

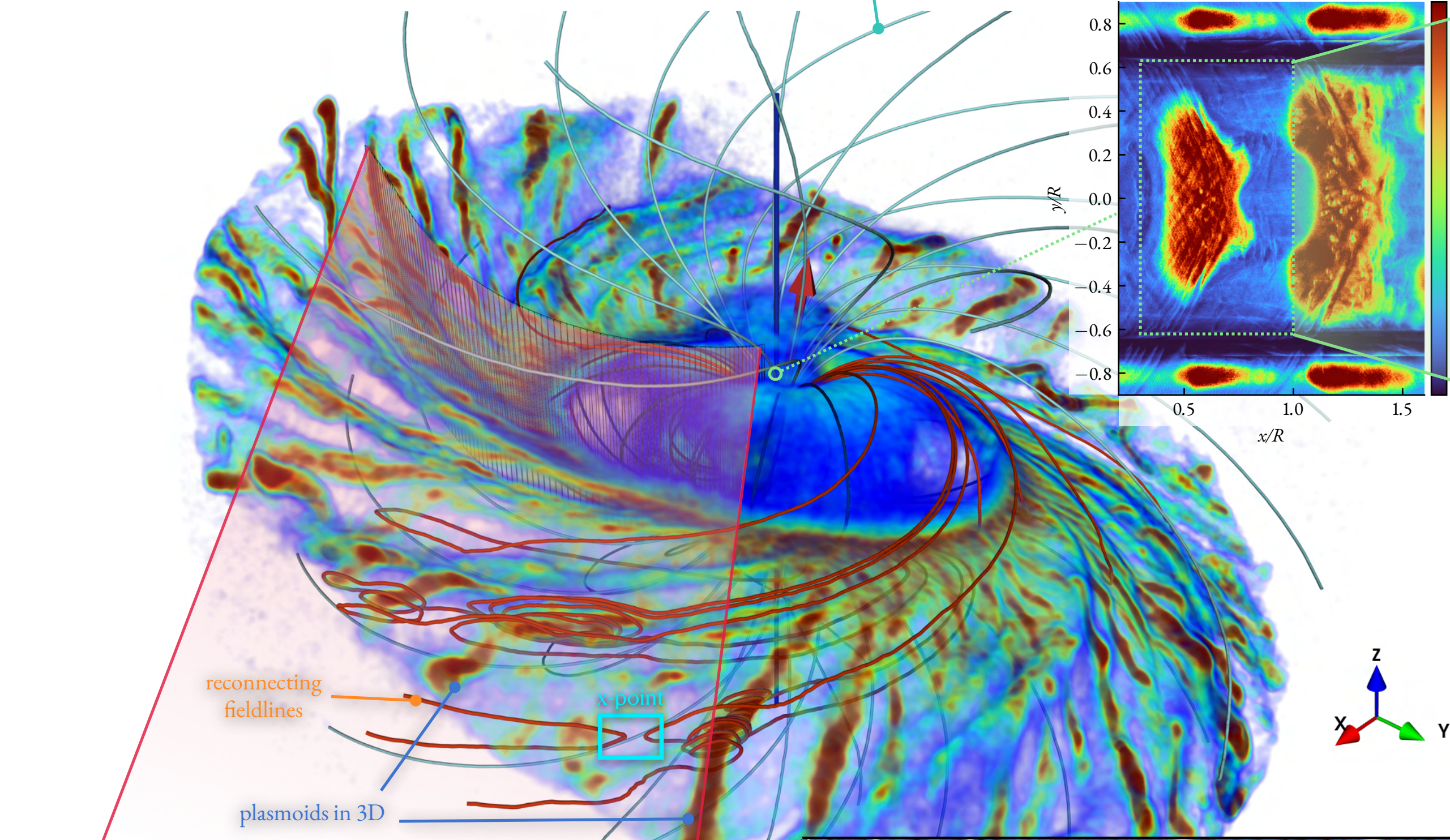
Pulsar magnetospheres and their radiation

Sasha Philippov (Maryland)

