

ニュートリノ質量、暗黒物質、 物質・反物質非対称性と 右巻きニュートリノ

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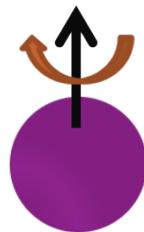
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Outline

- Phenomena beyond SM
 - 1. Neutrino mass
 - 2. Dark matter
 - 3. Matter and Anti-matter asymmetry
- Explanation by right-handed neutrinos

Mass in Standard Model

Left-handed:



Right-handed:



- Mass term in SM

electron:

$$-\mathcal{L}_{mass} = y \overline{L}_L \Phi e_R + h.c.$$

$$L_L \equiv \begin{pmatrix} \nu_L \\ e_L \end{pmatrix}$$

$$\langle \Phi \rangle = \begin{pmatrix} 0 \\ v \end{pmatrix}$$

➡ $m(\overline{e}_L e_R + \overline{e}_R e_L)$

neutrino:

$$-\mathcal{L}_{mass} = y \overline{L}_L \tilde{\Phi} \underline{\nu}_R + h.c.$$



ν_R is not in SM, so neutrino mass = 0

Motivation

ν mass by experiment:
 ('13 Cappozi et. al.)

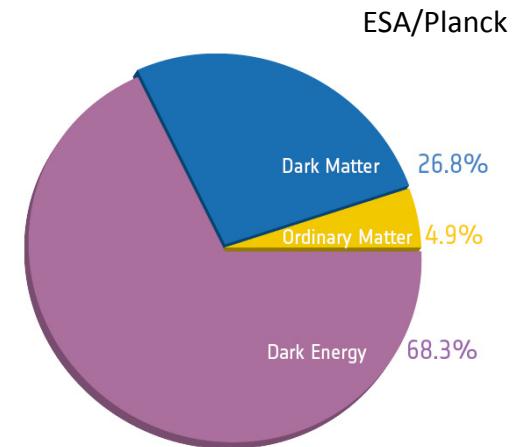
$$m_2^2 - m_1^2 = (0.00868\text{eV})^2$$

$$m_3^2 - (m_2^2 + m_1^2)/2 = (0.0494\text{eV})^2$$

ν_R can explain neutrino mass

Problem in cosmology:

- Dark matter
- Matter-anti-matter asymmetry



Q. All solved by ν_R ?

A. YES (Conclusion)

Neutrino mass

General mass term

$$-\mathcal{L}_{mass} = \frac{1}{2} \begin{pmatrix} \overline{\nu_L} & \overline{\nu_R^c} \end{pmatrix} \begin{pmatrix} M_L & yv \\ y^T v & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} + \text{h.c.}$$

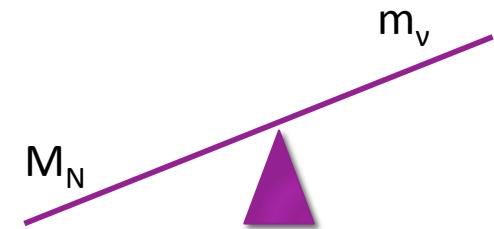
Block diagonalization ($M_L, yv \ll M_R$)

$$\begin{pmatrix} m_\nu & 0 \\ 0 & M_N \end{pmatrix} \quad \begin{aligned} m_\nu &= M_L - v^2 y M_R^{-1} y^T & (0) & : \text{light} \\ M_N &\simeq M_R & & : \text{heavy} \end{aligned}$$

m_ν and M_N are diagonalized by mixing matrix U, V

$$m_\nu^{\text{diag}} = U^\dagger m_\nu U^*$$

$$M_N^{\text{diag}} = V^\dagger M_R V^*$$



Interaction of ν_R

	Weak interaction	Electromagnetic interaction
e_L	○	○
e_R	×	○
ν_L	○	×
ν_R	×	×

- ν_R interact weakly by mixing with $\nu_L \rightarrow$ Dark matter

Life time: $\tau \sim 10^7 \tau_U$ (for $M_1 = 7\text{keV}$, $\Theta_1^2 = 10^{-10}$)

