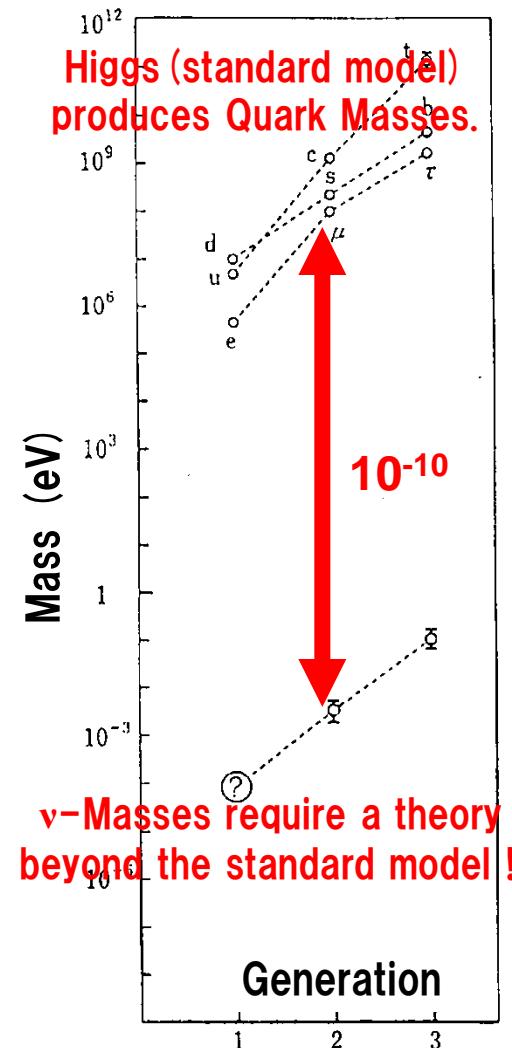


第27回宇宙ニュートリノ研究会
「宇宙と実験室から測定するニュートリノの絶対質量」
ICRR-UT, Jan. 20, 2014

CMB, BBN, 超新星爆発による
ニュートリノ質量および
振動への制限

梶野 敏貴
国立天文台 理論研究部
東大大学院 理学系研究科 天文学専攻

Challenge of the Century



Universe is flat and expanded acceleratingly.

$$\Omega_B + \Omega_{CDM} + \Omega_\Lambda = 1$$

■ What is CDM ($\Omega_{CDM} = 0.27$) and DE ($\Omega_\Lambda = 0.68$) ?

CMB & LSS including absolute ν -mass

■ Is BARYON sector ($\Omega_B = 0.05$) well understood ?

BBN ^7Li -Problem with DMs (Axion, SUSY ...)

SUSY-DM \Rightarrow beyond the Standard Model $\Rightarrow m_\nu \neq 0$, unique signal

Key Physics with $m_\nu \neq 0$ beyond the Standard Model :

- Unification, CP & L- & B- genesis, Dirac or Majorana ?
- Dark Matter & Big Bang Nucleosynthesis ?
- Explosion Mechanism of CC-SNe & Nucleosynthesis ?

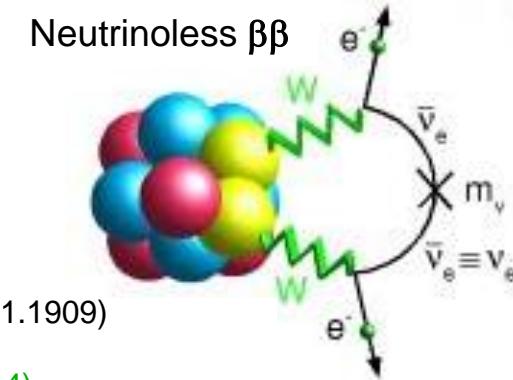
Purpose

is to constrain the total absolute ν -Mass, Hierarchy, and Nuclear EOS from ν -nucleus interactions in CMB, BBN, SNe and Relic SN ν 's.

Total ν -Mass constraint from Nuclear Physics and Cosmology

● $0\nu\beta\beta$ in COUORE, NEMO3, EXO, KamLAND Zen

$|\sum U_{e\beta}^2 m_\beta| < 0.3 \text{ eV}$: COUORE, NEMO3, EXO, KamLAND Zen (2012)



● CMB Anisotropies + LSS

$\sum m_\nu < 0.36 \text{ eV (95\%C.L.)}$: WMAP-7yr + HST + CMASS (Putter et al. arXiv:1201.1909)

$\sum m_\nu = 0.36 \pm 0.13 \text{ eV}$: WMAP-9yr + BOSS (Beauler, Saito & Seo: Jan. 20, 2014)

$\sum m_\nu < 0.2 \text{ eV (2σ, } B_\lambda < 2nG\text{)}$: Incl. Magnetic Field; Ymazaki, Kajino, Mathews & Ichiki, Phys. Rep. 517 (2012), 141; Phys. Rev. D81 (2010). 103519.

★ ν free-streaming effects

★ Integrated Sachs-Wolfe Effect !

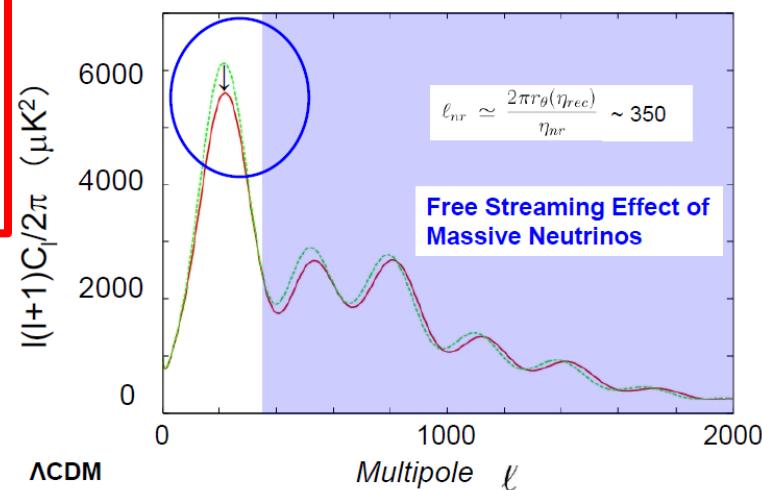
★ CMB anisotropies is generated even by:
compensation mode of ν -anisotropic stress (π_ν)
to primordial extra anisotropic stress (π_{ext}).

Magnetic field, dark radiation, Ekpyrotic universe, etc.

Effects of Neutrino Mass m_ν ($m_\nu = 2\text{eV}$)

Early Time Integrated Sachs-Wolfe
Effect, similar to CDM

S. Dodelson, E. Gates and A. Stebbins (1996)



Roles of ν -Anisotropic Stress with $m_\nu \neq 0$!

Standard inflationary cosmology needs
tuning primordial initial fluctuation !

Equations for CMB Calculations including Massive-v, and Primordial Magnetic Field

Ma & Bertschinger, 1995, ApJ, 455, 7

Linearized $\left\{ \begin{array}{l} \text{Energy-momentum conservation } T_{\mu,v}^v = 0 \\ \text{Boltzmann equation } P^\mu \frac{\partial f}{\partial x^\mu} - \Gamma_{\nu\lambda}^\mu P^{\nu\lambda} \frac{\partial f}{\partial P^\mu} = C[F] \end{array} \right.$

CDM $\dot{\delta}_c = -\frac{1}{2}\dot{h}$ (1)

$$\dot{\delta}_\nu = -\frac{4}{3}\theta_\nu - \frac{2}{3}\dot{h}$$
 (2)

Massless neutrino $\dot{\theta}_\nu = k^2 \left(\frac{1}{4}\delta_\nu - \sigma_\nu \right)$ (3)

neutrino $\dot{\theta}_\nu = 2\dot{\sigma}_\nu = \frac{8}{15}\theta_\nu - \frac{3}{5}kF_{\nu 3} + \frac{4}{15}\dot{h} + \frac{8}{5}\dot{\eta}$ (4)

$$\dot{F}_{\nu l} = \frac{k}{2l+1} (lF_{\nu(l-1)} - (l+1)F_{\nu(l+1)}), l \geq 3$$
 (5)

$$\dot{\Psi}_0 = \frac{qk}{\epsilon} \Psi_1 + \frac{1}{6}\dot{h} \frac{d \ln f_0}{d \ln q}$$
 (6)

$$\dot{\Psi}_1 = \frac{qk}{3\epsilon} (\Psi_0 - 2\Psi_2)$$
 (7)

Massive neutrino $\dot{\Psi}_2 = \frac{qk}{5\epsilon} (2\Psi_1 - 3\Psi_3) - \left(\frac{1}{15}\dot{h} + \frac{2}{5}\dot{\eta} \right) \frac{d \ln f_0}{d \ln q}$ (8)

$$\dot{\Psi}_l = \frac{qk}{(2l+1)\epsilon} (l\Psi_{l-1} - (l+1)\Psi_{l+1}), l \geq 3$$
 (9)

$$\dot{\delta}_\gamma = -\frac{4}{3}\theta_\gamma - \frac{2}{3}\dot{h}$$
 (10)

$$\dot{\theta}_\gamma = k^2 \left(\frac{1}{4}\delta_\gamma - \sigma_\gamma \right) + an_e \sigma_T (\theta_b - \theta_\gamma)$$
 (11)

photon $\dot{F}_{\gamma 2} = 2\dot{\sigma}_\gamma = \frac{8}{15}\theta_\gamma - \frac{3}{5}kF_{\gamma 3} + \frac{4}{15}\dot{h} + \frac{8}{5}\dot{\eta} - \frac{9}{5}an_e \sigma_T \sigma_\gamma + \frac{1}{10}an_e \sigma_T (G_{\gamma 0} + G_{\gamma 2})$ (12)

$$\dot{F}_{\gamma l} = \frac{k}{2l+1} (lF_{\gamma(l-1)} - (l+1)F_{\gamma(l+1)}) - an_e \sigma_T F_{\gamma l}, l \geq 3$$
 (13)

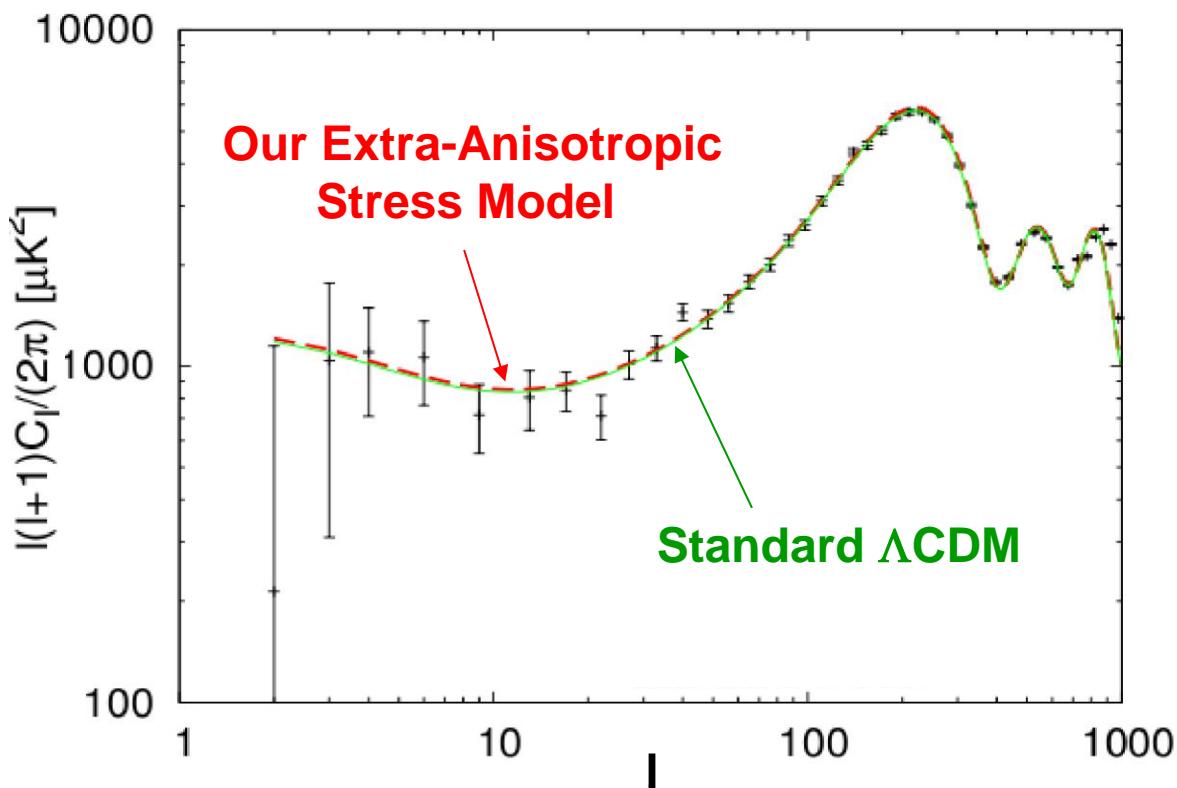
$\dot{G}_{\gamma l} = \frac{k}{2l+1} (lG_{\gamma(l-1)} - (l+1)G_{\gamma(l+1)}) + an_e \sigma_T \left[-G_{\gamma l} + \frac{1}{2}(F_{\gamma 2} + G_{\gamma 0} + G_{\gamma 2}) \left(\delta_{l0} + \frac{\delta_{l2}}{5} \right) \right], l \geq 3$ (14)

$$\dot{\delta}_b = -\theta_b - \frac{1}{2}\dot{h}$$
 (15)

baryon $\dot{\theta}_b = -\frac{\dot{a}}{a}\theta_b + c_s^2 k^2 \delta_b + \frac{4\bar{\rho}_\gamma}{3\bar{\rho}_b} an_e \sigma_T (\theta_\gamma - \theta_b)$ (16)

CMB from Neutrino & Extra Anisotropic Stress

K. Kojima, T. Kajino & G. J. Mathews,
JCAP **02** (2010), 0182009.



Spectral index & strength of primordial extra-anisotropic stress are set equal to be the CMB-best fit value.

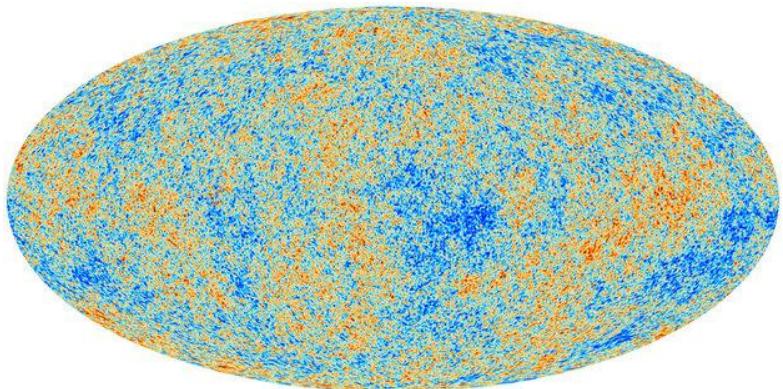
$$|\pi_{\text{ext}}| \sim 8.4 \times 10^{-6}$$

Our Extra-Primordial Anisotropic Stress Model is NOT an alternative to INFLATION !

ν -compensation mode (π_ν) plays a critical role in CMB with ν of finite mass !

- Curvature perturbation is generated by extra anisotropic stress π_{ext} and regulated later by ν –compensation mode π_ν after decoupling.
- It is desirable to know the cosmological origin of extra anisotropic stress π_{ext} and its generation epoch in the early universe.

