Cosmological constraint on neutrino properties I: Cosmic Microwave Background



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My today's tasks: in order to convince people here that cosmological constraints on neutrino properties are important,

- 1) explain How? From what? Why?
- 2) review **how strongly and robustly** we can constrain them in current cosmological observations

1st half: Cosmic Microwave Background (CMB)

well understood theoretically and experimentally precise measurements by WMAP, PLANCK etc are available

#### 2nd half: Large-Scale Structure (LSS)

has started to enter the precision cosmology era like CMB necessary to combine LSS with CMB to reach detection of O(0.1eV)



#### 1. Simple Theoretical Background

- 2. Observational Results
  - WMAP9
  - WMAP9 + SPT/ACT
  - Planck
  - (Planck reanalysis)



3. Summary & Future

#### References:

- Textbook: 松原 隆彦, "現代宇宙論", S. Dodelson, "Modern Cosmology"
- Review: W. Hu & S. Dodelson (2002)
- WMAP: Komatsu et al. (2010, 2012), Hinshaw et al. (2013)
- SPT: Hou et al. (2013)
- Planck Collaboration (2013 I, XVI)

# **Cosmic History**



Past

z ~ 3000 Rad-Matter equality

http://lambda.gsfc.nasa.gov

- z ~ 1300 Recombination
- z ~ 1100 Photon decoupling
- z ~ 10 Reionization
- z ~ 1 Dark energy dominant

 $I_z = 0$  Present  $\Omega_X = \rho_X / \rho_{crit}$ 

Scale factor a(t)

Redshift 
$$1+z = \frac{a(t_0)}{a(t)}$$

Hubble parameter  $H = \frac{\dot{a}}{a}$ 

flat ACDM model: 7 parametersdark energy $\Omega_{\Lambda}$ baryon $\Omega_{b}$ cold dark matter (CDM) $\Omega_{c}$ Hubble constant  $H_0 = 100h \, [\mathrm{km/s/Mpc}]$ initial perturbation $\Delta^2_{\mathcal{R}}(k_0), n_{\rm s}(k_0)$ reionization optical depth $\tau_{\rm re}$ 

### Cosmology constrains neutrinos via 'GRAVITY'

#### Einstein Eq.

$$H^{2} = H_{0}^{2} \left\{ \Omega_{\Lambda} + \Omega_{m} (1+z)^{3} + \Omega_{r} (1+z)^{4} \right\}$$

Angular diameter distance  $D_A(z) = \frac{1}{1+z} \int_0^z \frac{cdz}{H}$ 

Massless  $\rightarrow$  Massive neutrinos

- become non-relativistic at late times.

$$\Omega_{\nu} = 0.0217 \left( \frac{\sum m_{\nu}}{1 \,\mathrm{eV}} \right) \left( \frac{0.7}{h} \right)^2$$

-  $m_{
u,i} = \langle E_{
u,i} \rangle$  yields to

$$1 + z_{\rm nr,i} = 1890 \left(\frac{m_{\nu,i}}{1 \,\mathrm{eV}}\right)$$

- effective number of neutrinos (c.f. standard  $N_{
m eff}\sim 3.046$  )

$$\rho_r = \left[1 + \frac{7}{8} \left(\frac{4}{11}\right)^{4/3} N_{\text{eff}}\right] \rho_\gamma$$

§1. Simple Theoretical Background

### **Cosmic Microwave Background**

At early times, photon interacts with baryon via Thomson scattering



Sound waves traveling in the baryon-photon fluid

$$r_s = \int_0^{t_{LS}} dt \frac{c(1+z)}{\sqrt{3/(1+3\rho_{\rm b}/(4\rho_{\gamma}))}}$$

used as a 'standard ruler' in the BAO analysis (2nd talk)



Right after recombination ( $z\sim1300$ ), photon becomes decoupled.  $\rightarrow$  observed as CMB photon of 2.726K