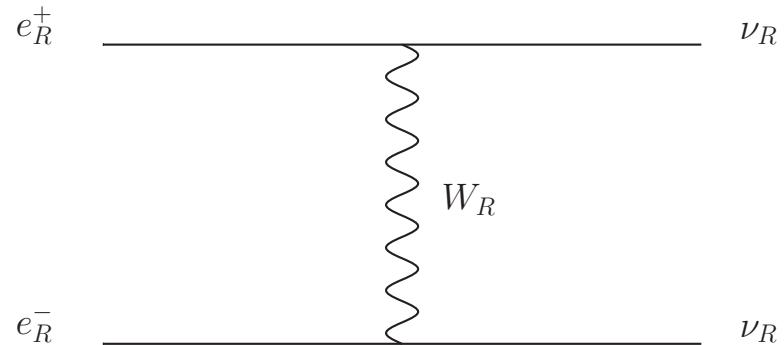
Production of ν_{R1} DM

Assumption 1: ν_R were in thermal equilibrium

-> Introduce new gauge interaction

(e.g., Left-right model, GUT)

-> DM is thermal relic



Advantage of thermal relic:

1. No need for assuming initial abundance of ν_R
2. Simple calculation of abundance today
3. Colder than other mechanism ('94 Dodelson, Widrow)
 ->weaker constraint from Ly- α

Production of ν_{R1} DM

- Problem of thermal relic: Over production

$$\frac{n_{N_1}}{s} = \frac{1}{g_{*f}} \frac{135\zeta(3)}{4\pi^4} \quad : \text{constant}$$

➡ $\Omega_{N_1} \gtrsim 10\Omega_{DM}$
 $\gg \Omega_{DM} \sim 0.3$

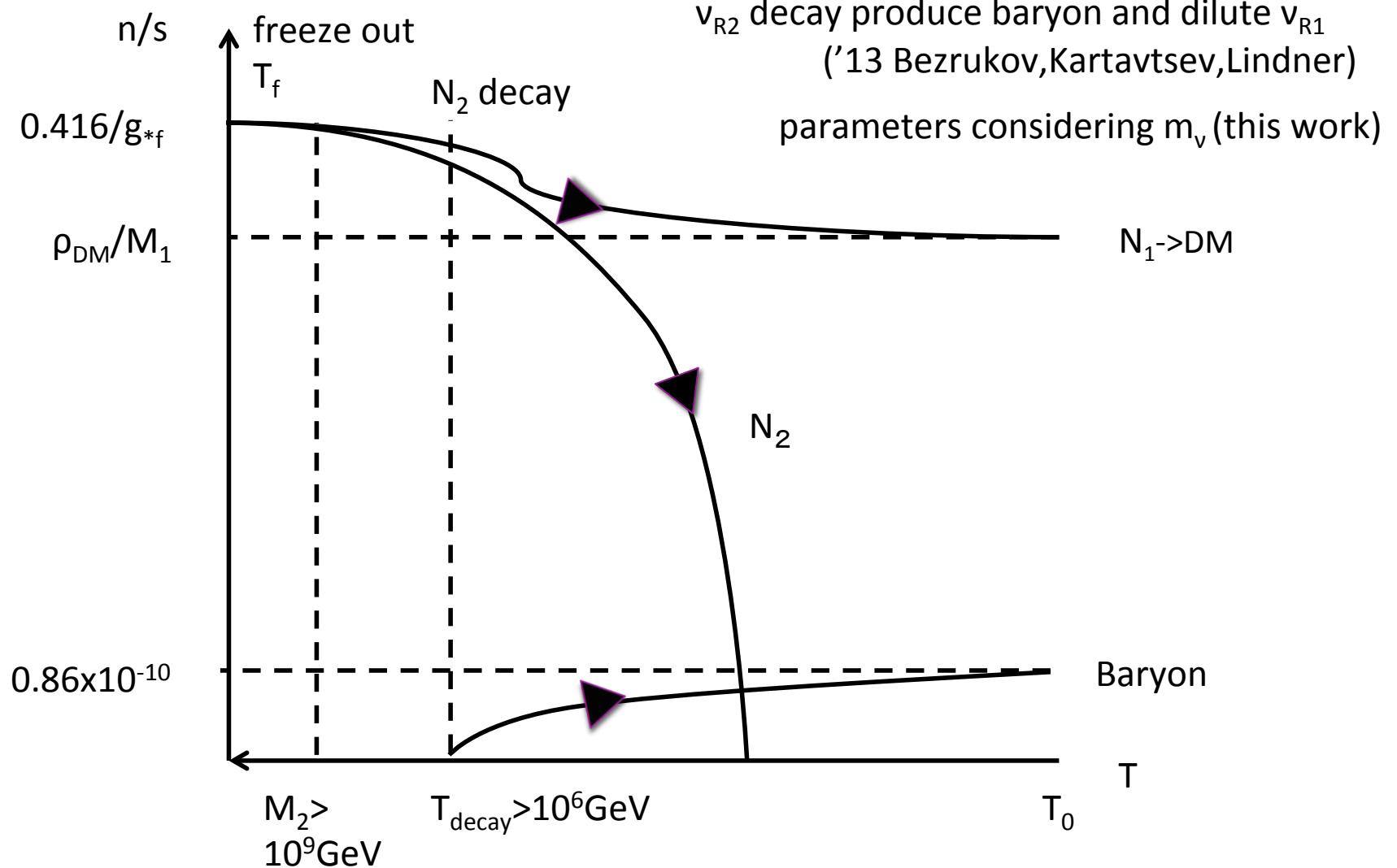
- Solution: Diluted by Entropy production $s \rightarrow \frac{S_{after}}{S_{before}} s$