CMB Temp. & Polar. including Cosmological Magnetic Field → Neutrino Mass Constraint

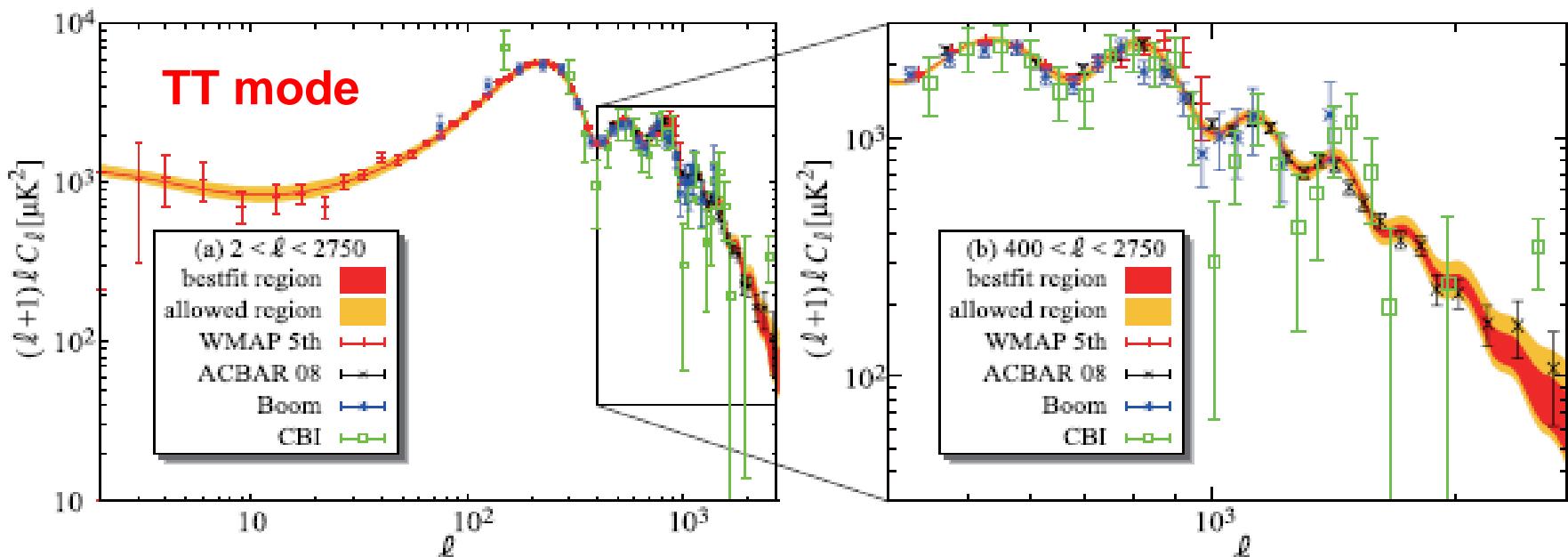


An example Primordial Anisotropic Stress

Lewis 2004; Mack 2002;
 Challinor 2004; Kahnashvili & B. Ratra 2005;
 Kosowsky et al. 2005;
 Ichikawa, Fukugita & Kawasaki 2005;
 Yamazaki et al. 2005 - 2012; Kojima et al. 2009 – 2010.

MCMC fit to CMB anisotropies
 PR D81 (2010), 023008; PR D81 (2010), 103519; Phys. Rep. 517 (2012), 141.

D. Yamazaki, K. Ichiki, T. Kajino, G. J. Mathews,
 PR D81 (2010), 023008; PR D81 (2010), 103519; Phys. Rep. 517 (2012), 141.

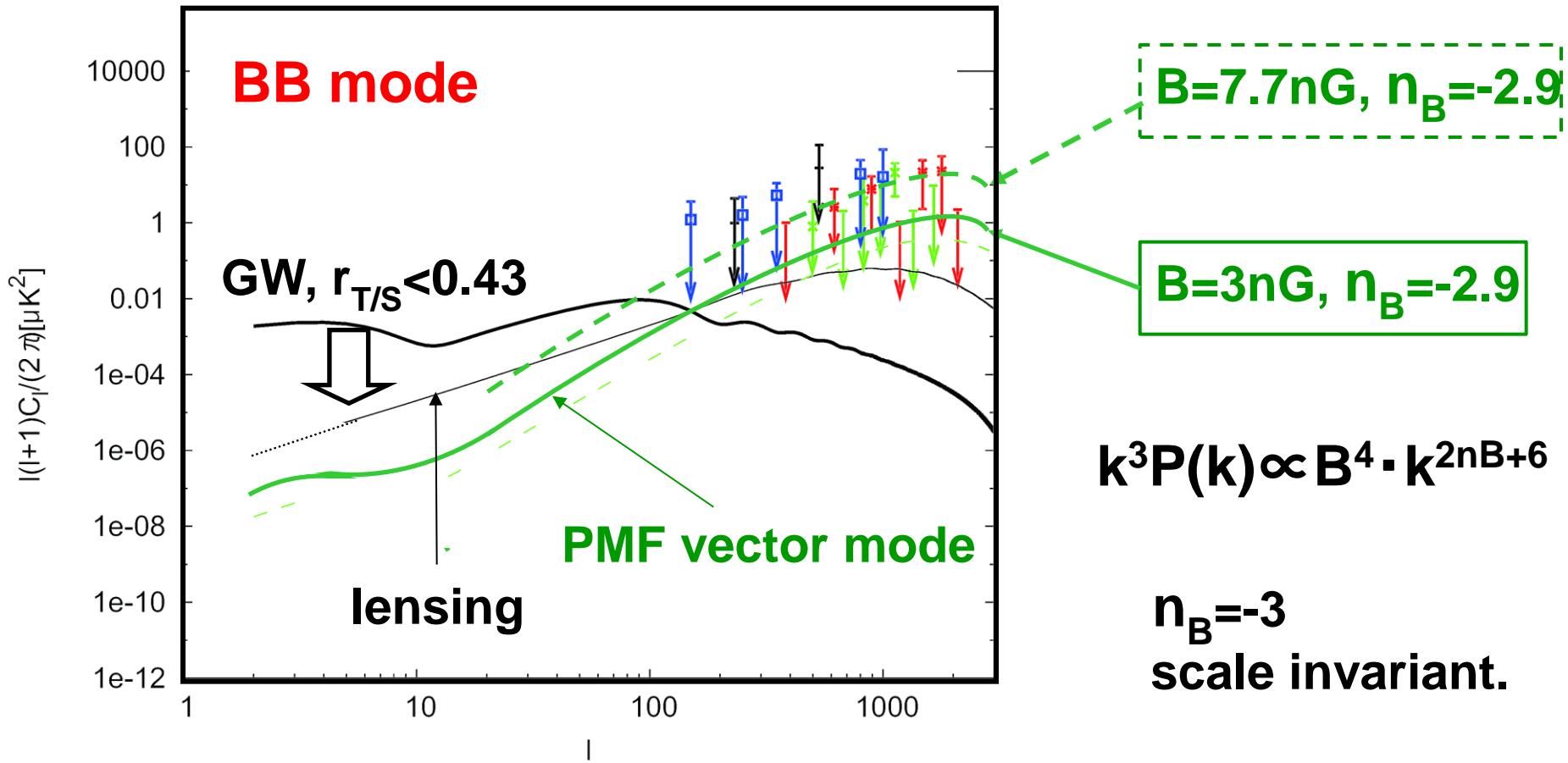


Expected Presence of Primordial Magnetic Field (PMF)

Yamazaki, Ichiki & Kajino, ApJ 825 (2006), L1

Yamazaki, Ichiki, Kajino, Mathews, PRD, 77, 043005 (2008)

Upper limit: $B = 3nG$, $m_v = 0$

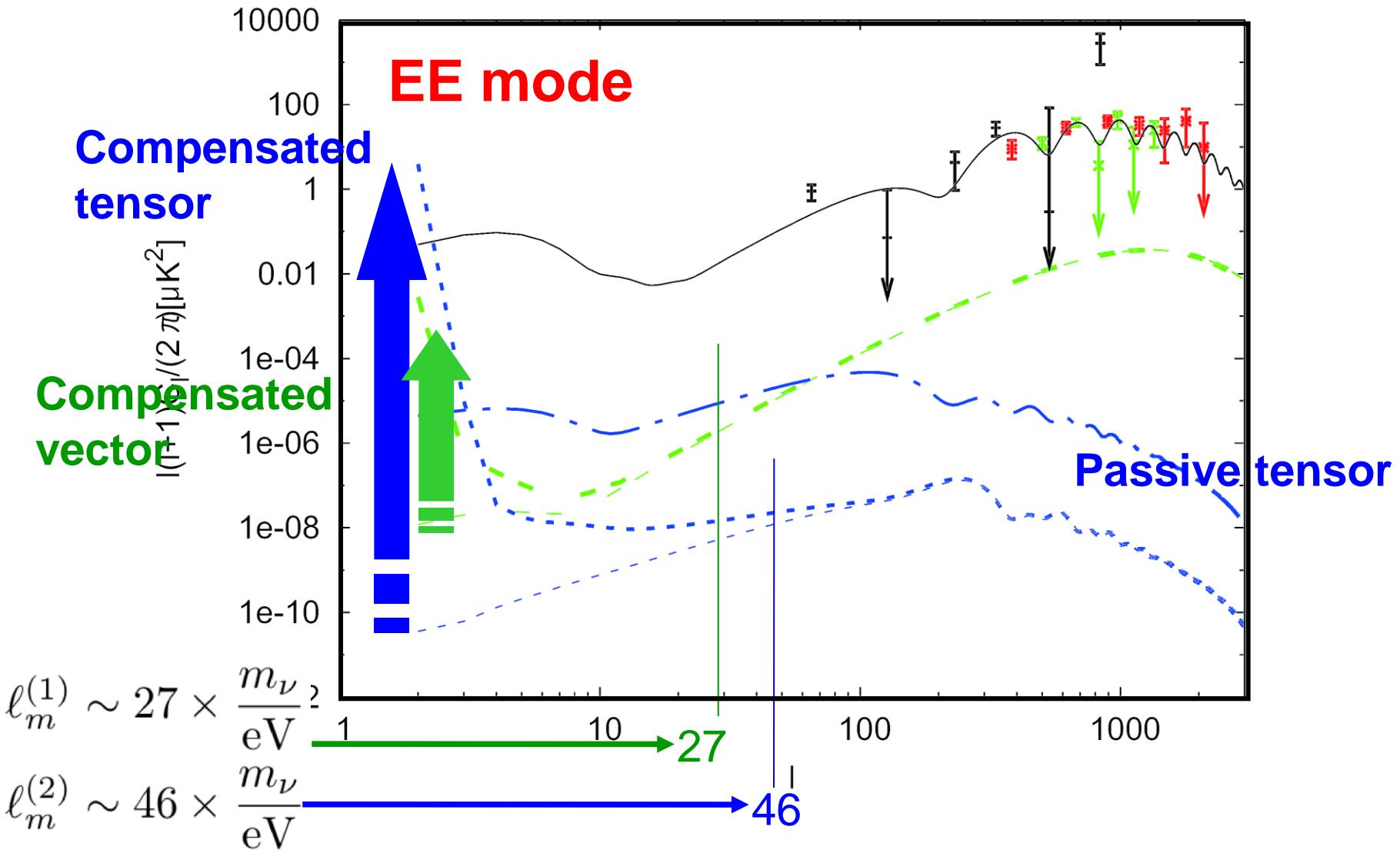


EE Polarization Mode with massive ν

$m_\nu = 1\text{eV}$

Kojima, Ichiki, Yamazaki, Kajino & Mathews,
 PR D78 (2008), 045010;
 Yamazaki, Kajino, Mathews & Ichiki, Phys. Rep.
 517 (2012) 141.

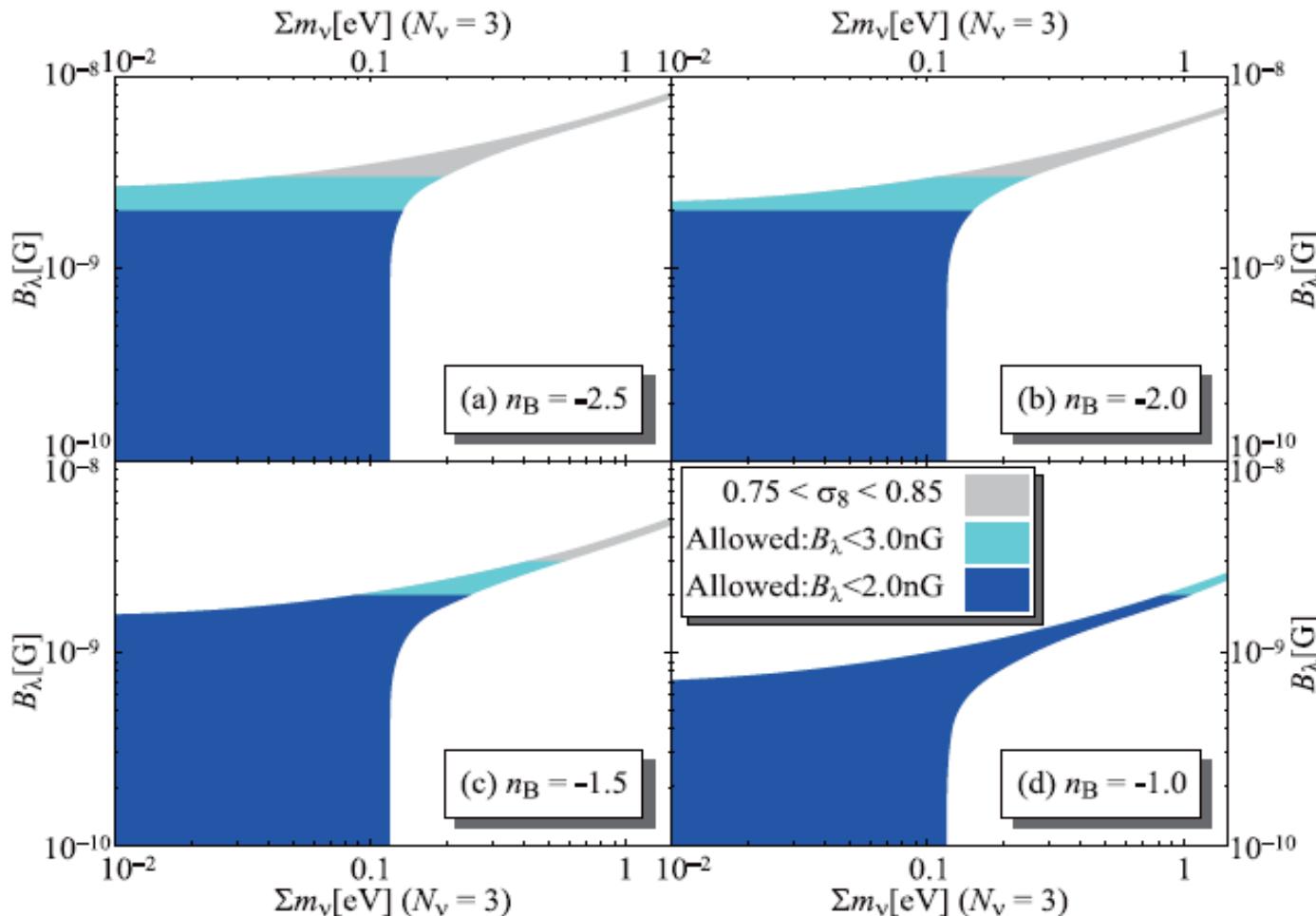
Extra Anisotropic-Stress of Primordial
 Magnetic Field; $B=3nG$, $n_B=-2.9$



Neutrino Mass Constraints from CMB + LSS

$\Sigma m_\nu < 0.2 \text{ eV}$

Yamazaki, Ichiki, Kajino, and Mathews,
PR D81 (2010), 103519;
D. Yamazaki, T. Kajino, G. J. Mathews,
and K. Ichiki, Phys. Rep. 517 (2012) 141.