# Statistical Tool to analyze CMB fluctuation





For the modes beyond the horizon at decoupling (|<100),

$$\frac{\Delta T}{T}\Big|_{\rm SW} = \frac{1}{3}\Phi + 2\int dt \,\dot{\Phi}$$
  
SW Integrated SW

Sachs-Wolfe effect: since  $\Delta_{\ell}(k, t_*) = j_{\ell}(k\chi_*)$ ,  $\ell(\ell+1)C_{\ell} \propto \text{Const}$ 

Integrated Sachs-Wolfe effect:  $\Phi = \text{Const}$  if matter-dominant 1) early time ISW: not fully matter dominant at recombination important to understand *neutrino mass effect!* 

2) late time ISW: dark-energy dominant era



#### Hu & Dodelson (2002)





## (2) Acoustic Oscillation

For the modes within the horizon at decoupling (I>100),

$$\ddot{\delta_{\gamma}} + \frac{c_s^2}{a^2}k^2\delta_{\gamma} = -\frac{4k^2}{3a^2}\Phi$$

This is just a harmonic oscillator and its solution is

$$\frac{1}{4}\delta_{\gamma} + \Phi = \left[\frac{1}{4}\delta_{\gamma 0} + \left(1 + \frac{3\rho_{\rm b}}{4\rho_{\gamma}}\right)\right]\cos(kr_s) - \frac{3\rho_{\rm b}}{4\rho_{\gamma}}\Phi$$

odd-peak only depends on baryon density!

Sound waves traveling in the baryon-photon fluid

$$r_s = \int_0^{t_{LS}} dt \frac{c(1+z)}{\sqrt{3/(1+3\rho_{\rm b}/(4\rho_{\gamma}))}}$$

One of the best measured quantities in cosmology:

$$\frac{r_s(z_*)}{D_A(z_*)} = 1.19355 \pm 0.00078 \deg$$
Planck Collaboration 2013

§1. Simple Theoretical Background

1000

(c) Baryons

 $\Omega_{\rm b}h^2$ 

10

100

100

80

00 (NU) TA

# (3) Silk damping

Recombination is not instantaneous, and hence diffused photon can be affected by friction due to the velocity difference with photon

At smaller scales than photon's diffusion scale, the temperature fluctuation is washed out



This scale as  $1/\sqrt{Hn_{\rm e}}$ , while the sound horizon scales as 1/Hbecause  $r_d \sim \sqrt{3ct/\lambda_C}\lambda_C$  where photon's mean free path  $\lambda_C = 1/(\sigma_{\rm T}n_{\rm e})$ Note  $n_{\rm e} = (1 - Y_{\rm He})n_{\rm b}$ 

## **Reionization & CMB polarization**

The universe is known to be reioniozed at z ~ 10, and the CMB photon become scattered and obscured.

The temperature power spectrum is reduced by a factor of  $\exp[-2\tau_{\rm re}]$ 

$$\tau_{\rm re} = \int n_{\rm e} \sigma_T dt$$

degenerate with fluctuation parameters,  $\Delta_{\mathcal{R}}^2(k_0), n_s(k_0)$ 

Polarization can help!





Primary CMB is lensed by the foreground large-scale structure (LSS)



Hu & Okamoto (2002)

Smooth out the acoustic peak (~deg scale) Lewis & Challinor (2006)



Massive neutrinos can suppress LSS (2nd talk)  $\rightarrow$  less lens effect

Note the 4pt estimator is complementary Schmittfull et al. (2013)

CMB Constraint on >  $m_{\nu}$ 

increasing m means increasing matter at decoupling

Relativistic neutrino's energy at recombination  $E_i \sim 0.58 \,\mathrm{eV}$ In order to be non-relativistic at that time,

$$\sum m_{\nu} > 0.58 \times 3 = 1.74 \,\mathrm{eV}$$

One could reach beyond the limit by a little bit. e.g. ~1.5eV

Ichikawa, Fukugida, Kawasaki (2005)



# CMB Constraint on $N_{\rm eff}$

- increasing N<sub>eff</sub> means increasing radiation at decoupling
  - nicely illustrated by WMAP9 paper



## Successful observations of CMB

- **COBE** nicely measured the Planck distribution  $2.725 \pm 0.001 K$ 
  - full-sky satellite: WMAP, PLANCK
  - terrestrial small-sky but high-resolution: **SPT, ACT, POLARBEAR** c.f. 西野さんのトーク