Assumption 2: ν_{R2} decay produce entropy

$$\frac{S_{after}}{S_{before}} \sim 10$$

Thermal history

v_{R1}, v_{R2} was in equilibrium



Entropy production

Ratio of entropy
before and after ν_{R2} decay:

$$\frac{N_I}{\tilde{y}_{\alpha I} \equiv (yV^*)_{\alpha I}} \xrightarrow{\Phi} L_{L\alpha}$$

$$10 \sim \frac{S_{after}}{S_{before}} \propto \frac{1}{\sqrt{\Gamma_{N_2}}}$$

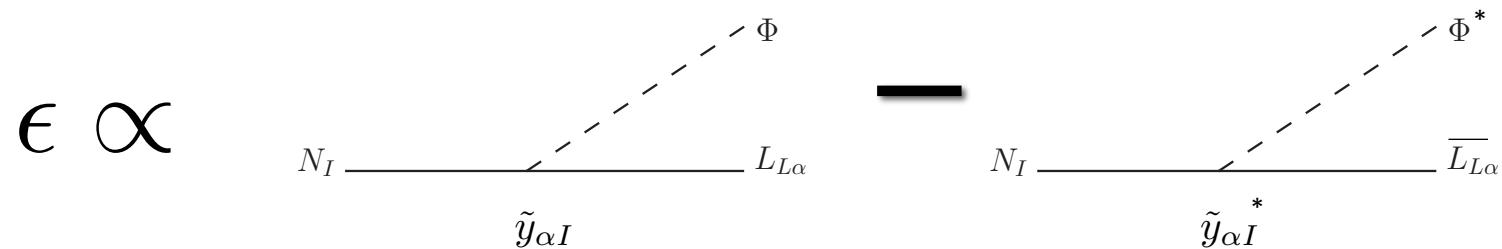
$$\Gamma_{N_2} = \frac{(\tilde{y}^\dagger \tilde{y})_{22}}{8\pi} M_2$$



$$(\tilde{y}^\dagger \tilde{y})_{22} = 1.1 \times 10^{-14} \left(\frac{\text{keV}}{M_1} \right)^2 \left(\frac{M_2}{10^9 \text{GeV}} \right) \quad (1)$$

Matter-antimatter asymmetry

N_2 decays into lepton and anti-lepton



Ratio: 1,000,000,000 : 1,000,000,001

Baryon number:

$$\frac{n_B}{s} = -1.4 \times 10^{-4} \epsilon \left(\frac{\text{keV}}{M_1} \right)$$

By 1-loop calculation of ϵ ,

$$\text{Im}[(\tilde{y}^\dagger \tilde{y})_{32}^2] = 1.66 \times 10^{-19} \frac{\text{keV}}{M_1} \frac{M_2}{10^9 \text{GeV}} \frac{1}{g(M_3^2/M_2^2)}$$