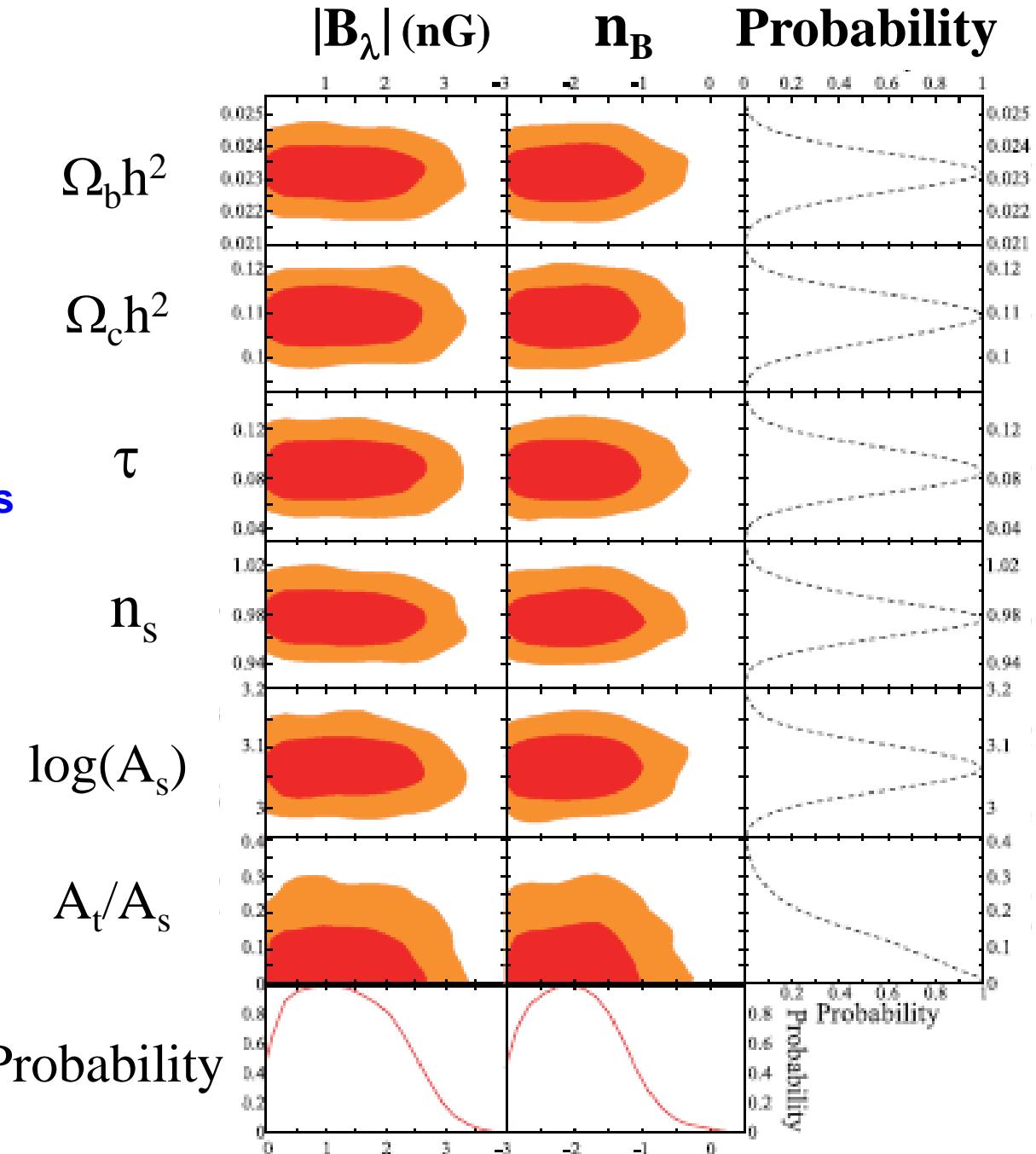
Likelihoodness and Probability

CMB Temperature and
Polarization Anisotropies
including
Primordial Magnetic Fields
and
Neutrino Mass

$\Sigma m_\nu < 0.2 \text{ eV}$

D. Yamazaki, K. Ichiki, T. Kajino
& G.J. Mathews,
PR D81 (2010), 103519.

D. Yamazaki, T. Kajino, G.J.
Mathews & K. Ichiki,
Phys. Rep. 517 (2012) 141.



BBN Constraint on Magnetic Moment of Massive Neutrinos

M. Kusakabe, A. B. Balantekin, T. Kajino & Y. Pehlivan, PR D87 (2013), 085045.

Cosmological: $\nu \rightarrow \nu' + \gamma_{\text{NT}}$ Parent ν could be massive!
 sterile ($m_\nu \neq 0, \mu_\nu \neq 0$) sterile/active

γ_{NT} & BBN: $D(\gamma, p) n$, $^6\text{Li}(\gamma, d) \alpha$, $^7\text{Li}(\gamma, t) \alpha$

Current Constraints

Laboratory: $\mu_\nu < 2.9 \times 10^{-11} \mu_B$

Astrophysical: $\mu_\nu < 3 \times 10^{-12} \mu_B$

Magnetic Moment of massive neutrino X

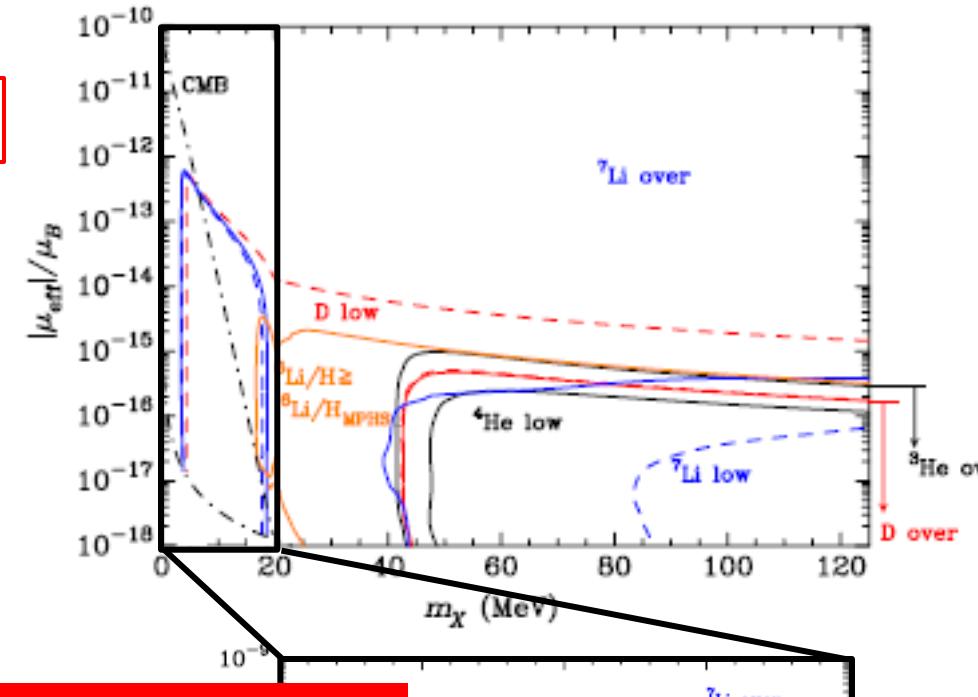
$$|\mu_{\text{eff}}|^2 = |\mu_{ij}|^2 + |\epsilon_{ij}|^2.$$

$$\boxed{\mu_\nu} \quad \boxed{\tau_X^{-1}} = \frac{|\mu_{ij}|^2 + |\epsilon_{ij}|^2}{8\pi} \left(\frac{m_i^2 - m_j^2}{m_i} \right)^3 \\ = 5.308 \text{ s}^{-1} \left(\frac{\mu_{\text{eff}}}{\mu_B} \right)^2 \left(\frac{m_i^2 - m_j^2}{m_i^2} \right)^3 \left(\frac{m_i}{\text{eV}} \right)^3$$

Decoupling Temp. is Max [1 MeV, $m_X/20$]

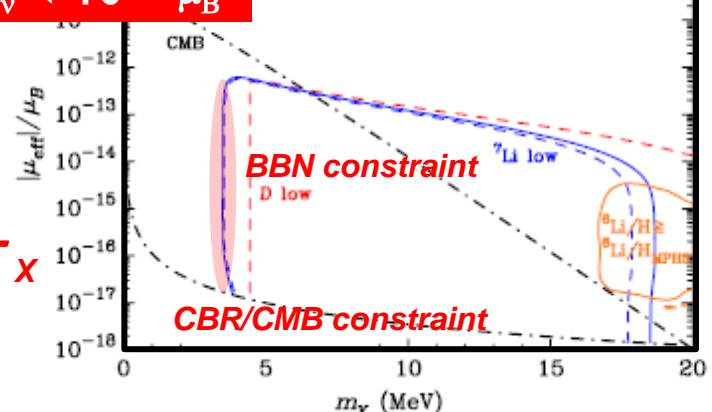
$$\boxed{\frac{n_X}{n_\gamma}} = \frac{4}{11} \frac{n_{dX}(m_X)}{n_\gamma(T_d)} = \frac{2\pi^2}{11\zeta(3)} \frac{n_{dX}(m_X)}{T_d^3}.$$

$$n_{dX}(m_X) = \frac{g_X}{2\pi^2} \int_0^\infty dp \frac{p^2}{\exp\left[\sqrt{p^2 + m_X^2}/T_d(m_X)\right] + 1}$$

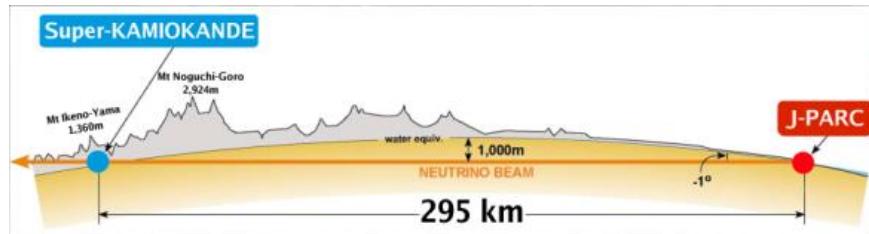


$$10^{-17} \mu_B < \mu_\nu < 10^{-12} \mu_B$$

longer τ_X



Long Baseline ν — T2K & MINOS (2011—)



$$\sin^2 2\theta_{13} = 0.1$$

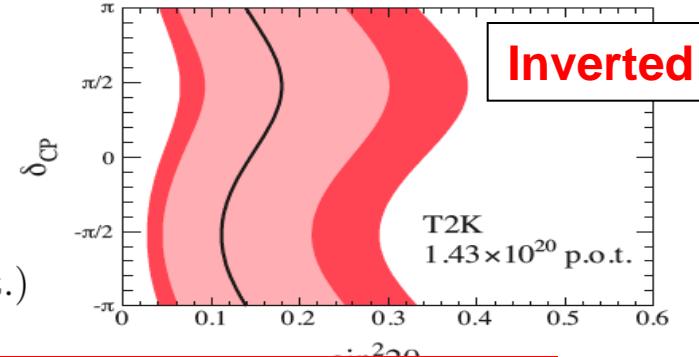
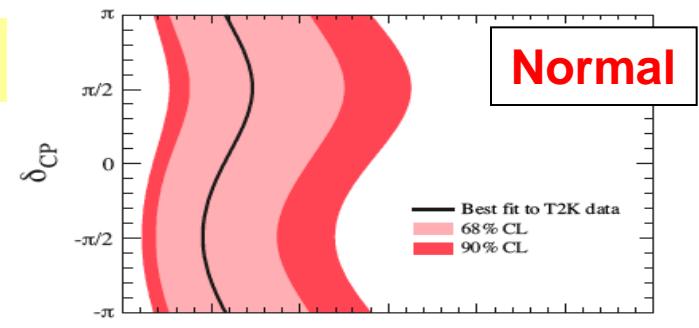
Daya Bay 2012 $\sin^2 2\theta_{13} = 0.092 \pm 0.016(\text{stat}) \pm 0.005(\text{syst})$

Rino (2012) $\sin^2 2\theta_{13} = 0.113 \pm 0.013(\text{stat.}) \pm 0.019(\text{syst.})$

Double Chooz
(2012)

Minos (2012)

T2K (2012) $0.03(0.04) < \sin^2 2\theta_{13} < 0.28(0.34)$

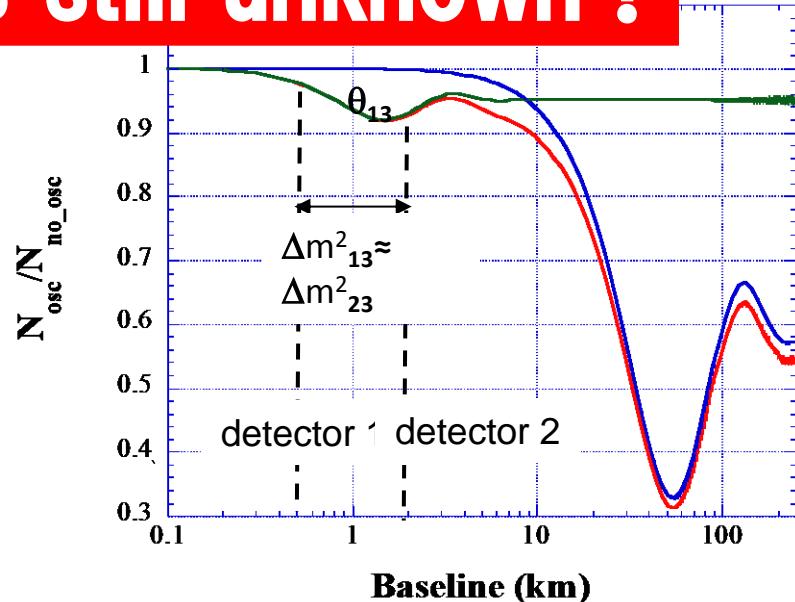


Mass hierarchy is still unknown !

