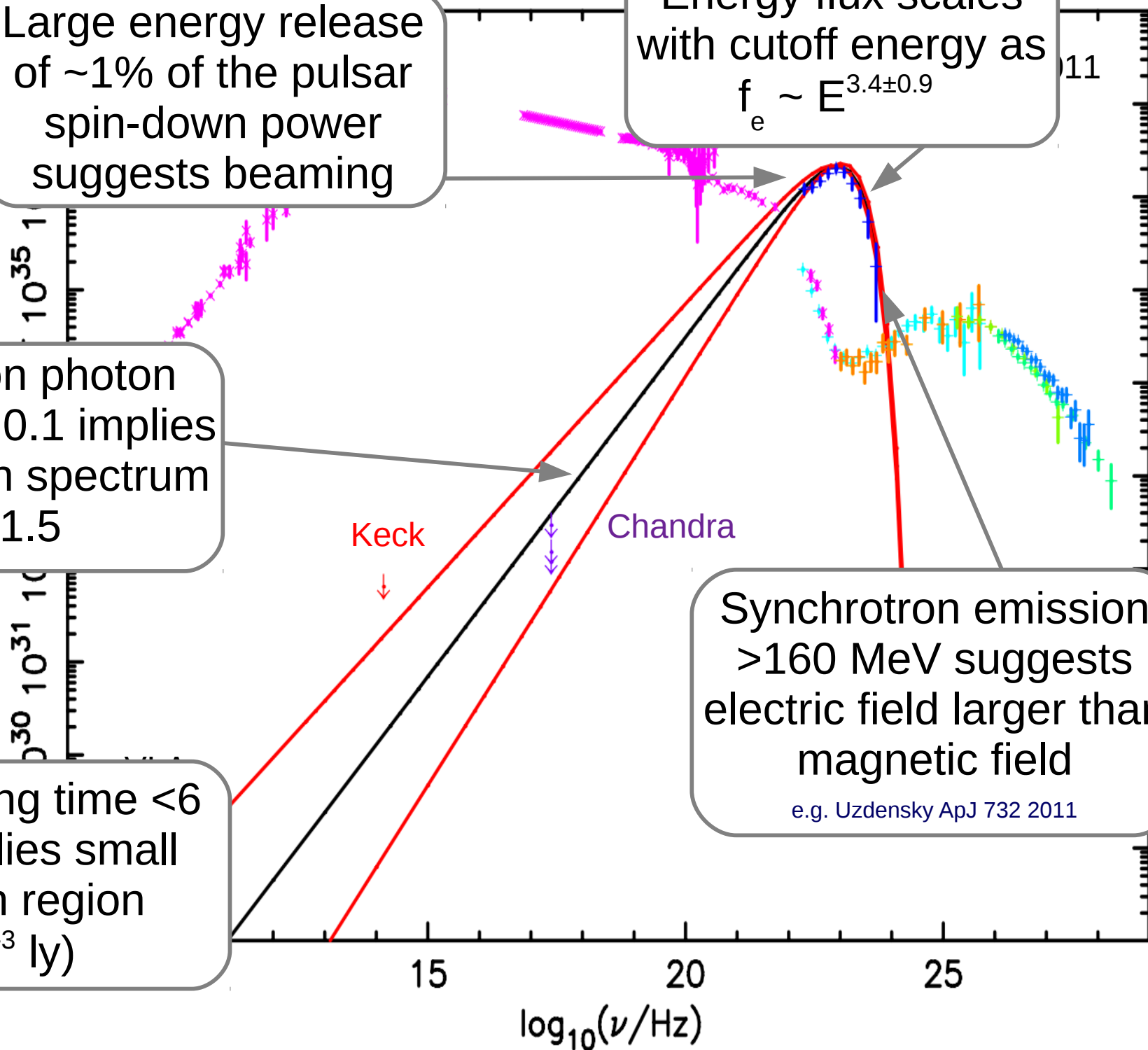
Energy flux scales
with cutoff energy as
 $f_e \sim E^{3.4 \pm 0.9}$

Synchrotron photon
index 1.24 ± 0.1 implies
hard electron spectrum
 $s \sim 1.5$

Flux doubling time < 6
hours implies small
emission region
($< 10^{-3}$ ly)

Synchrotron emission
 > 160 MeV suggests
electric field larger than
magnetic field