(2)

Parameter

$$-\mathcal{L}_{mass} = \frac{1}{2} \begin{pmatrix} \overline{\nu_L} & \overline{\nu_R^c} \end{pmatrix} \begin{pmatrix} M_L & yv \\ y^T v & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} + \text{h.c.}$$

Condition:

1. Neutrino mass

$$m_\nu = M_L - v^2 y M_R^{-1} y^T \quad (0)$$

2. Dark matter

$$(\tilde{y}^\dagger \tilde{y})_{22} = 1.1 \times 10^{-14} \left(\frac{\text{keV}}{M_1} \right)^2 \left(\frac{M_2}{10^9 \text{GeV}} \right) \quad (1)$$

3. Matter asymmetry

$$\text{Im}[(\tilde{y}^\dagger \tilde{y})_{32}^2] = 1.66 \times 10^{-19} \frac{\text{keV}}{M_1} \frac{M_2}{10^9 \text{GeV}} \frac{1}{g(M_3^2/M_2^2)} \quad (2)$$

Result

Masses of ν_R : $M_1 \sim \text{keV}, M_2 \gtrsim 10^9 \text{GeV}$

Yukawa coupling: $\tilde{y}v = iV_\nu(X_\nu^{\text{diag}})^{1/2}R(M_R^{\text{diag}})^{1/2}$

$$X_\nu \equiv m_\nu - M_L, \quad X_\nu^{\text{diag}} = V_\nu^\dagger X_\nu V_\nu^* \quad R \sim \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