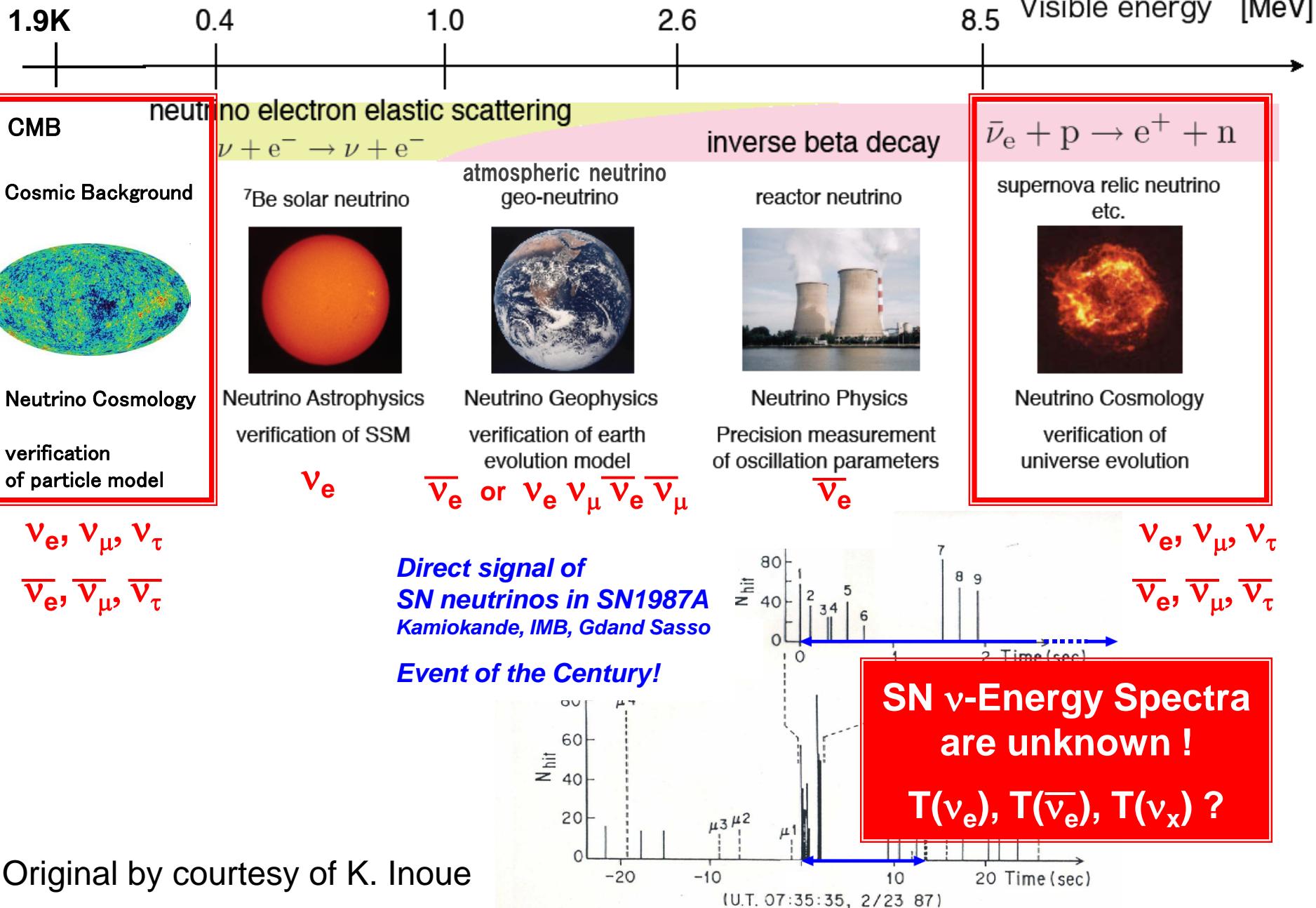
$\sin^2 2\theta_{23} = 1$

Reactor ν — RENO, Daya Bay & Double Chooz (2012—)

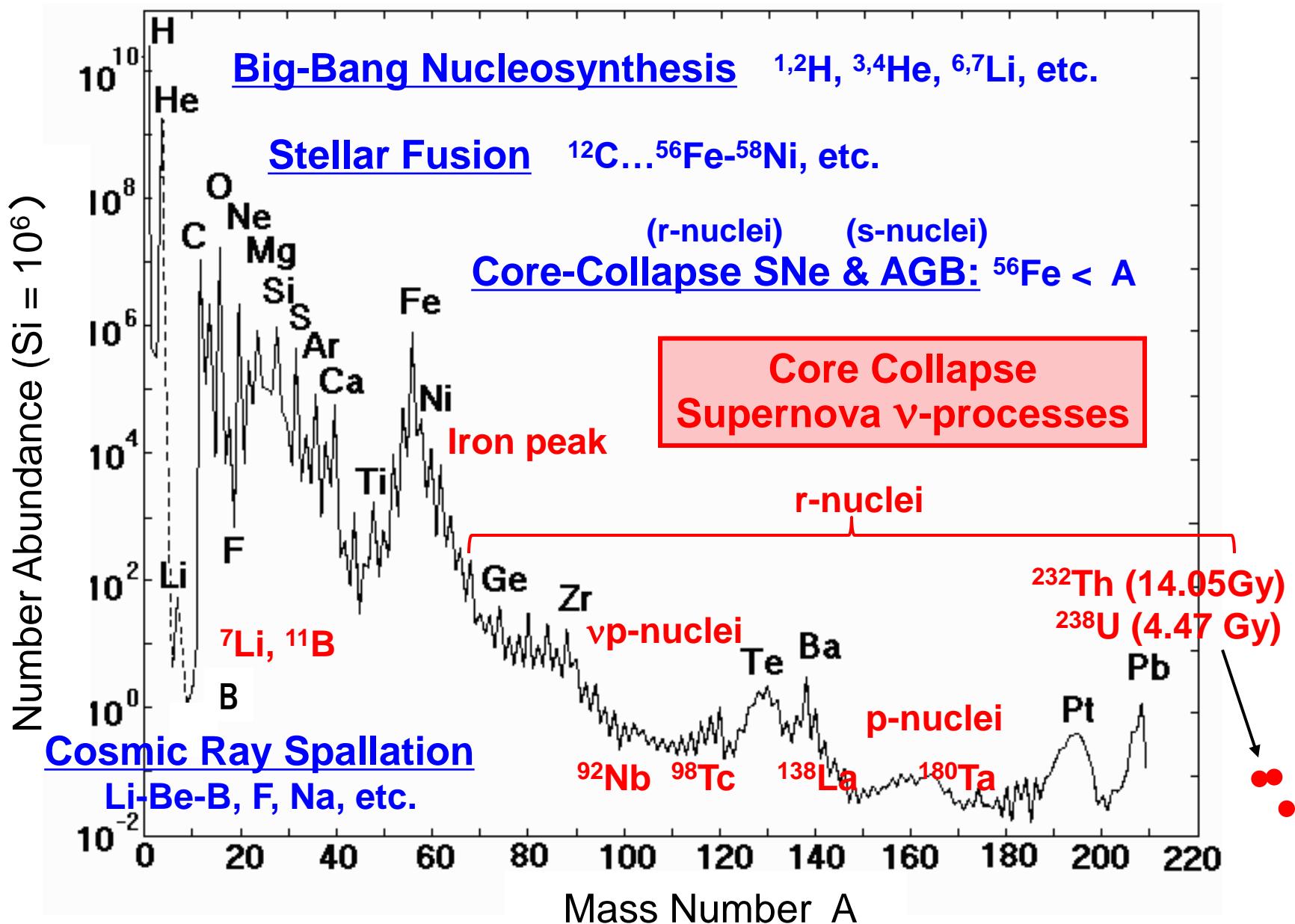
$$P_{ee} \approx 1 - \sin^2 2\theta_{13} \sin^2 \left(\frac{\Delta m_{31}^2 L}{4E_\nu} \right) - \cos^4 \theta_{13} \sin^2 2\theta_{12} \sin^2 \left(\frac{\Delta m_{21}^2 L}{4E_\nu} \right)$$



Various Neutrino-Sources in Nature/Culture



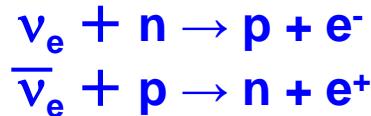
Solar System Abundance



R-process Nucleosynthesis

Otsuki, Tagoshi, Kajino and Wanajo, ApJ 533 (2000) ,424; Wanajo, Kajino, Mathews and Otsuki, ApJ 554 (2001),578.

Neutron-rich condition for successful r-process: $0.1 < Y_e < 0.5$



$$Y_e = \frac{p}{n+p} \approx \left(1 + \frac{L_{\bar{\nu}_e}}{L_{\nu_e}} \times \frac{\epsilon_{\bar{\nu}_e} - 2\Delta + 1.2\Delta^2/\epsilon_{\bar{\nu}_e}}{\epsilon_{\nu_e} + 2\Delta + 1.2\Delta^2/\epsilon_{\nu_e}} \right)^{-1}$$

$$\epsilon_\nu = 3.15 T_\nu$$

$$T_{\nu e} = 3.2 \text{ MeV}, \quad T_{\bar{\nu} e} = 4 \text{ MeV}$$

Theoretical Challenge:

1) Astrophysical Sites ?

- ν-wind SNe
- MHD jet SNe
- NS mergers (short GRB)
- long GRBs

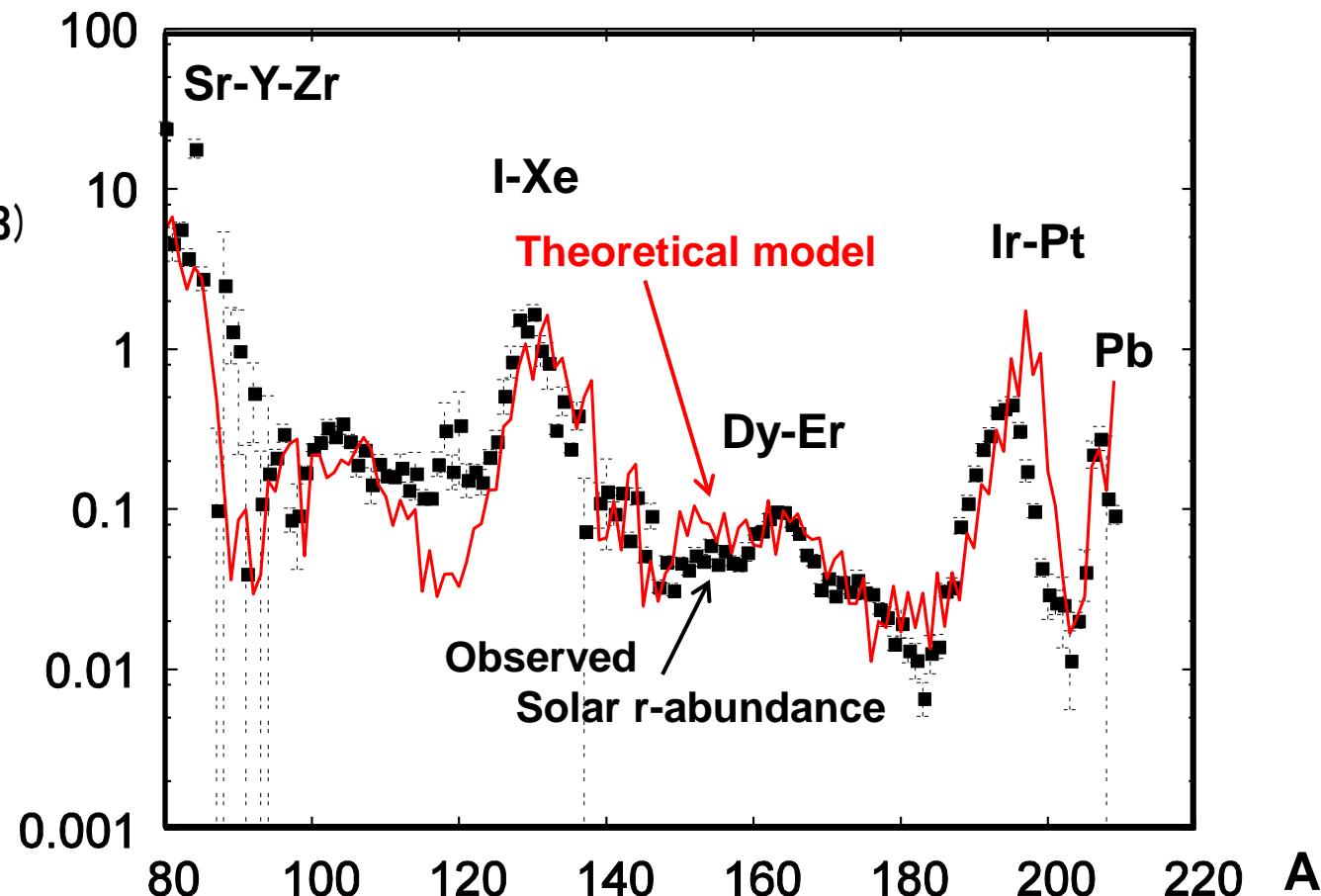
2) Neutrino effects ?

$$Y_e > 0.5 ?$$

Roberts, Reddy and Shen
(PR C86, 065803, 2012)
pointed out

$$Y_e < 0.5 !$$

for nucleon potential
and Pauli blocking
effects.



^{92}Nb also has SN- ν Origin !

Hayakawa et al., ApJ 778 (2013) L1.

$^{138}\text{La}, ^{180}\text{Ta}$, too !

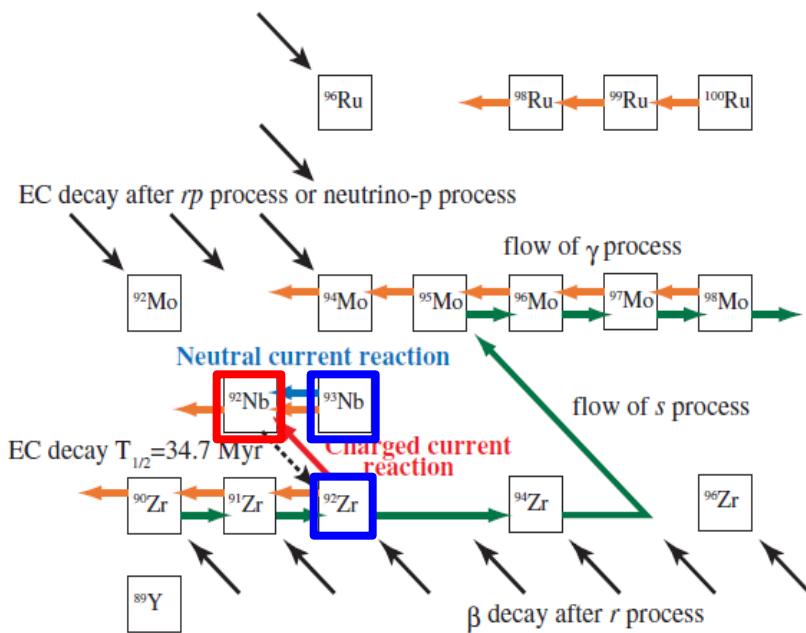
Hayakawa, et al., PR C81 (2010) 052801®,
Hayakawa et al., PR C82 (2010) 058801.