WMAP9 + SPT

• LCDM + 
$$\sum m_{\nu}$$
  
WMAP9 only  $\sum m_{\nu} < 1.3 \text{ eV} (95\% \text{ C.L.})$   
WMAP9 + ACT/SPT  $\sum m_{\nu} < 1.5 \text{ eV} (95\% \text{ C.L.})$ 

 $LCDM + N_{eff} + Y_{He}$ 

WMAP9 only

 $N_{\rm eff} > 1.7 \,(95\% \,{\rm C.L.})$ 

WMAP9 + ACT/SPT  $N_{\rm eff} = 2.92 \pm 0.79$ 





- LCDM + 
$$\sum m_{\nu}$$

Planck + WP + ACT/SPT  $\sum m_{\nu} < 0.66 \,\text{eV} \,(95\% \,\text{C.L.})$ 

+ CMB lensing marginalized  $\sum m_{\nu} < 1.08 \, {\rm eV} \, (95\% \, {\rm C.L.})$ 

We do see a discrepancy with CMB lensing (4pt)

 $N_{\rm eff} = 3.33^{+0.59}_{-0.83}$  (68%; *Planck*+WP+highL),  $Y_{\rm P} = 0.254^{+0.041}_{-0.033}$  (68%; *Planck*+WP+highL).



Even in ACDM model, there is a discrepancy b/w WMAP & PLANCK

## Systematics in PLANCK?



# Reanalyzed by D. Spergel



PLANCK

- conservative 24% sky
- marginalized over FG template

Their work

- less conservative 47% sky
- subtracted FG using 545GHz







D. Spergel et al. (2013)

 $N_{\rm eff} = 3.34 \pm 0.35$ 

# CMB: Summary & Future

CMB plays a very important role in cosmology

- precisely measure the basic parameters in  $\Lambda\text{CDM}$  universe

Neutrino Mass & N<sub>eff</sub> can be determined separately

-  $\sum m_{\nu}$  > 2.0eV is ruled out by CMB only

- CMB lensing is potentially powerful but has a discrepancy

-  $N_{\text{eff}}$  is close to the standard value 3.046

- need other probes!  $\rightarrow$  2nd talk!

#### Future

- Planck's next data release (w/ polarization) will be summer in 2014
- Many terrestrial experiments coming! c.f. 西野さんのトーク

Cosmological constraint on neutrino properties II: Large-Scale Structure



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### Contents

- 1. How to surpass the CMB constraints?
  - add low-redshift distance information
  - add large-scale structure information
- 2. Brief review of LSS observables
   Pros & Cons
- 3. Neutrino mass from **Galaxy Clustering** 
  - Baryon Acoustic Oscillations
  - Shape of galaxy power spectrum
  - Redshift-Space Distortion
- 4. Summary & Conclusion



## Neutrino Suppression



(CDM)

**b** 





# Neutrino Suppression in Nonlinear theory



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#### Abazajian et al. (2011, 2013)

Probe	Current $\sum m_{\nu}$ (eV)	Forecast $\sum m_{\nu}$ (eV)	Key Systematics	Current Surveys	Future Surveys
CMB Primordial	1.3	0.6	Recombination	WMAP, Planck	None
CMB Primordial + Distance	0.58	0.35	Distance measure- ments	WMAP, Planck	None
Lensing of CMB	$\infty$	0.2 - 0.05	NG of Secondary anisotropies	Planck, ACT [39], SPT [96]	EBEX [57], ACTPol, SPTPol, POLAR- BEAR [5], CMBPol [6]
Galaxy Distribution	0.6	0.1	Nonlinearities, Bias	SDSS [58, 59], BOSS [82]	DES [84], BigBOSS [81], DESpec [85], LSST [92], Subaru PFS [97], HET- DEX [35]
Lensing of Galaxies	0.6	0.07	Baryons, NL, Photo- metric redshifts	CFHT-LS [23], COS- MOS [50]	DES [84], Hy- per SuprimeCam, LSST [92], Euclid [88], WFIRST[100]
Lyman $\alpha$	0.2	0.1	Bias, Metals, QSO continuum	SDSS, BOSS, Keck	BigBOSS[81], TMT[99], GMT[89]
21 cm	$\infty$	0.1 - 0.006	Foregrounds, Astro- physical modeling	GBT [11], LOFAR [91], PAPER [53], GMRT [86]	MWA [93], SKA [95], FFTT [49]
Galaxy Clusters	0.3	0.1	Mass Function, Mass Calibration	SDSS, SPT, ACT, XMM [101] Chan- dra [83]	DES, eRosita [87], LSST