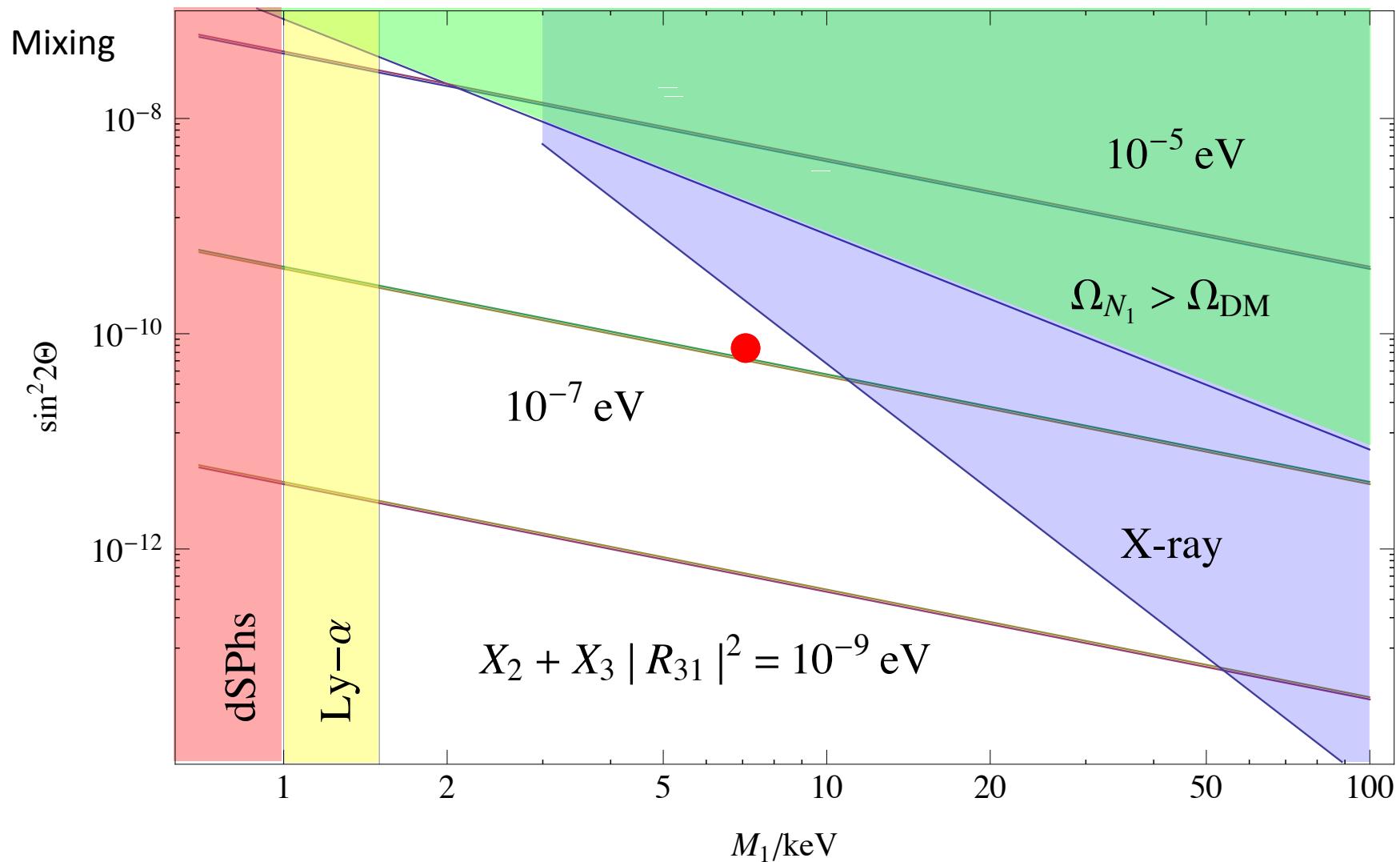
Mixing of N_1

$$\Theta_1^2 = \frac{(\tilde{y}y)_{11}v^2}{M_1^2} = \frac{X_2 + X_3|R_{31}|^2}{M_1}$$



Constraints on ν_{R1} DM

Detection?@7keV
(arXiv: 1402.2301, 1402.4119)



Summary

Model:

- neutrino mass explained by heavy ν_R
- ν_{R1} DM produced as thermal relic (over production)
- Decay of ν_{R2} dilute ν_{R1} and produce matter-anti-matter asymmetry

Conclusion:

- m_ν , DM, asymmetry of matter can be explained simultaneously

ニュートリノの混合

弱い相互作用の固有状態

$$\nu_{L\alpha} = U_{\alpha i} \nu'_{Li} + \theta_{\alpha I} \nu'^c_{RI}$$

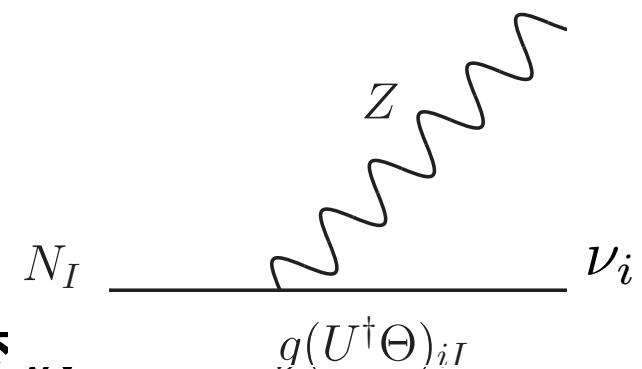
$$\theta_{\alpha I} \equiv (m_D M^{-1})_{\alpha I}$$