^{92}Nb ($\tau_{1/2}=3.47\times 10^7$ y): Unique Chronometer of SN ν -Process

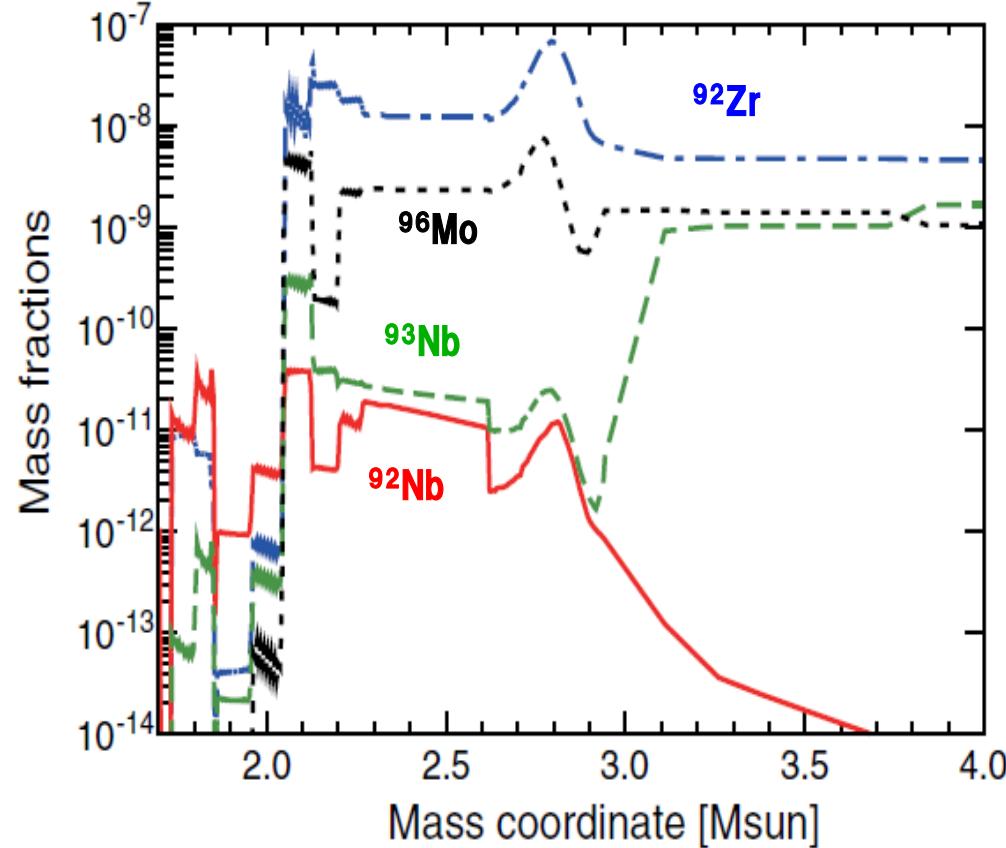
Isotopic anomaly in meteoritic $^{92}\text{Zr}/^{93}\text{Nb}$:

$$\Delta = 1 \times 10^6 - 3 \times 10^7 \text{ y}$$

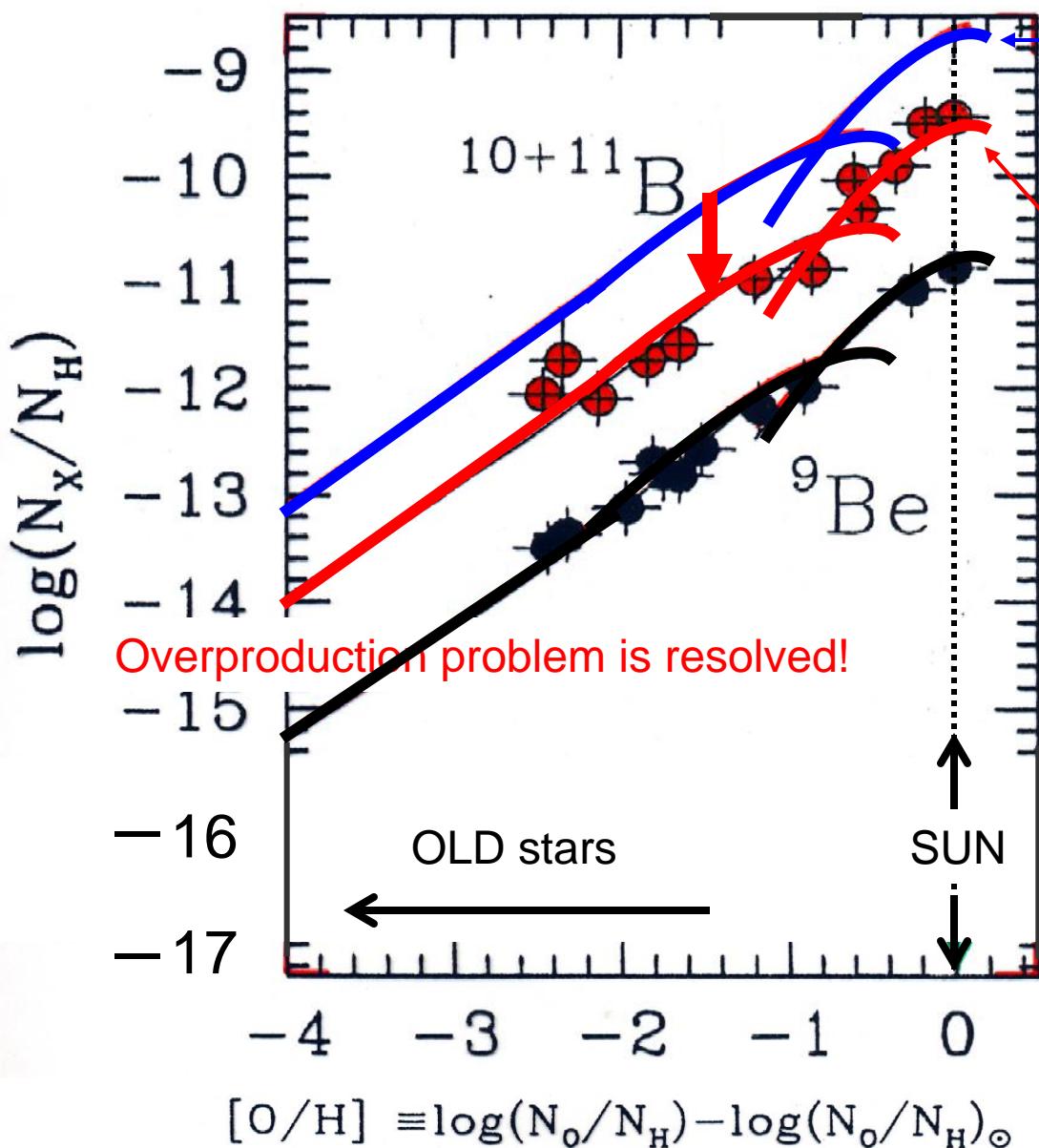
Time duration after the last nearby Supernova to the Solar-System (protosolar cloud) formation



$$T_{\nu e} = 3.2 \text{ MeV}, \quad T_{\bar{\nu} e} = 4 \text{ MeV}$$



Galactic Chemical Evolution of ${}^9\text{Be}$ & ${}^{10,11}\text{B}$



Livermore Model

$$T_{\nu_{\mu,\tau}} = 8 \text{ MeV}$$

Woosley -Weaver 1995, ApJS 101, 181.

$$\sigma \propto E_\nu^2$$

$$T_{\nu_{\mu,\tau}} = 6 \text{ MeV}$$

Consistent with SN1987A

Yoshida, Kajino & Hartmann 2005,
PRL 94 (2005), 231101.

Consistent with SN1987A

${}^9\text{Be}$:

– Galactic Cosmic Rays

${}^{10,11}\text{B}$ + ${}^{11}\text{B}$:

– Galactic Cosmic Rays

– Supernova ν -process

Yoshii, Kajino, Ryan, 1997, ApJ 486, 605.
Ryan, Kajino, Suzuki, 2001, ApJ 549, 55.

Supernova ν -Process: $^7\text{Li}, ^{11}\text{B}, ^{92}\text{Nb}, ^{138}\text{La}, ^{180}\text{Ta}$

^{14}N



Shell Model:

Yoshida, Suzuki, Chiba, Kajino, et al., ApJ 686 (2008), 448:

Suzuki and Kajino, J. Phys. G40 (2013), 083101:

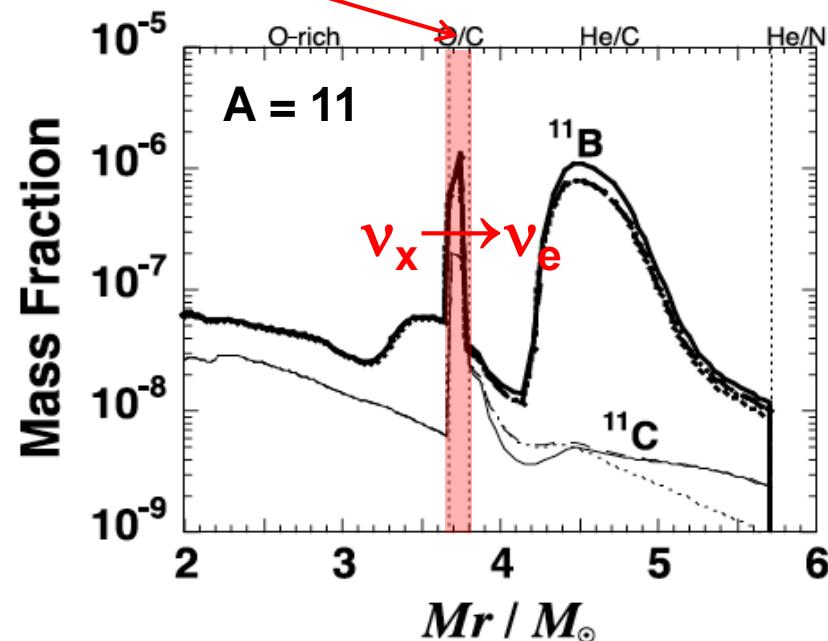
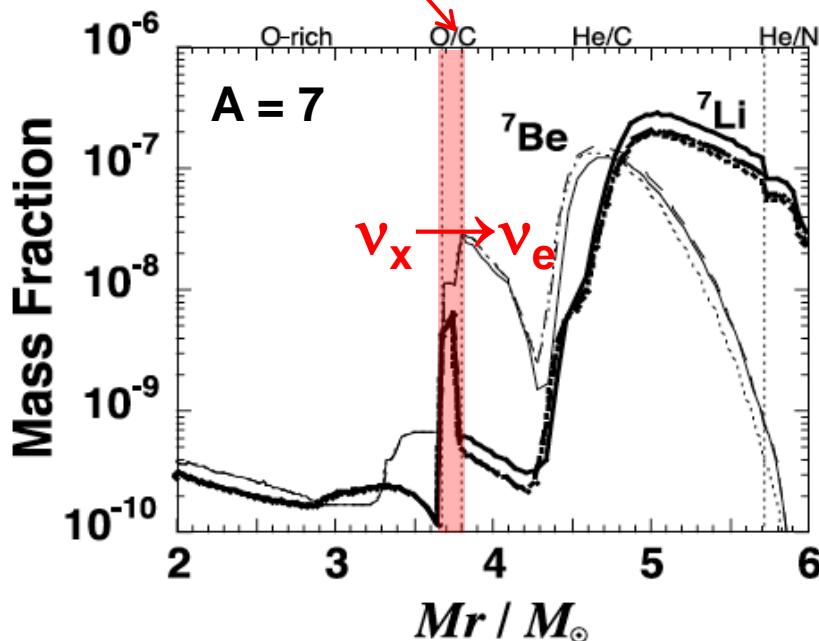
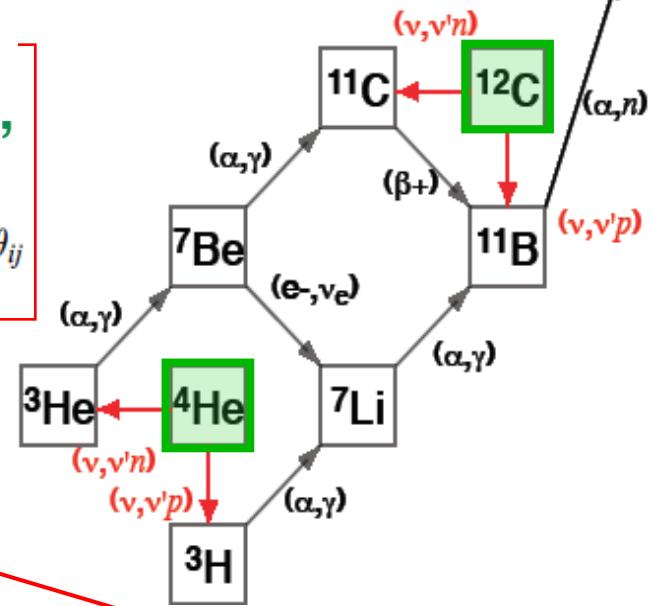
QRPA:

Choun, et al.

J. Phys. G37 (2010), 055101; PRC82 (2010), 035504;

MSW high-density resonance is located at the bottom of He/C shell.

$$= 6.55 \times 10^6 \left(\frac{\Delta m_{ji}^2}{1 \text{ eV}^2} \right) \left(\frac{1 \text{ MeV}}{\varepsilon_\nu} \right) \cos 2\theta_{ij}$$



Nucleosynthetic Constraints on ν -Temperatures!

- R-process (neutron-richness) $\rightarrow T\nu_e = 3.2 \text{ MeV}, T\bar{\nu}_e = 4 \text{ MeV}$

Otsuki, Tagoshi, Kajino and Wanajo, *Astrophys. J.* 533 (2000) 424.

Wanajo, Kajino, Mathews and Otsuki, *Astrophys. J.* 554 (2001) 578.

Roberts, Reddy and Shen, *Phys. Rev. C* 86 (2012) 065803.

- P-process; $^{180}\text{Ta}/^{138}\text{La}, ^{92}\text{Nb}$ (CC- ν) $\rightarrow T\nu_e = T\bar{\nu}_e = 4 \text{ MeV}$

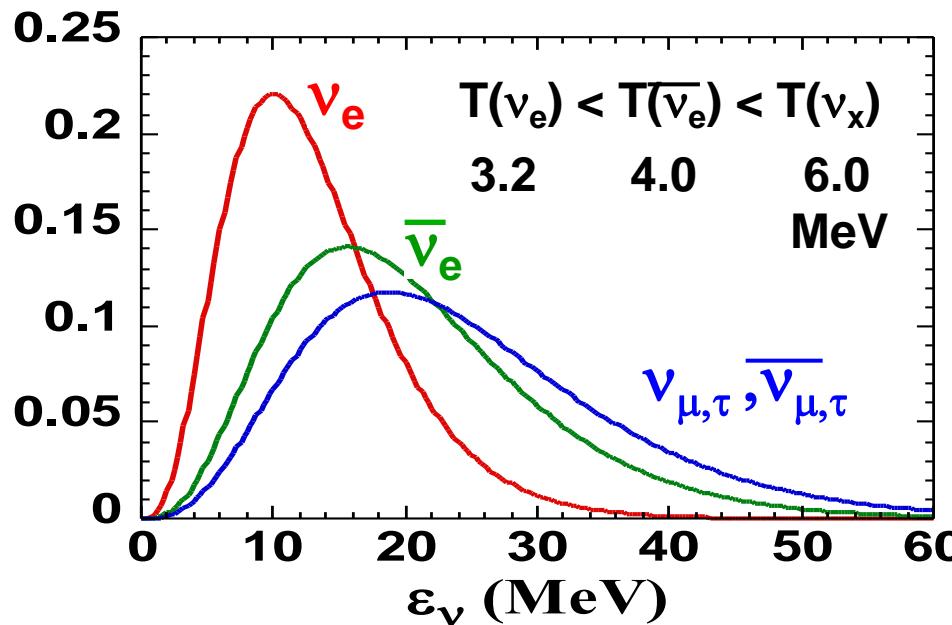
Hayakawa, et al., *Phys. Rev. C* 81 (2010) 052801®; *Phys. Rev. C* 82 (2010) 058801.

Hayakawa et al., *Astrophys. J. Lett.* 778 (2013) L1.

- GCE; $^{6,7}\text{Li}-^9\text{Be}-^{10,11}\text{B}$ & Meteoritic $^{11}\text{B}/^{10}\text{B}$ (NC- ν) $\rightarrow T\nu_{x=\mu,\tau} = 6 \text{ MeV}$

Yoshida, Kajino & Hartman, *Phys. Rev. Lett.* 94 (2005) 231101.