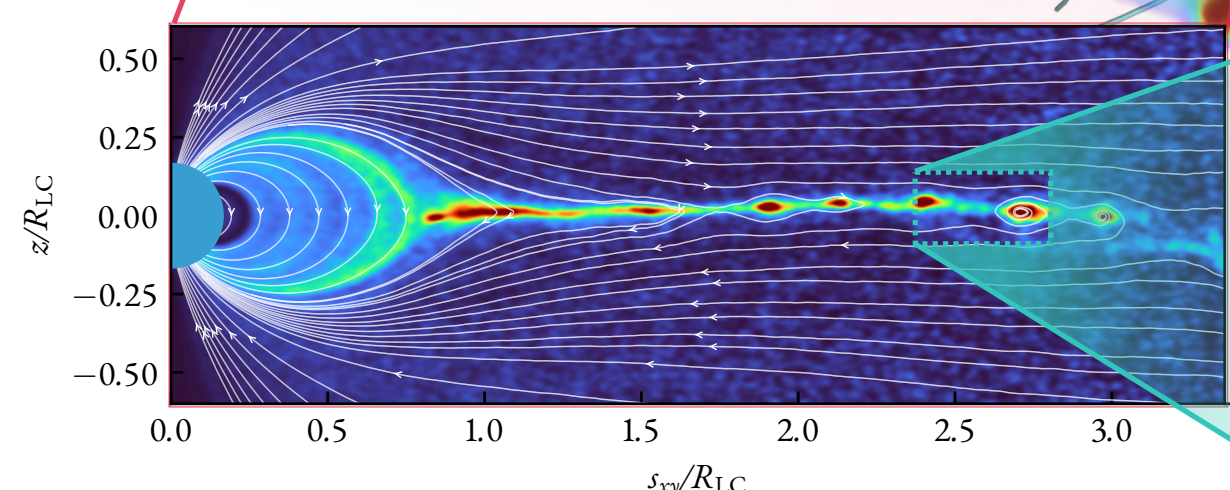
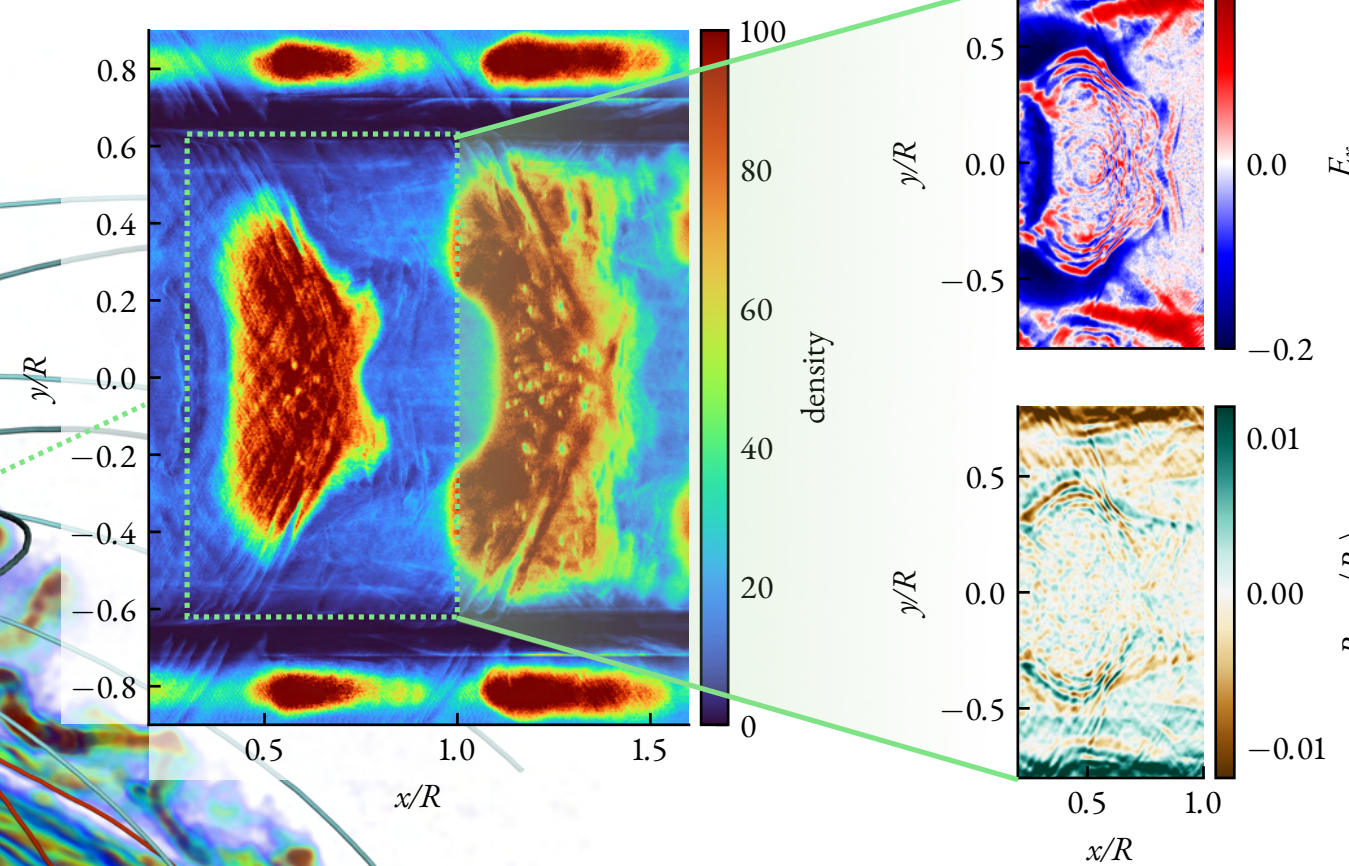
with:

- Benoit Cerutti (*Grenoble*)
- Sasha Chernoglazov (*Maryland*)
- Sam Gralla (*Arizona*)
- Hayk Hakobyan (*Columbia*)
- Anatoly Spitkovsky (*Princeton*)
- Andrey Timokhin (*Zielona Gora*)
- Libby Tolman (*IAS, Flatiron*)
- Dmitri Uzdensky (*Colorado*)

a) 3D PIC simulation of the global pulsar magnetosphere

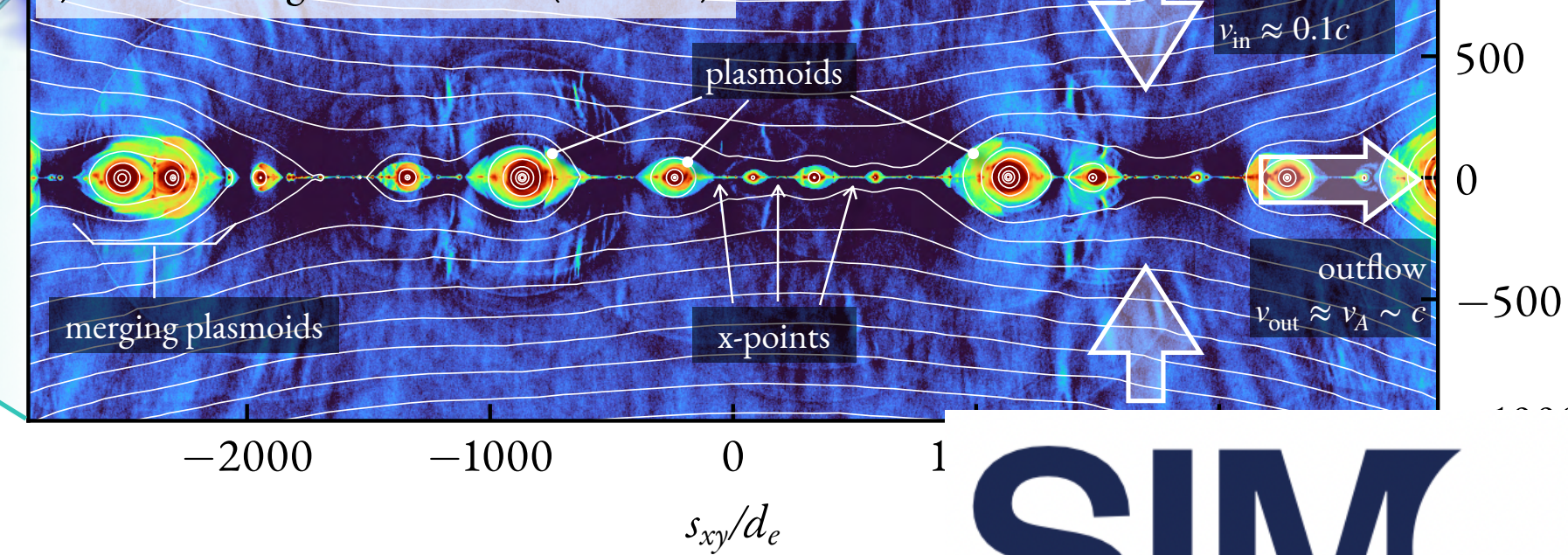


d) intermittent polar cap discharge (2D PIC)