重いニュートリノは、混合により弱い相互作用をする

強さ: $\sim g\theta_I$

$$\theta_I^2 \equiv \sum_{\alpha=e,\mu,\tau} |\theta_{\alpha I}|^2$$

重いニュートリノの反応性を決定する変数



• $g(x)$

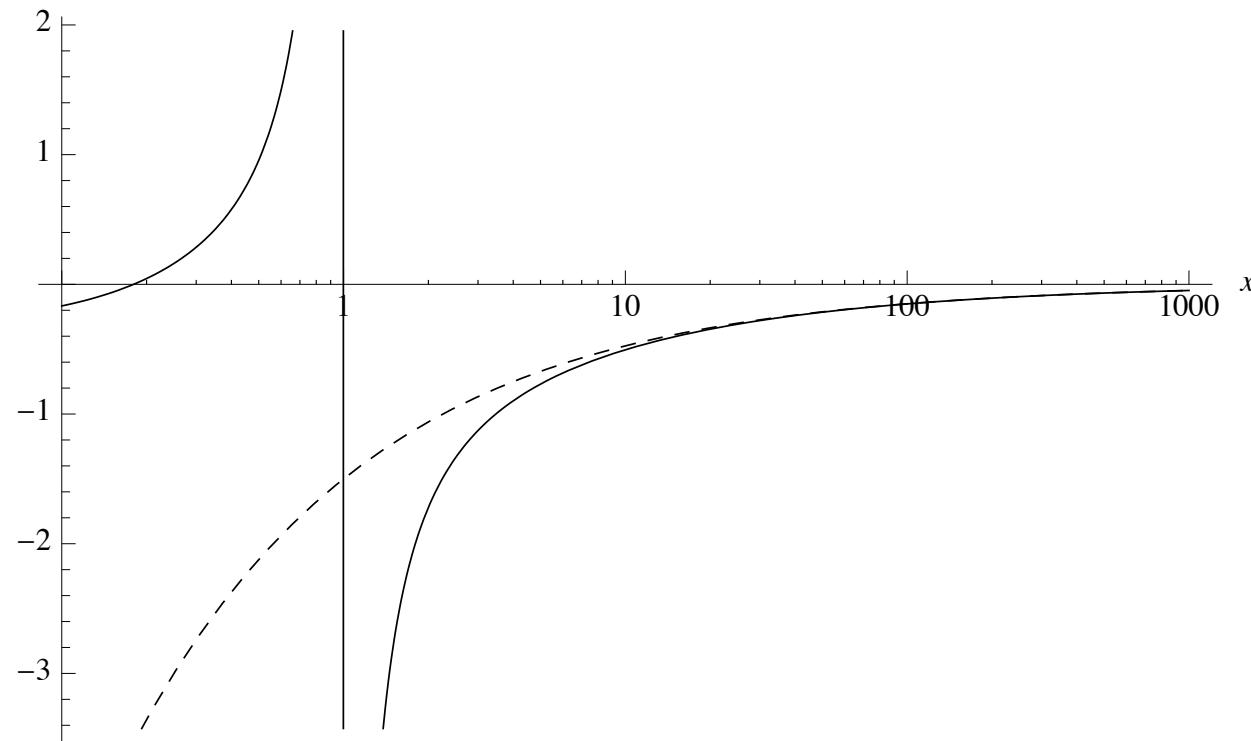


図 3.2: $g(x)$ (実線) およびその漸近式 $-\frac{3}{2\sqrt{x}}$ (点線)。

$$g(M_3^2/M_2^2) \simeq -\frac{3}{2} \frac{M_2}{M_3} \quad (M_3/M_2 \gtrsim 3)$$

Tremaine-Gunn境界

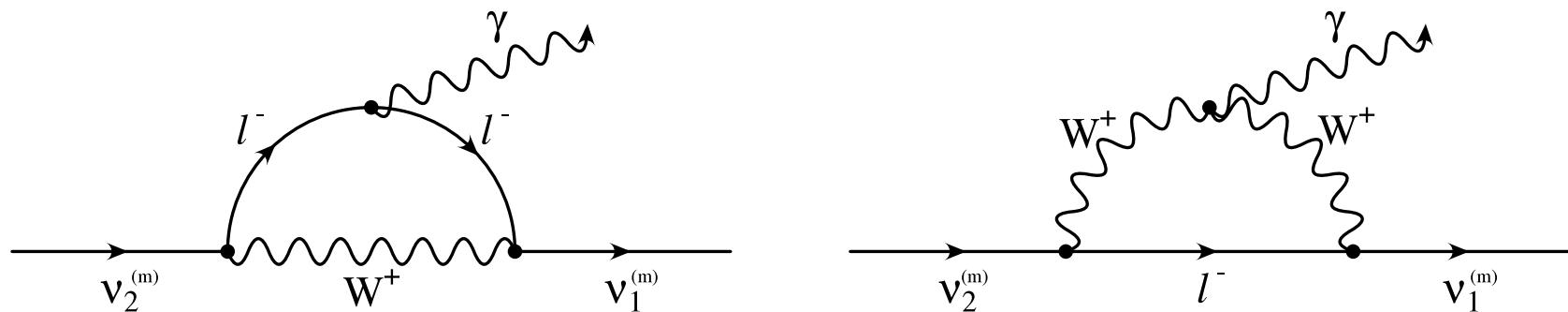
$$\begin{aligned}
 Q &\equiv \frac{\rho}{\langle v^2/3 \rangle^{3/2}} \\
 &= 3^{3/2} m_{DM}^4 \frac{n}{\underline{\langle p^2 \rangle^{3/2}}} \\
 &< \max f(p) = \frac{g_{\text{DM}}}{(2\pi)^3}
 \end{aligned}$$