Suzuki & Kajino, *J. Phys. G* 37 (2010), 055101.



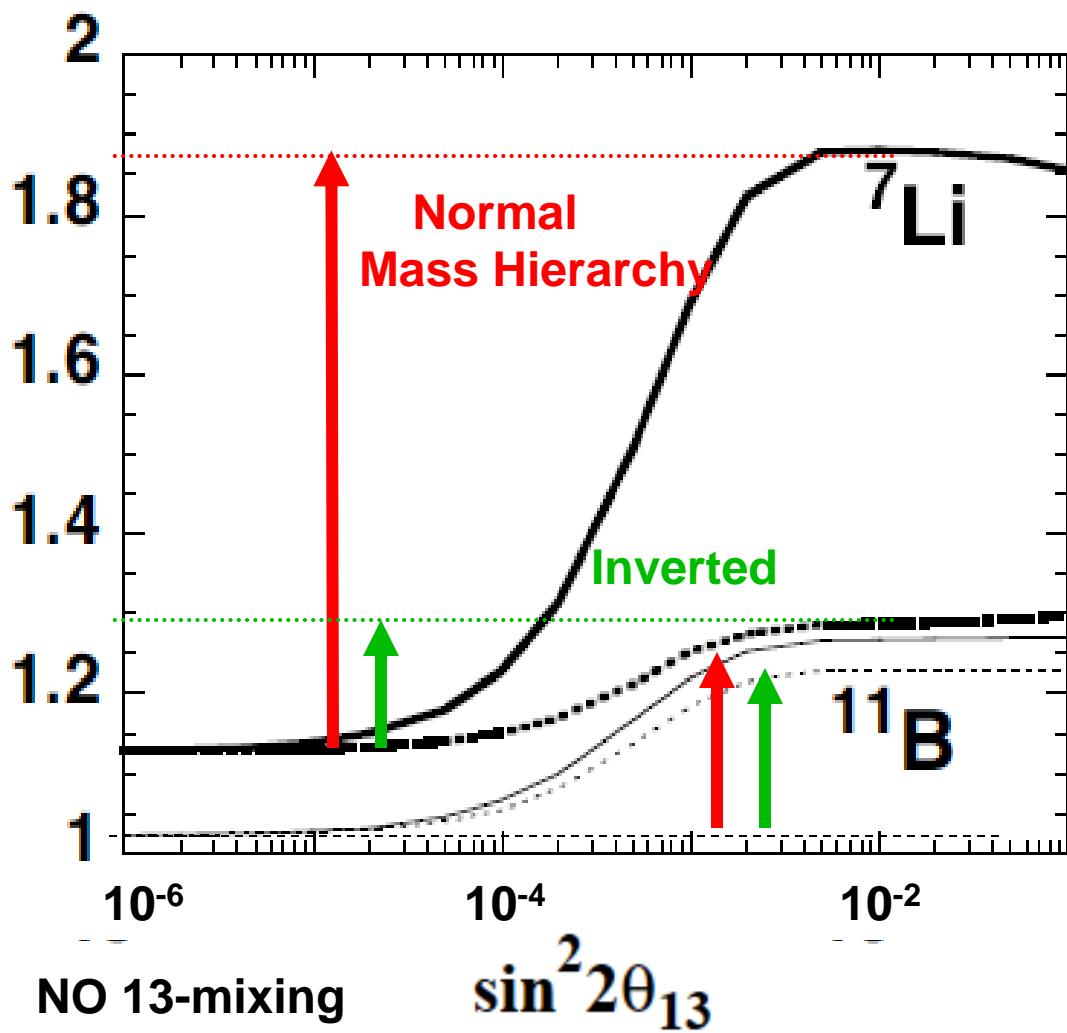
Variation of T's from different Supernova Models is taken into account !

Mathews, Kajino, Aoki, Fujita & Pitts,
Phys. Rev. D 85 (2012) 105023.

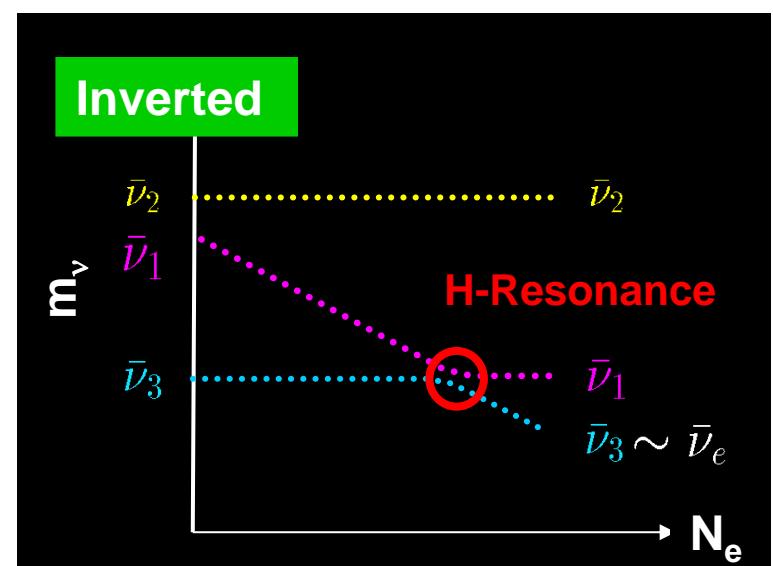
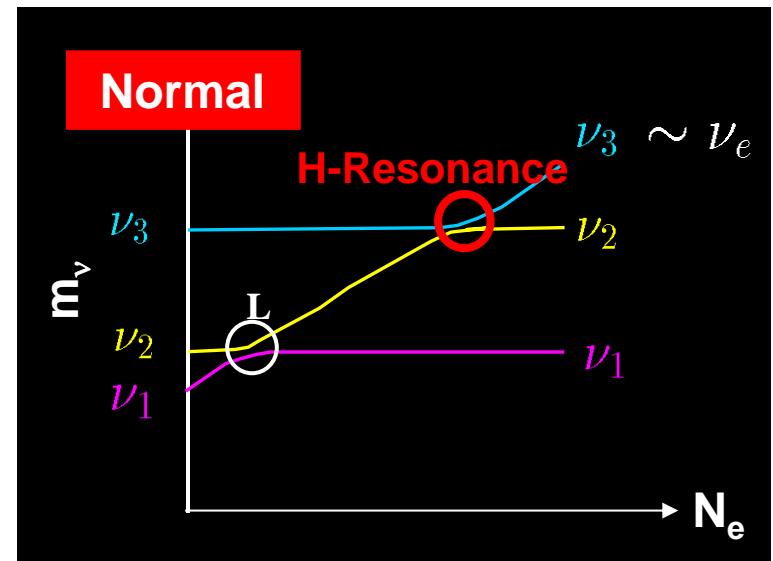
larger effect !

$$T_{\nu e} < T_{\bar{\nu} e} < T_{\nu \mu \tau, \bar{\nu} \mu \tau}$$

smaller effect !



Yoshida, Kajino, Yokomakura, Kimura, Takamura & Hartmann,
PRL 96 (2006) 09110; ApJ 649 (2006), 349.



**Exploring the neutrino mass hierarchy probability with meteoritic supernova material,
 ν -process nucleosynthesis, and θ_{13} mixing**

G. J. Mathews,^{1,2} T. Kajino,^{2,3} W. Aoki,² W. Fujiya,⁴ and J. B. Pitts⁵

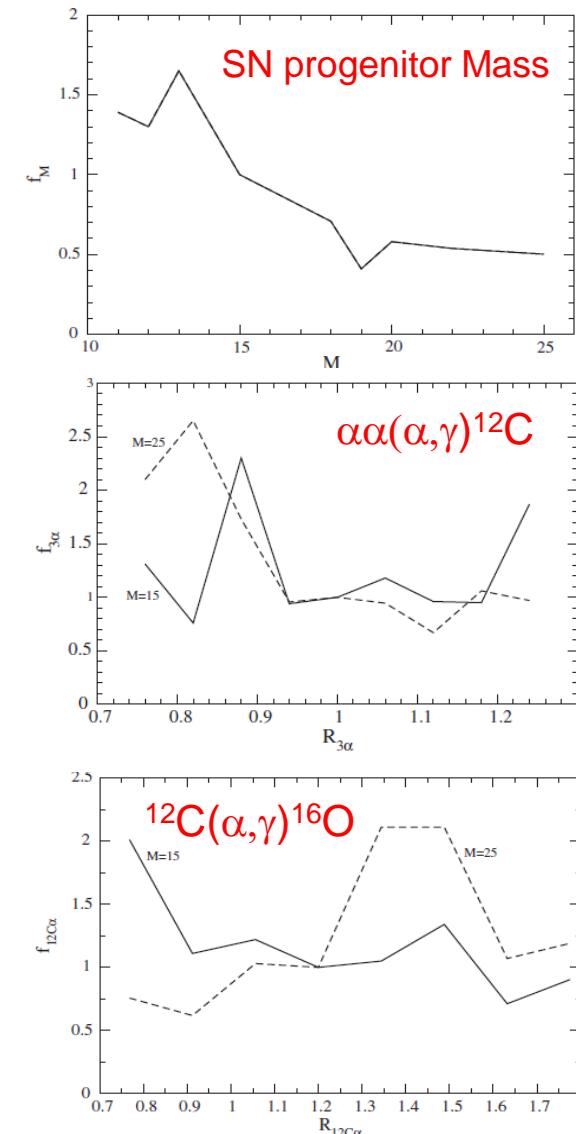
Bayesian Analysis, including astrophysical model dependence on SN progenitor masses, ν -temps. ($T_{\nu e}, T_{\nu \bar{e}}, T_{\nu \mu \tau}, \overline{\nu \mu \tau}$) and nuclear input data.

$$P(M_i|D) = \frac{P(D|M_i)P(M_i)}{\sum_j P(D|M_j)P(M_j)}$$

$$\begin{aligned} P(D|M_i) &= \int dE dZ da_k P(E, Z, D|M_i, a_k) P(a_k|M_i) \\ &= \int dE dZ da_k P(D|M_i, a_k, E, Z) P(Z, E|M_i, a_k) P(a|M_i) \end{aligned}$$

TABLE I: Parameter likelihood functions $P(a_k|M_i)$.

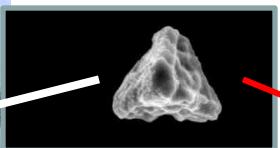
Parameter a_k	prior			reference
$\sin^2 2\theta_{13}$	$e^{-(x-x_0)/2\sigma_x^2}$	$x_0 = 0.92$	$\sigma_x = 0.017$	[7]
$R_{3\alpha}$	$e^{-(x-x_0)/2\sigma_x^2}$	$x_0 = 1.0$	$\sigma_x = 0.12$	[35]
$R_{12C\alpha}$	$e^{-(x-x_0)/2\sigma_x^2}$	$x_0 = 1.2$	$\sigma_x = 0.25$	[36]
$M_{prog}(\text{M}_\odot)$	$m^{-2.65}$	$m_{min} = 10$	$m_{max} = 25$	[37]
$T_\nu(\text{MeV})$	Top hat	$T_\nu = 3.2 - 6.5$	(see text)	[15]



Murchison Meteorite



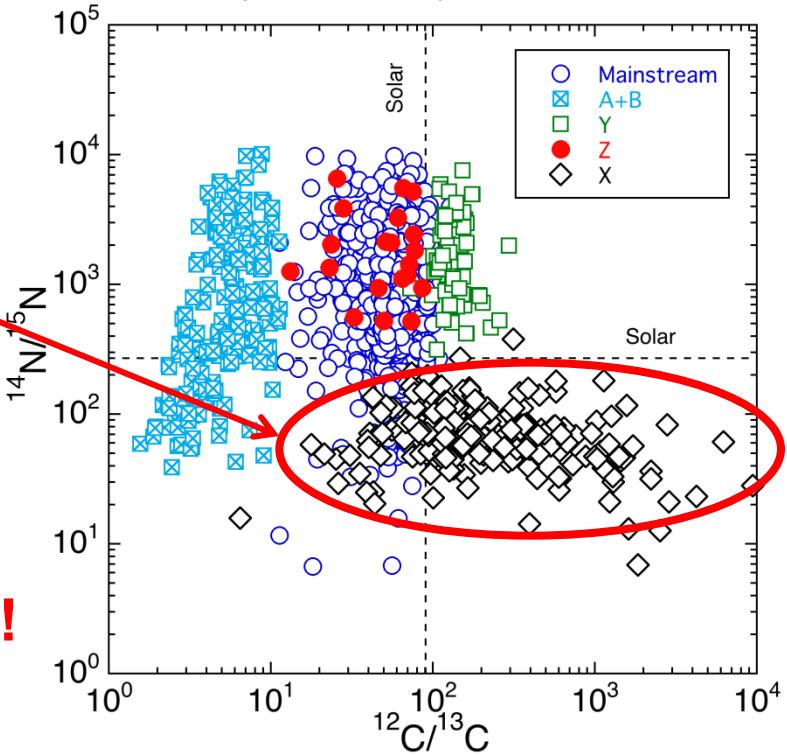
SiC X-grains



- $^{12}\text{C}/^{13}\text{C} > \text{Solar}$
- $^{14}\text{N}/^{15}\text{N} < \text{Solar}$

- Enhanced ^{28}Si
- Decay of ^{26}Al ($t_{1/2}=7 \times 10^5 \text{ yr}$), ^{44}Ti ($t_{1/2}=60 \text{ yr}$)

By courtesy of S. Amari



SiC X-grains are made of cc-SN Dust !

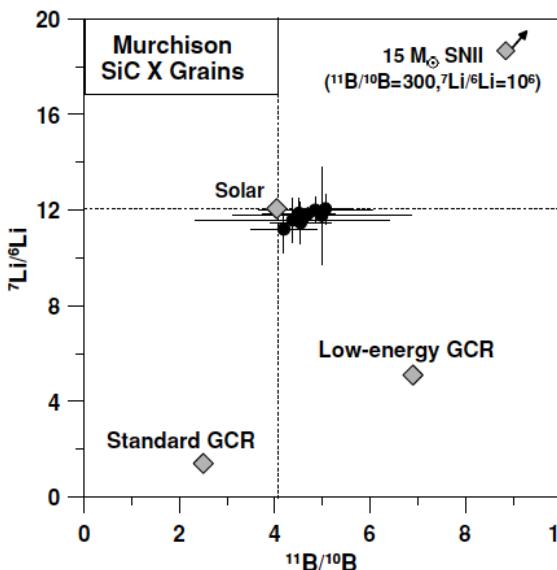
Fujiya, Hoppe and Ott (2011, ApJ 730, L7)
discovered ^{11}B and ^7Li isotopes in 13 SiC X-grains.