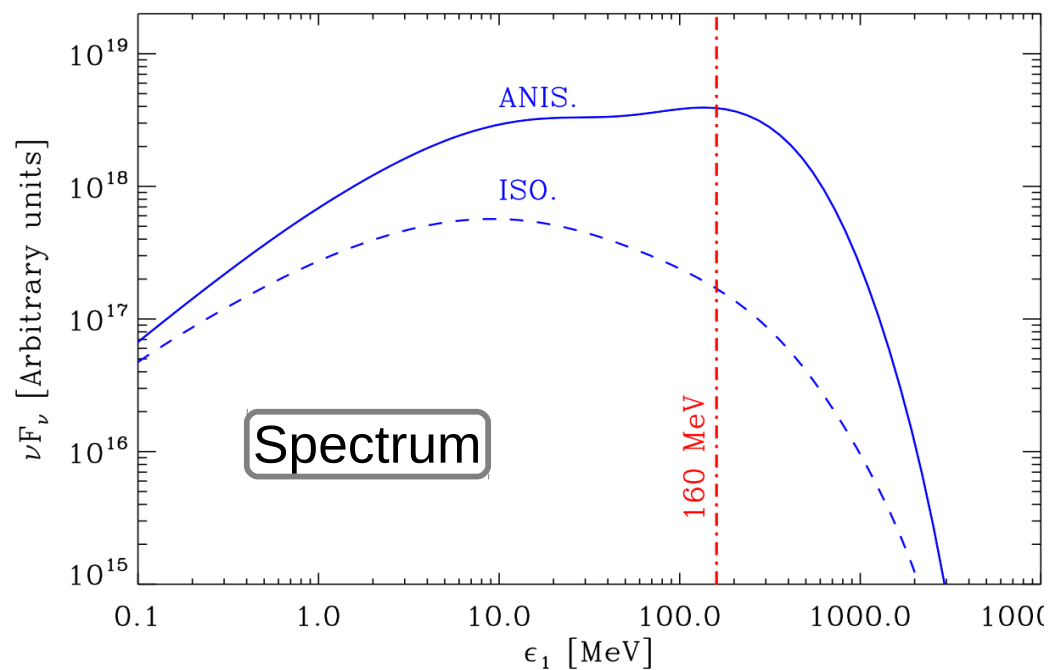
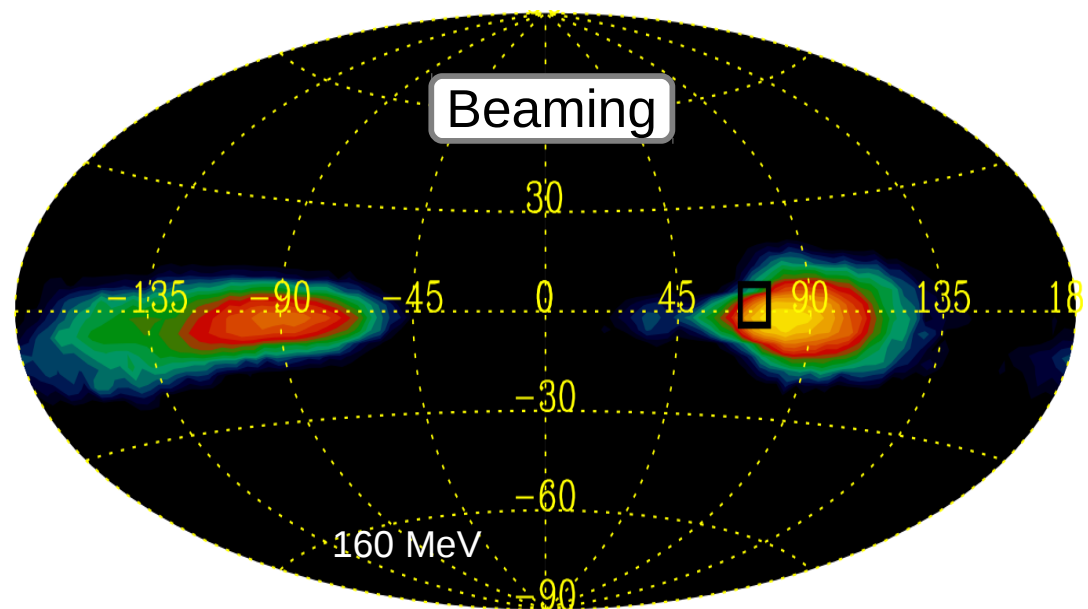
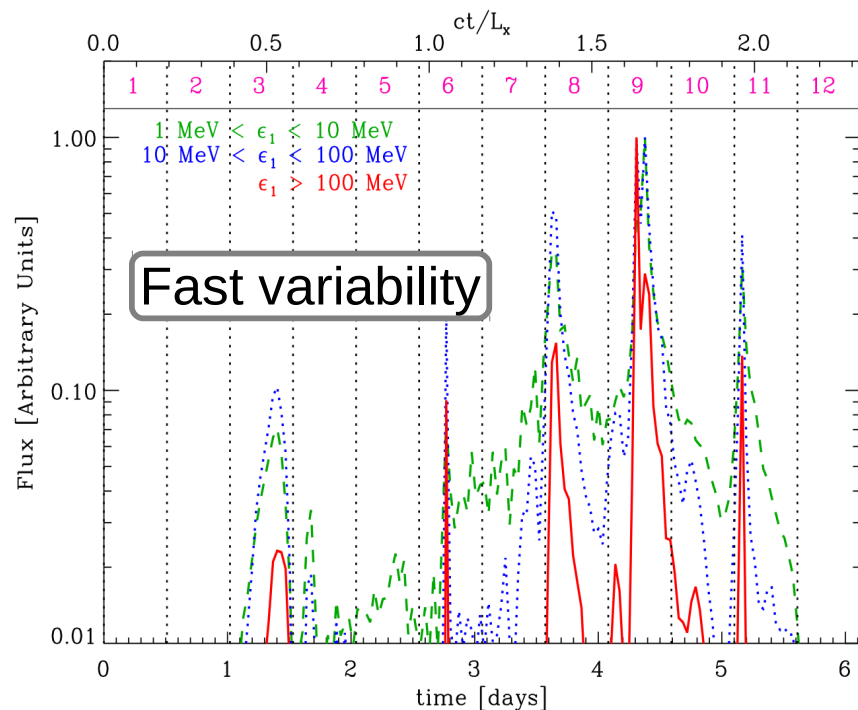
e.g. Uzdensky ApJ 732 2011



Flares from reconnection?

3D Particle In Cell simulations
can reproduce the flare
properties

Cerutti et al. ApJ 770 2013
Cerutti et al. ApJ 782 2014



Summary

At the contact point of plasma flows, shock acceleration and magnetic reconnection can accelerate non-thermal particles

Shock acceleration is inefficient at magnetized relativistic shocks. Magnetic reconnection appears more natural in relativistic flows (as AGN, PWN or GRBs)

The Crab nebula flare properties can be well explained as explosive reconnection events in Particle In Cell simulations