dSph(Coma Berenices
等)の観測:

$$\begin{aligned}
 q &= 5 \times 10^{-3} \sim 2 \times 10^{-2}, \\
 Q &= q \frac{M/\text{pc}^3}{(\text{km/s})^3}
 \end{aligned}$$

$$m > 1.0 \text{ keV} \left(\frac{q}{5 \times 10^{-3}} \right)^{1/4}$$

X-ray bound



$$\begin{aligned}\Gamma_{\nu_s \rightarrow \gamma \nu_a} &= \frac{9}{256\pi^4} \alpha_{\text{EM}} G_F^2 \sin^2 \theta m_s^5 \\ &= \frac{1}{1.8 \times 10^{21} \text{ s}} \sin^2 \theta \left(\frac{m_s}{\text{keV}}\right)^5\end{aligned}$$

バリオン数生成

- CP: ν_R の湯川相互作用
- 熱平衡からのずれ：
熱浴の外にいる ν_R の崩壊

作られるレプトン数：

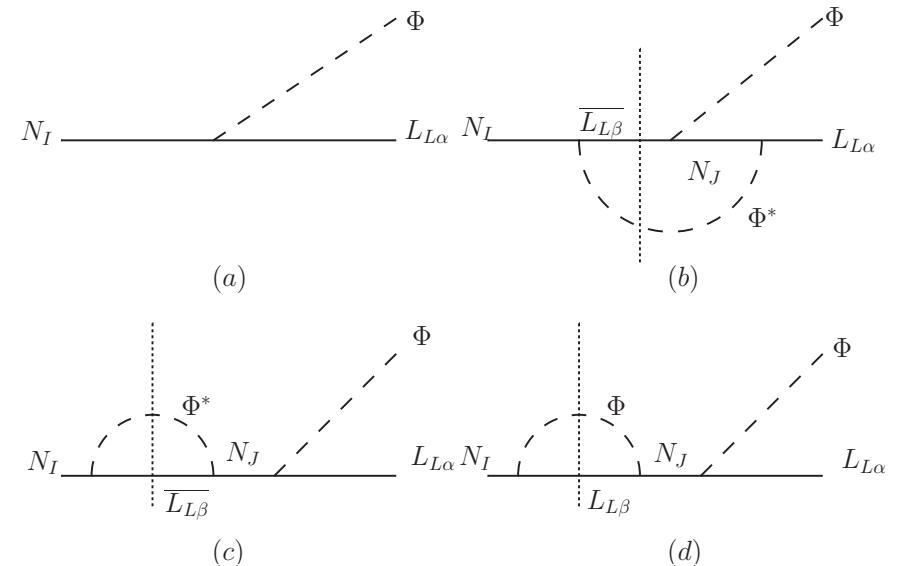
$$Y_{L_\alpha}|_{\text{崩壊後}} \propto \epsilon_{I\alpha} Y_I|_{\text{崩壊前}}$$

CPの破れ：

$$\epsilon_{I\alpha} \equiv \frac{\langle \Gamma(N_I \rightarrow L_\alpha \Phi) \rangle - \langle \Gamma(N_I \rightarrow \bar{L}_\alpha \Phi^*) \rangle}{\sum_\alpha \{ \langle \Gamma(N_I \rightarrow L_\alpha \Phi) \rangle + \langle \Gamma(N_I \rightarrow \bar{L}_\alpha \Phi^*) \rangle \}}$$