Table 1

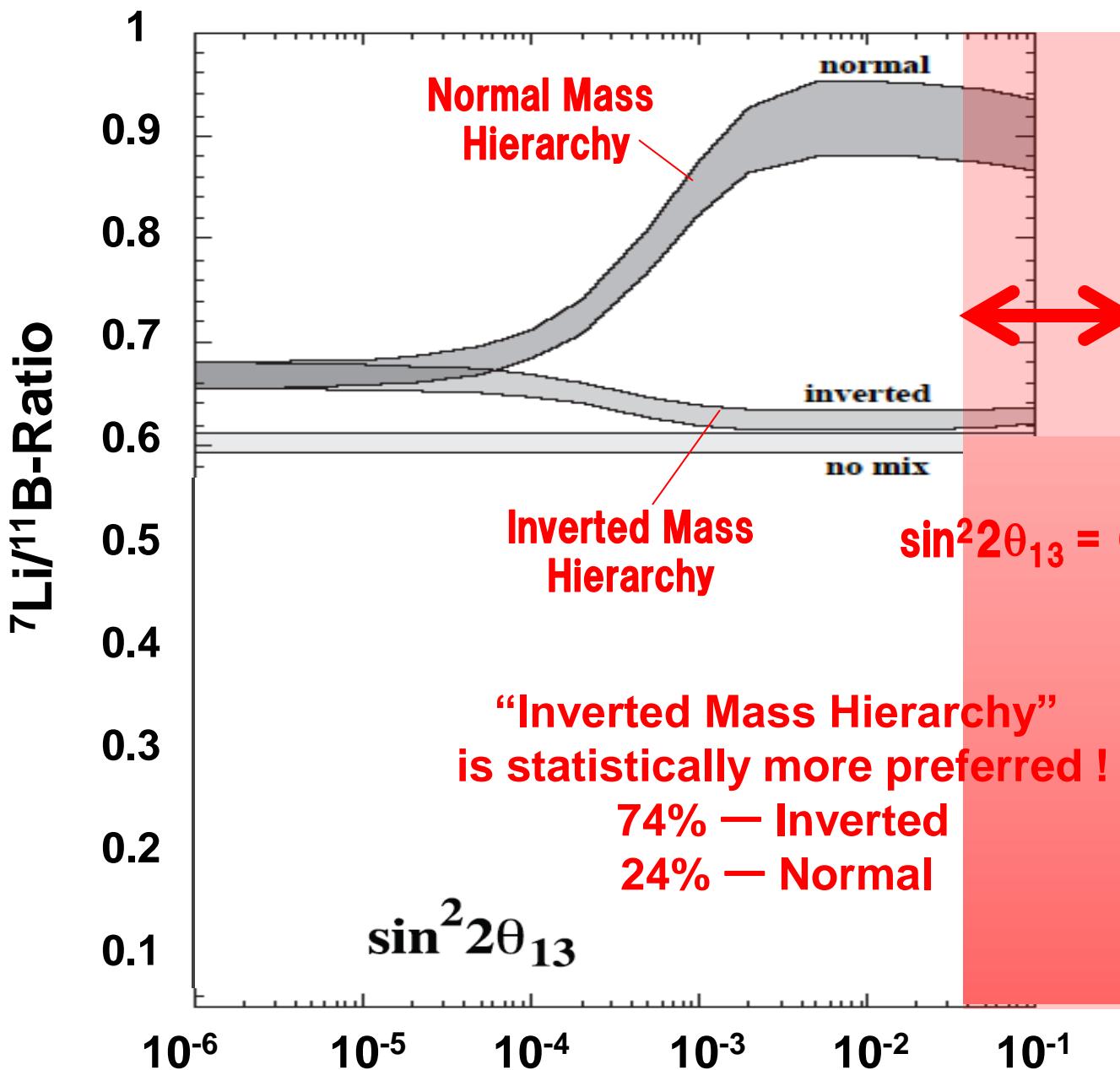
C-, Si-, Li-, and B-isotopic Compositions of SiC X Grains from the Murchison Meteorite

Grain	Size (μm)	$^{12}\text{C}/^{13}\text{C}$	$\delta^{29}\text{Si}^{\text{a}}$ (‰)	$\delta^{30}\text{Si}^{\text{a}}$ (‰)	$^7\text{Li}/^6\text{Li}$	$^{11}\text{B}/^{10}\text{B}$	Li/Si (10^{-5})	B/Si (10^{-5})
Single X grains								
X1	0.6	114 ± 2	-178 ± 11	-265 ± 9	11.87 ± 0.63	4.51 ± 0.77	9.69	3.33
X2	1.2	128 ± 2	-377 ± 11	-261 ± 10	12.06 ± 0.62	5.06 ± 0.58	23.8	18.8
X3	1.5	244 ± 5	-205 ± 10	-297 ± 7	11.48 ± 0.86	4.54 ± 0.63	1.76	1.92
X4	1.0	241 ± 6	-556 ± 10	-245 ± 9	12.00 ± 0.56	4.85 ± 1.19	24.8	3.31
X9	0.6	38 ± 1	-361 ± 10	-394 ± 8	11.20 ± 1.01	4.19 ± 0.70	10.8	11.4
X11	0.8	326 ± 14	-358 ± 12	-432 ± 11	11.78 ± 2.03	4.99 ± 1.88	3.66	3.00
X13	0.7	345 ± 6	-261 ± 10	-424 ± 7	11.59 ± 0.93	4.37 ± 2.04	10.7	1.14
Average					11.83 ± 0.29	4.68 ± 0.31		
X grains + other nearby/attached SiC grains								
X5	34 \pm 1	-226 ± 11	-120 ± 10	12.21 ± 0.41	4.36 ± 0.40	40.2	18.8	
X6	88 \pm 1	-236 ± 11	-189 ± 9	13.06 ± 1.36	3.83 ± 0.27	2.15	14.2	
X7	78 \pm 1	-281 ± 11	-208 ± 10	11.20 ± 2.40	11.47 ± 6.36	8.28	9.48	
X8	76 \pm 1	-223 ± 10	-266 ± 8	11.29 ± 0.64	4.27 ± 0.29	4.80	12.4	
X12	83 \pm 1	-271 ± 11	-242 ± 10	11.54 ± 0.52	4.13 ± 0.46	24.3	14.2	
Average					11.90 ± 0.28	4.16 ± 0.17		
Solar	89	0	0	12.06	4.03	5.6	1.9	

Note. ${}^a\delta\text{Si} = [({}^i\text{Si}/{}^{28}\text{Si})/({}^i\text{Si}/{}^{28}\text{Si})_{\odot} - 1] \times 1000$.



MSW Effect & ν Mass Hierarchy



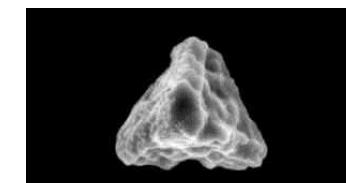
Mathews, Kajino, Aoki & Fujiya, Phys. Rev. D85, 105023 (2012).

Long Baseline Exp. in 2011:

- T2K (Kamioka)
- MINOS

Reactor Exp. in 2012:

- Double CHOOZ
- Daya Bay
- First Detection of ${}^7\text{Li}$ RENO (KOREA)
- SNO grains



W. Fujiya, P. Hoppe, & U. Ott, ApJ 730, L7 (2011).

A New Method to constrain EOS & ν -Oscllation

G.J. Mathews, J. Hidaka, T. Kajino & J. Suzuki, ApJ (2014), submitted.

THE ASTROPHYSICAL JOURNAL, 738:154 (16pp), 2011 September 10

THE COSMIC CORE-COLLAPSE SUPERNOVA RATE DOES NOT MATCH THE MASSIVE-STAR FORMATION RATE

SHUNSAKU HORIUCHI^{1,2}, JOHN F. BEACOM^{1,2,3}, CHRISTOPHER S. KOCHANEK^{2,3}, JOSE L. PRIETO^{4,5},
K. Z. STANEK^{2,3}, AND TODD A. THOMPSON^{2,3,6}

Supernova Rate Problem/Discrepancy

SFR of Massive Stars at birth

SNR: Supernova Explosions at death !

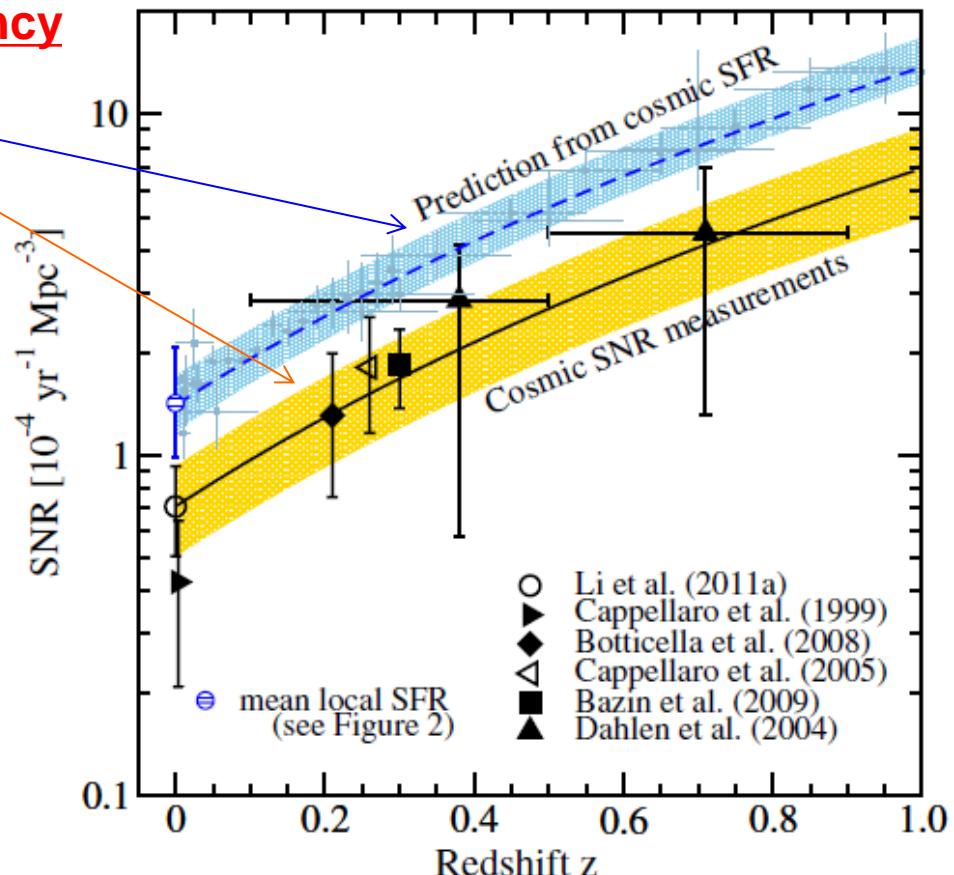
50% Massive Stars, missing !

Expected Reasons:

Half was evolved into too dark SNe
to detect!

1. Failed SNe ($< 25 M_\odot$ BH formation)
2. Faint ONeMg-SNe (8-10 M_\odot)

or the mass function changed!



Electron-capture SN (Faint SnNe)

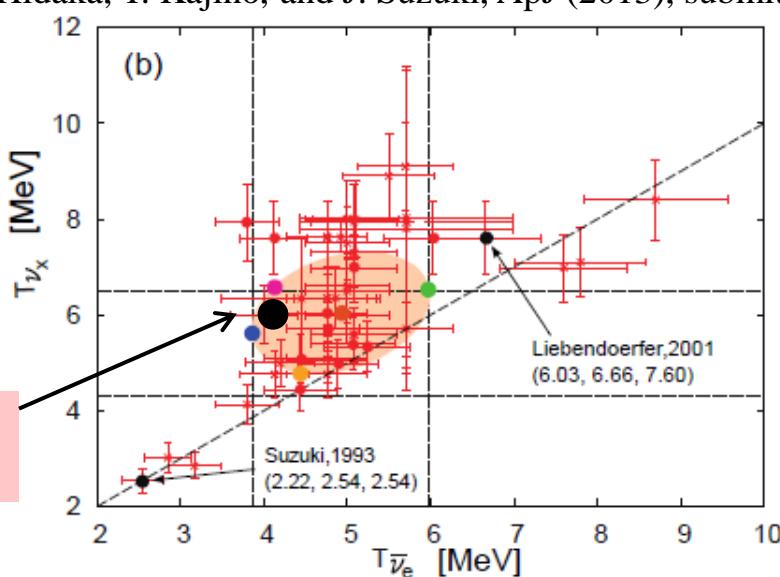
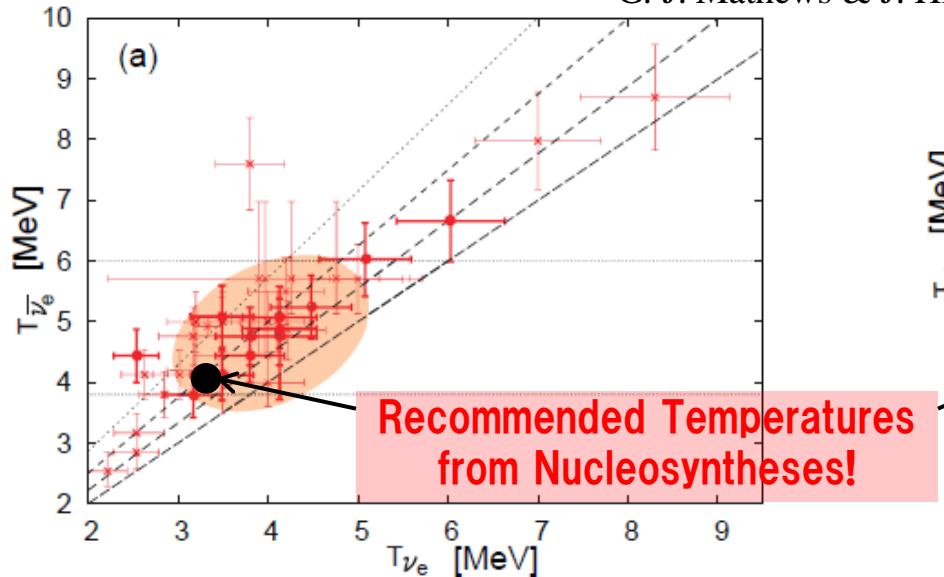
Normal SNe (Neutron Star formation)

Failed SNe (Black Hole formation)

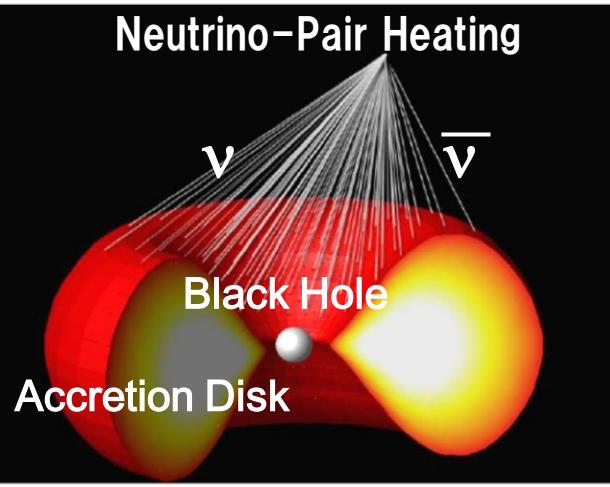
detail	ONeMg SN	CC-SN	fSN(SH EOS)	fSN(LS EOS)	GRB
mass(M_{\odot})	(8 ~ 10)	8 ~ 25(10~25)	25 ~ 125 (99.96%)	25 ~ 125 (99.96%)	25 ~ 125 (0.04%)
Remnant Phenomenon	Neutron Star Supernova	Neutron Star Supernova	Black Hole	Black Hole	Black Hole
T_{ν_e} (MeV)	3.0	3.2	5.5	7.9	3.2
$T_{\bar{\nu}_e}$ (MeV)	3.6	4.0	5.6	8.0	5.3
T_{ν_x} (MeV)	3.6	6.0	6.5	11.3	4.4
$E_{\nu_e}^{total}$ (erg)	3.3×10^{52}	5.0×10^{52}	5.5×10^{52}	8.4×10^{52}	1.7×10^{53}
$E_{\bar{\nu}_e}^{total}$ (erg)	2.7×10^{52}	5.0×10^{52}	4.7×10^{52}	7.5×10^{52}	3.2×10^{53}
$E_{\nu_x}^{total}$ (erg)	1.1×10^{53}	5.0×10^{52}	2.3×10^{52}	2.7×10^{52}	1.9×10^{52}
Δt	few s	few s	$\sim 0.5s$	$\sim 0.5s$	$\sim 10s$