Important to trace characteristic time and scale dependence



Key issue in each observable is how to control systematics

# PLANCK CMB Lensing

Planck XVII (2013)







### Why galaxy clustering?

 $\begin{array}{l} \text{Galaxy clustering} \quad \delta_{\mathrm{g}}(\mathbf{x}) = \frac{n_{\mathrm{g}}(\mathbf{x}) - \bar{n}_{\mathrm{g}}}{\bar{n}_{\mathrm{g}}} \\ \text{2pt statistics} \\ \text{Power Spectrum} \text{ in Fourier Space} \\ \langle \delta_{\mathrm{g}}(\mathbf{k}) \delta_{\mathrm{g}}(\mathbf{k}') \rangle = (2\pi)^{3} \delta_{D}(\mathbf{k} + \mathbf{k}') P_{\mathrm{g}}(\mathbf{k}) \\ \text{Correlation Function} \text{ in Configuration Space} \\ \langle (1 + \delta_{\mathrm{g}}(\mathbf{x}))(1 + \delta_{\mathrm{g}}(\mathbf{x} + \mathbf{r})) \rangle = 1 + \xi(\mathbf{r}) \end{array}$ 



Reasons why 3D large-scale galaxy clustering is interesting:

- 1. By eye, looks like  $\,\delta_{\rm g}\sim {\cal O}(1)$  (c.f. CMB:  $\delta T/T\sim {\cal O}(10^{-5})$  )
- 2. Traces underlying matter distribution  $P_{\rm g}(k) \approx b_1^2 P_{\rm m}(k)$ 
  - Baryon Acoustic Oscillations ~150Mpc → *dark energy*
  - mildly nonlinear regime → *initial condition, neutrino mass*
  - Redshift Space Distortions → *modified gravity*
- 3. Observation is not technically challenging

## Baryon Acoustic Oscillations (BAOs)

 $\text{og}_{10}~\text{P(k)} \ / \ h^{-3}\text{Mpc}^3$ 

Sunyaev & Zel'devich 1970, Peebles & Yu 1970...

Sound waves traveling in the baryon-photon fluid

$$r_s = \int_0^{t_{LS}} dt \frac{c(1+z)}{\sqrt{3/(1+3\rho_{\rm b}/(4\rho_{\gamma}))}}$$

The sound horizon at recombination can be seen as acoustic peaks in CMB

$$\frac{r_s(z_*)}{D_A(z_*)} = 1.19355 \pm 0.00078 \deg$$
Planck Collaboration 201

After baryon-dragging epoch, BAOs are imprinted into matter

 $r_d = 147.49 \pm 0.59 \, \mathrm{Mpc}$  Hu & Sugiyama 1996

But suppressed by  $\frac{\Omega_{\rm b}}{\Omega_{\rm m}} \sim \mathcal{O}(0.1)$ 



## BAO as a 'standard ruler'

When making a 3D map, d-z relation is necessary Alcock & Paczynski (1979)









true cosmology: *isotropic* 

fiducial (wrong) cosmology: anisotopic

In the BAO business: very robust against systematics **1D BAO**: spherically-averaged  $D_{\rm V} = \left[cz(1+z)^2 \frac{D_{\rm A}^2}{H}\right]^{1/3}$ 

**2D BAO (AP test)**:  $(D_A \times H)$  Padmanabhan & White 2006

\* shape of P(k) can be also useful for the AP signal

#### **Redshift Space Distortion (RSD)**



 A distance to a galaxy is measured by "redshift" which cannot be distinguished from its peculiar velocity

redshift space 
$$\vec{s} = \vec{r} + \frac{\vec{v} \cdot \hat{z}}{aH(z)}$$
 line of sight direction