量子異常によりレプトン数は₈₈⁷⁹バリオン数に変換

$$Y_B = \frac{79}{88} Y_L$$



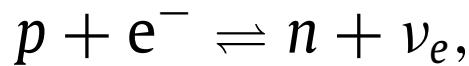
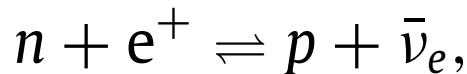
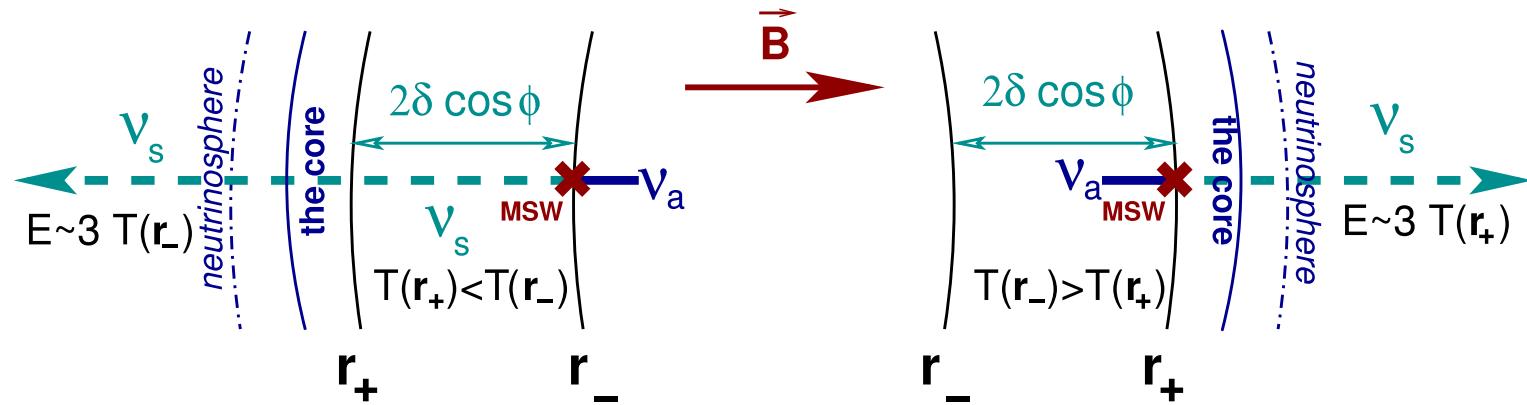
非等方性

$$E \sim 3 \times 10^{53} \text{ erg}, \quad \begin{aligned} & (09 \text{ Kusenko}) \\ & 1\text{erg}=10^{-7}\text{J} \\ & M_{\text{sun}}=2.0*10^{30}\text{kg} \end{aligned}$$

$$p_{\nu, \text{total}} \sim 1 \times 10^{43} \text{ g cm/s.}$$

$$\begin{aligned} p_* = (1.4M_{\odot})v &\approx 3 \times 10^{41} \left(\frac{v}{1000 \text{ km/s}} \right) \text{ g cm/s} \\ &\approx 0.03 \left(\frac{v}{1000 \text{ km/s}} \right) p_{\nu, \text{total}}. \end{aligned}$$

1%~5%くらいの非等方性で説明できる



$$\sigma(\uparrow e^-, \uparrow \nu) \neq \sigma(\uparrow e^-, \downarrow \nu)$$

$$V(\nu_s) = 0,$$

$$V(\nu_e) = -V(\bar{\nu}_e) = V_0(3Y_e - 1 + 4Y_{\nu_e}),$$

$$V(\nu_{\mu,\tau}) = -V(\bar{\nu}_{\mu,\tau}) = V_0(Y_e - 1 + 2Y_{\nu_e}) + \frac{eG_F}{\sqrt{2}} \left(\frac{3N_e}{\pi^4} \right)^{1/3} \frac{\vec{k} \cdot \vec{B}}{|\vec{k}|},$$

$$\frac{\Delta k_s}{k} = \frac{1}{6} \frac{\Delta k_s}{k_s} = \frac{4e\sqrt{2}}{3\pi^2} \frac{\mu_e \mu_n^{1/2}}{m_n^{3/2} T^2} B$$

$$= 0.01 \left(\frac{\mu_e}{100 \text{ MeV}} \right) \left(\frac{\mu_n}{80 \text{ MeV}} \right)^{1/2} \left(\frac{20 \text{ MeV}}{T} \right)^2 \left(\frac{B}{3 \times 10^{16} \text{ G}} \right).$$