- ONeMg SN Hudepohl, et al., PRL 104 (2010) Shen-EOS LS-EOS
- fSN (failed SN) Sumiyoshi, et al., ApJ 688 (2008), 1176; Shen et al. (1998), Lattimer & Swesty (1991)
- CC-SN Yoshida, et al., ApJ 686 (2008), 448; Suzuki & Kajino, J. Phys. G40 (2013) 83101 +

G. J. Mathews & J. Hidaka, T. Kajino, and J. Suzuki, ApJ (2013), submitted.



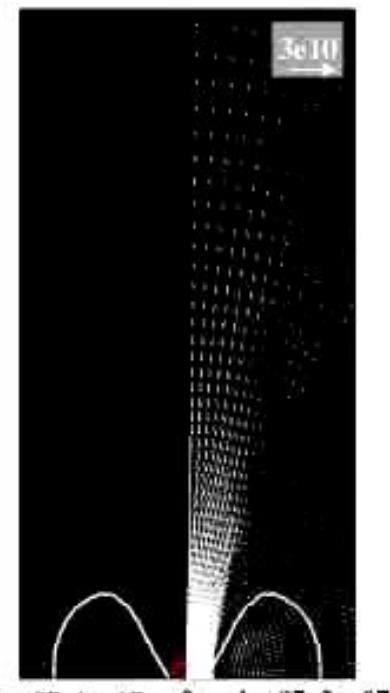
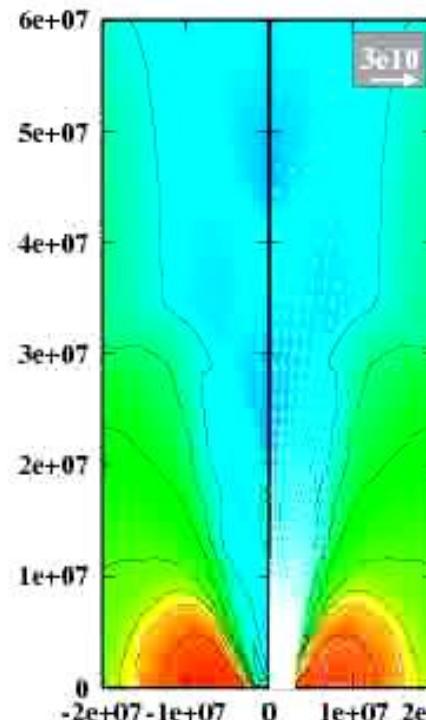
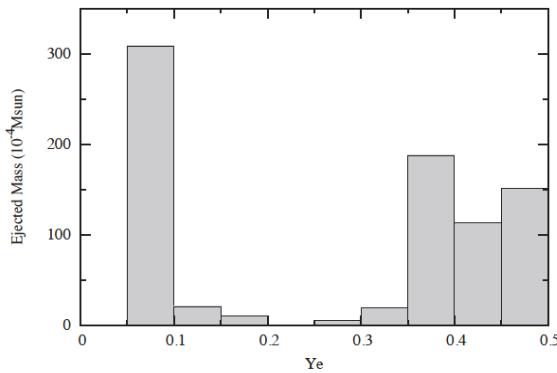
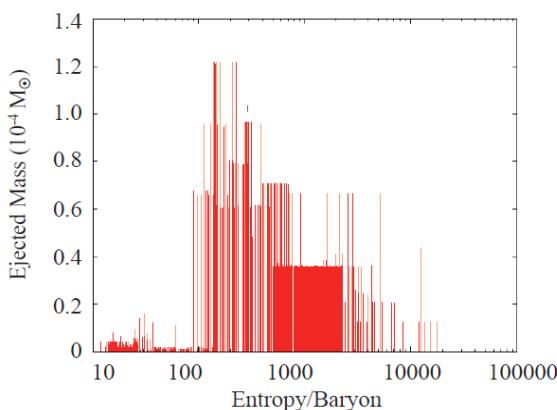
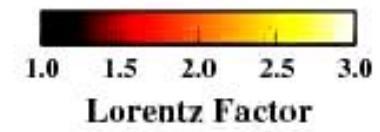
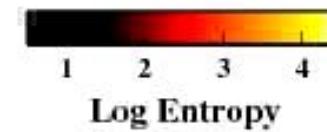
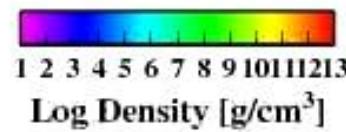
Neutrino-Pair Heating



Collapsar Model for Long Gamma-Ray Bursts

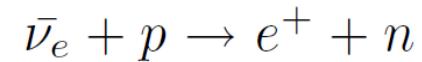
Harikae et al., ApJ 704 (2009), 354; 713 (2010) 304.

Nakamura, Kajino, Mathews, Sato & Harikae, IJMP 22 (2013), 1330022.



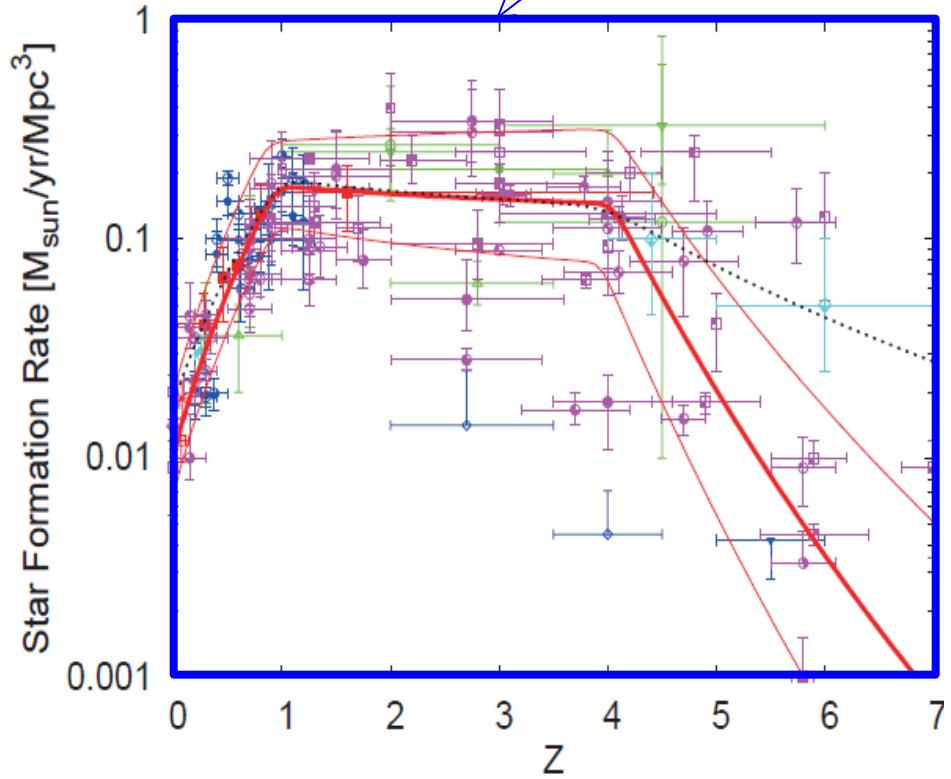
Spectrum of Relic Supernova Neutrinos (RSNs)

for Hyper-Kamiokande (Mega-ton): Water Cherenkov

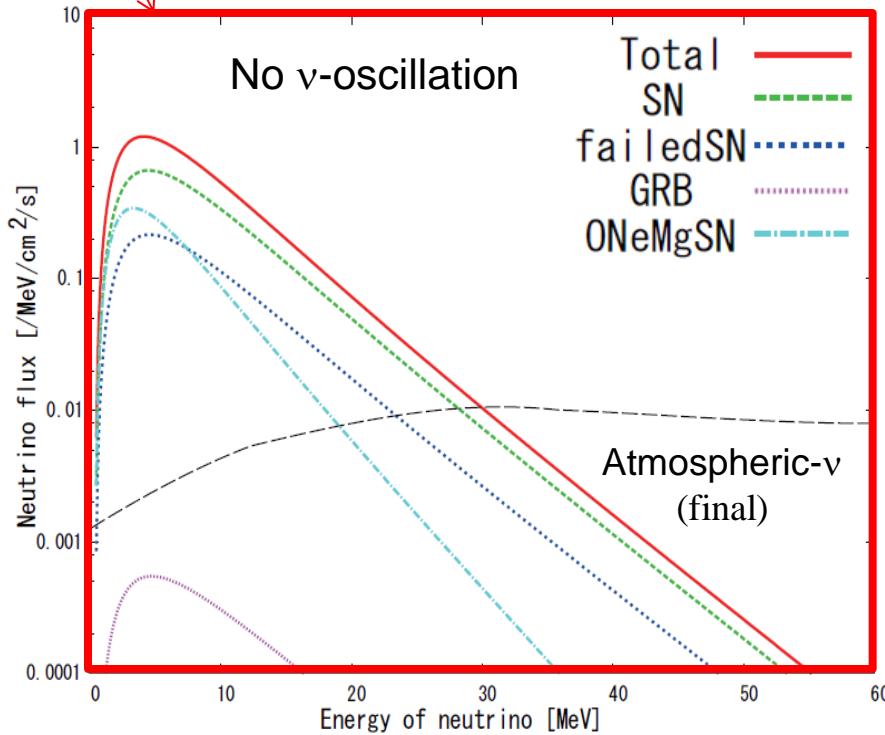


$$\frac{dN_\nu}{dE_\nu} = \frac{c}{H_0} \int_0^{z_{max}} R_{SN}(z) \frac{dN_\nu(E'_\nu)}{dE'_\nu} \times \frac{dz}{\sqrt{(\Omega_m)(1+z)^3 + \Omega_\Lambda}}$$

SN Rate x Volume



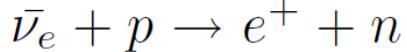
ν -spectrum at Various SNe & GRB



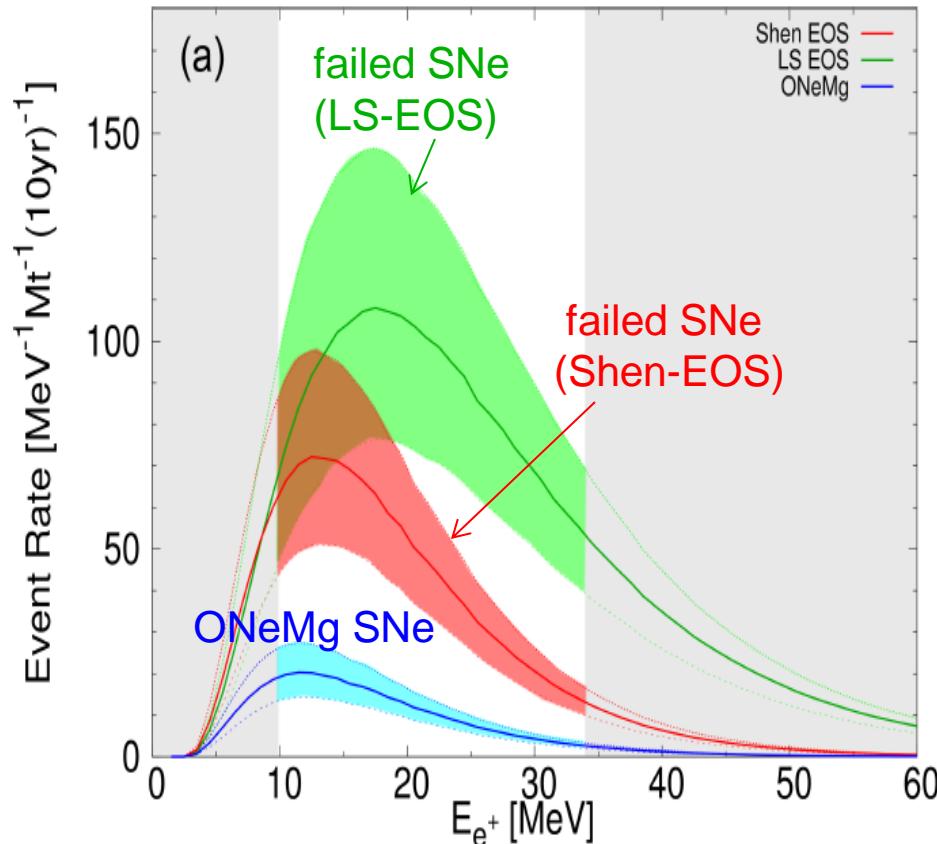
Relic Supernova Neutrinos (RSNs)

G. J. Mathews, J. Hidaka, T. Kajino, and J. Suzuki, ApJ (2014), submitted.

Hyper-Kamiokande (Mega-ton, 10y), Gd-loaded Water Cherenkov Detector



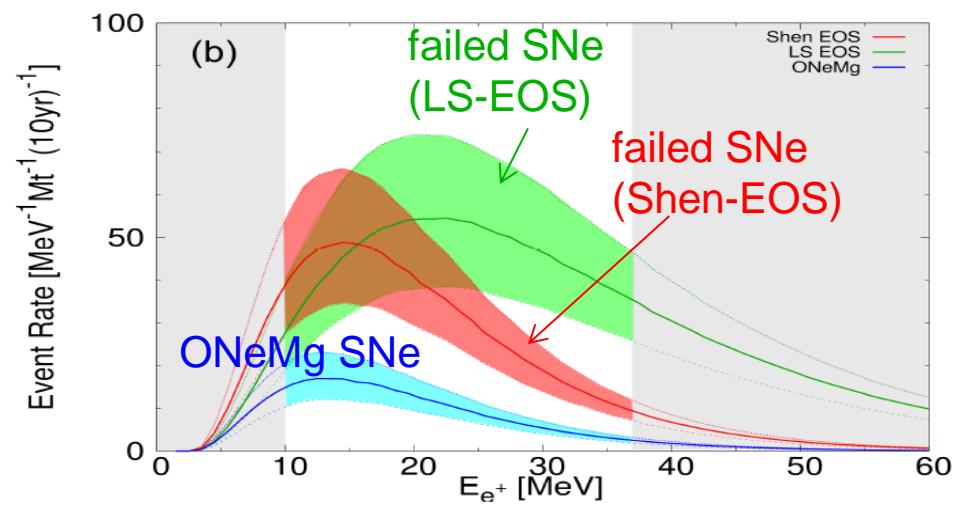
Non-Adiabatic Matter (MSW) Oscillation



Assuming Horiuchi, Beacom et al. 2011
100% missing SNe:

Standard + **MISSING**
SNe (ONeMg+CC+fSN) + GRB

Adiabatic Matter (MSW) Oscillation



SUMMARY

Total ν -mass:

Total ν -mass is constrained to be $\sum m_\nu < 0.2$ eV for the primordial magnetic field $B < 3nG$.

ν -Mass hierarchy:

Supernova ν -process could determine the mass hierarchy Δm_{13}^2 and $\sin^2 \theta_{13} \sim 0.1$ simultaneously. Inverted hierarchy is more preferred statistically.

Relic Supernova- ν :

Future observation of Relic Supernova ν 's in megaton Hyper-Kamiokande (i.e. Gd-loaded Water Cherenkov detector in 10y run) could identify the missing SN component and discriminate EoS of proto-neutron star and neutrino oscillation pattern.