### RSD in linear theory

In linear theory, the anisotropic P(k) in redshift space is

$$P_{\rm g}(k,\mu) = b^2 \left(1 + \frac{f}{b}\mu^2\right) {}^2 P^L(k)$$
 Kaiser (1987)

Legendre multipole expansion is useful to characterize P(k, $\mu$ )  $\vec{v} \cdot \hat{z}$ 

$$P(k,\mu) = \underbrace{\sum_{\ell}}_{\ell} \overline{P_{\ell}}(k) \underbrace{\frac{v \cdot z}{\widehat{aH}(z)}}_{aH(z)} \hat{z}$$

Notice that, in linear theory, we only have **monopole (I=0), quadrupole (I=2), hexadecapole (I=4)** 

• The definition of 'f=dlnD/dlna' comes from the velocity field 
$$f\sigma_8$$
  
linear Euler equation  $\frac{\partial \mathbf{u}(\mathbf{x},\tau)}{\partial \tau} + \mathcal{H}(\tau) \mathbf{u}(\mathbf{x},\tau) = -\nabla \Phi(\mathbf{x},\tau)$   
gravity test  
 $f = \Omega_{\mathrm{m}}(z)^{\gamma} \quad \gamma = 0.545 \text{ (GR)} \quad \gamma = 0.68 \text{ (DGP)}$ 

Note that RSDs are complementary to weak lensing  $(\Phi+\Psi)$ 

### Galaxy Redshift Surveys

Many gigantic galaxy redshift surveys for **BAO** and **RSD** 

Completed: CfA2, 2dF, SDSS-II, WiggleZ

Ongoing: **BOSS**, VIPERS, FastSound

Future: HETDEX, SuMiRe PFS, DESI, Euclid

Complementary to imaging (photo-z) surveys

surveys: DES, HSC, LSST etc...

designed for weak lensing

angular clustering is possible, but little information on RSD



### **Baryon Oscillation Spectroscopic Survey**

- A part of Sloan Digital Sky Survey III (2009-2014) Eisenstein et al. 2011
  - 2.5m telescope in Apache Point Observatory in NM, USA
  - image (u,g,r,i,z) bands with r~22.5 covering ~14,000deg<sup>2</sup> Fukugida et al. 1996
  - BOSS (DR9-DR12)
    - 1) 1.5 million luminous galaxies
    - 2) 150,000 quasars (Ly-a forest) e.g. Slosar et al. 2012



## Huge improvement of 'CMASS' from DR9 to DR11



### Nonlinear Issues in modeling P(k)

Accurate prediction of the nonlinear galaxy P(k) is necessary:



Taruya, Nishimichi, S.S. (2010), Nishimichi & Taruya (2012), Taruya et al. (2013) Matsubara (2008a, 2008b, 2011, 2013) Seljak & McDonald (2012), Okumura et al. (2012a, 2012b), Vlah et al. (2012,2013) Reid & White (2010), Carlson et al. (2012), Wang et al. (2013)

galaxy bias  $\delta_{\rm g}/\delta_{\rm h} \leftrightarrow \delta_{\rm m}$   $\delta_{\rm g}(k) = b_0 \frac{1 + A_1 k^2}{1 + A_0 k} \delta_{\rm m}(k)$ 

e.g. Nishimichi & Taruya (2011), Oka, S.S. et al. (2013)

## DR9 analysis to use 'shape'

#### G. Zhao, SS et al. (2013)

Use just monopole. Model is based on Perturbation Theory.



# Want to do for DR11 but...

#### SS et al., in prep

![](_page_40_Picture_2.jpeg)

We could do but I personally don't want to redo the same thing...

- a lot of rooms to improve our modeling
- now can be compared in detail against simulations with neutrinos
- need to check if our model can return the input value of  $\Sigma m v$

![](_page_40_Picture_7.jpeg)

#### 1% measurement of cosmological distance!

![](_page_41_Figure_1.jpeg)

#### In a flat ACDM universe

![](_page_42_Figure_1.jpeg)

### The first measurement of the BOSS P<sub>d</sub>(k)

![](_page_43_Figure_1.jpeg)

- These are the anisotropic P(k) firstly measured in the BOSS sample.
- Impressive precision level! ~1% for monopole & ~10% for quadrupole.
  - The difference b/w **north** and **south** is due to the survey window.

### Performance of our model

![](_page_44_Figure_1.jpeg)

1.9%

1.6%

9.1%

8.3%

4.0%

statistical error

3.1%

Results

fitting range	0.01 -	$0.15h/{ m Mpc}$	$0.01$ - $0.20h/{ m Mpc}$		
	best fit	mean $\pm 1\sigma$	best fit	mean $\pm 1\sigma$	
$f_{  }$	1.008	$1.005\pm0.057$	1.014	$1.018\pm0.036$	
$f_{\perp}^{"}$	1.026	$1.029\pm0.023$	1.029	$1.029\pm0.015$	
$f(z_{ m eff})\sigma_8(z_{ m eff})$	0.420	$0.423\pm0.052$	0.422	$0.419 \pm (0.042 + 0.014)$	.N!!
$b_1\sigma_8(z_{ m eff})$	1.221	$1.222\pm0.044$	1.221	$1.224 \pm 0.031$	
$b_2\sigma_8(z_{ m eff})$	1.7	$0.7\pm1.2$	-0.21	$-0.09 \pm 0.62$	
$\sigma_v$	$4.6{ m Mpc}/h$	$4.3\pm1.3{\rm Mpc}/h$	$4.63{ m Mpc}/h$	$4.65\pm0.81{\rm Mpc}/h$	
Ν	$1030  [{ m Mpc}/h]^3$	$1080 \pm 620  [\mathrm{Mpc}/h]^3$	$1890  [{ m Mpc}/h]^3$	$1690 \pm 600  [\mathrm{Mpc}/h]^3$	
$D_V(z_{ m eff})/r_s(z_d)$	13.83	$13.85\pm0.27$	13.88	$13.89\pm0.18$	
$F_{ m AP}(z_{ m eff})$	0.684	$0.686\pm0.046$	0.683	$0.679 \pm 0.031$	
$H(z_{ m eff})r_s(z_d)/r_s^{ m fid}(z_d)$	$94.0\mathrm{km/s/Mpc}$	$94.1\pm5.4\mathrm{km/s/Mpc}$	$93.5\mathrm{km/s/Mpc}$	$93.1\pm3.3\mathrm{km/s/Mpc}$	
$D_A(z_{ m eff})r_s^{ m fid}(z_d)/r_s(z_d)$	$1385\mathrm{Mpc}$	$1389\pm31{\rm Mpc}$	$1389{ m Mpc}$	$1388\pm22{\rm Mpc}$	
$\beta$	0.344	$0.346\pm0.043$	0.346	$0.342 \pm 0.037$	
$b_1  imes (0.8/\sigma_8)$	2.035	$2.037\pm0.073$	2.035	$2.040 \pm 0.052$	
$b_2  imes (0.8/\sigma_8)$	2.8	$1.2\pm2.0$	-0.4	$-0.2 \pm 1.0$	

![](_page_45_Figure_2.jpeg)

![](_page_45_Figure_3.jpeg)

#### Consistency check in ACDM+GR

![](_page_46_Figure_1.jpeg)

#### Abandon GR?

$$\int f = \Omega_{
m m}(z)^{\gamma}$$
  $\gamma = 0.545$  (GR)

![](_page_47_Figure_2.jpeg)

Interestingly, RSD+AP leads to larger γ, i.e., weaker gravity than GR at ~2σ level, independently of the CMB prior choice

Remark that this is just a LCDM+GR consistency check! In general, model-dependent analysis would be necessary

e.g., Taruya et al. 2013

![](_page_48_Picture_0.jpeg)

The low redshift measurements are necessary to break CMB limit to constrain the neutrino mass

Galaxy survey provides us a unique opportunity to measure

- distance from BAO
- growth from P(k) shape & RSD

Given the discrepancy b/w WMAP & Planck, it seems early to conclude anything.