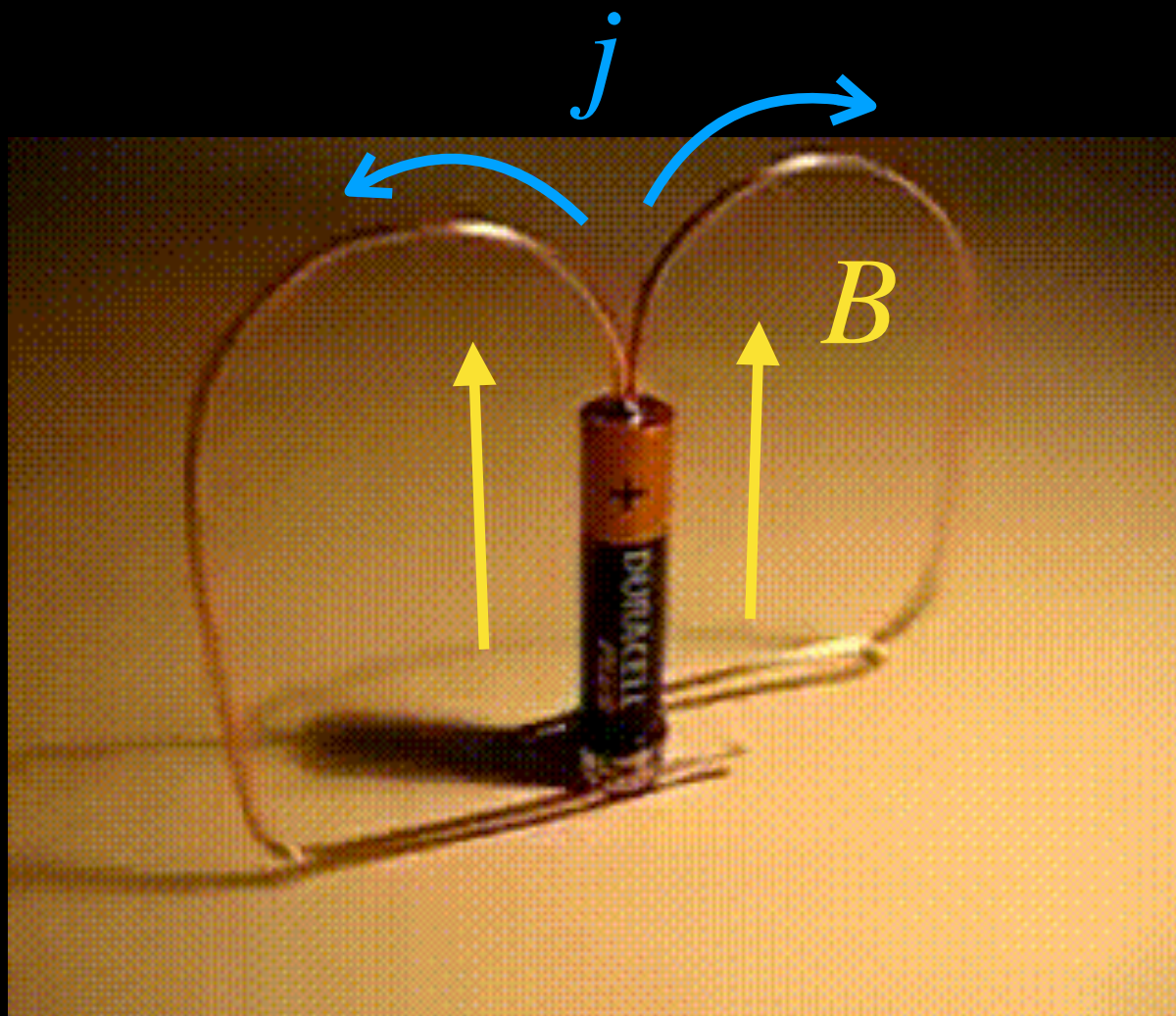
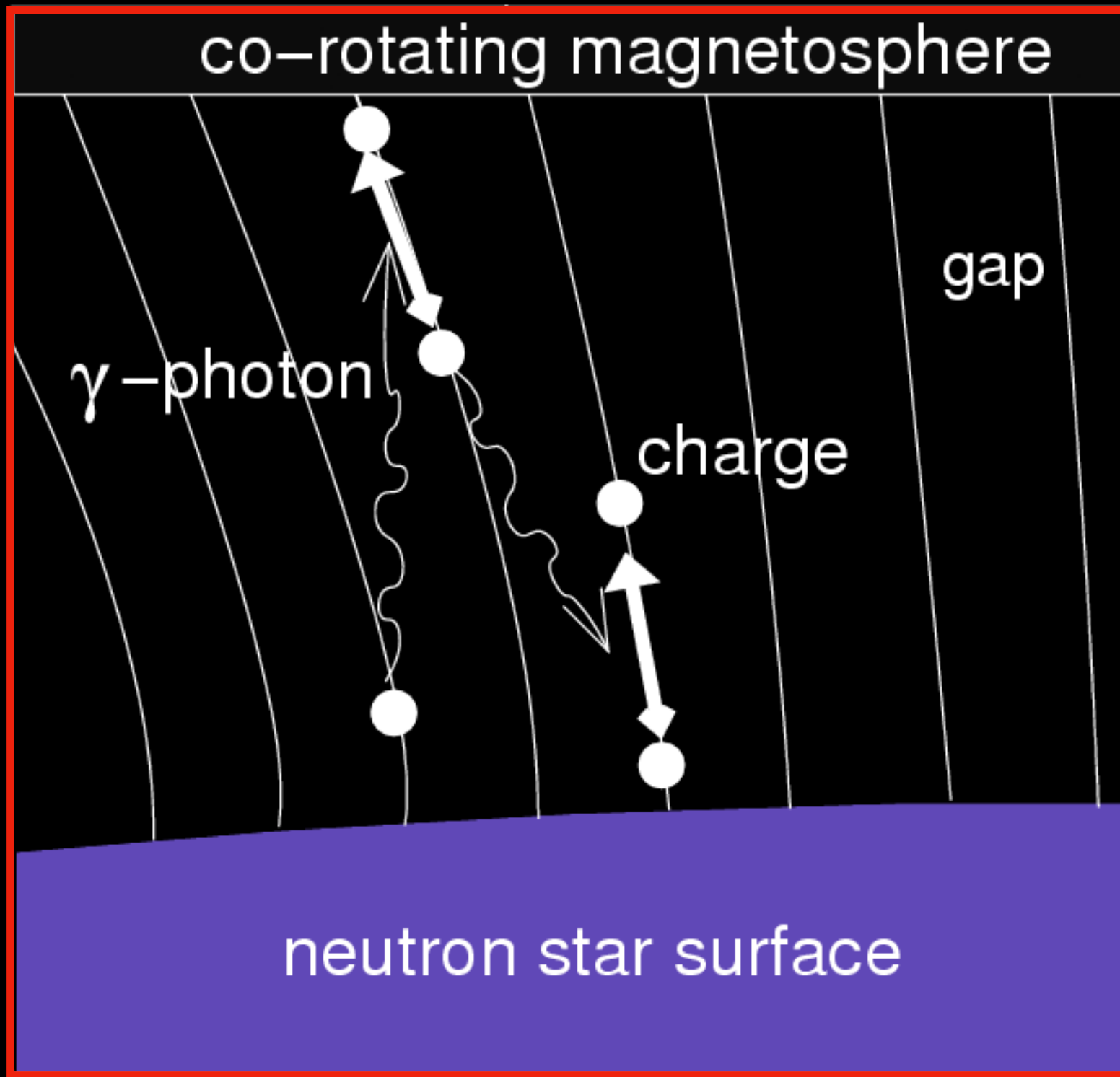


b) magnetic reconnection in the plasmoid unstable current sheet (slice from global 3D PIC)

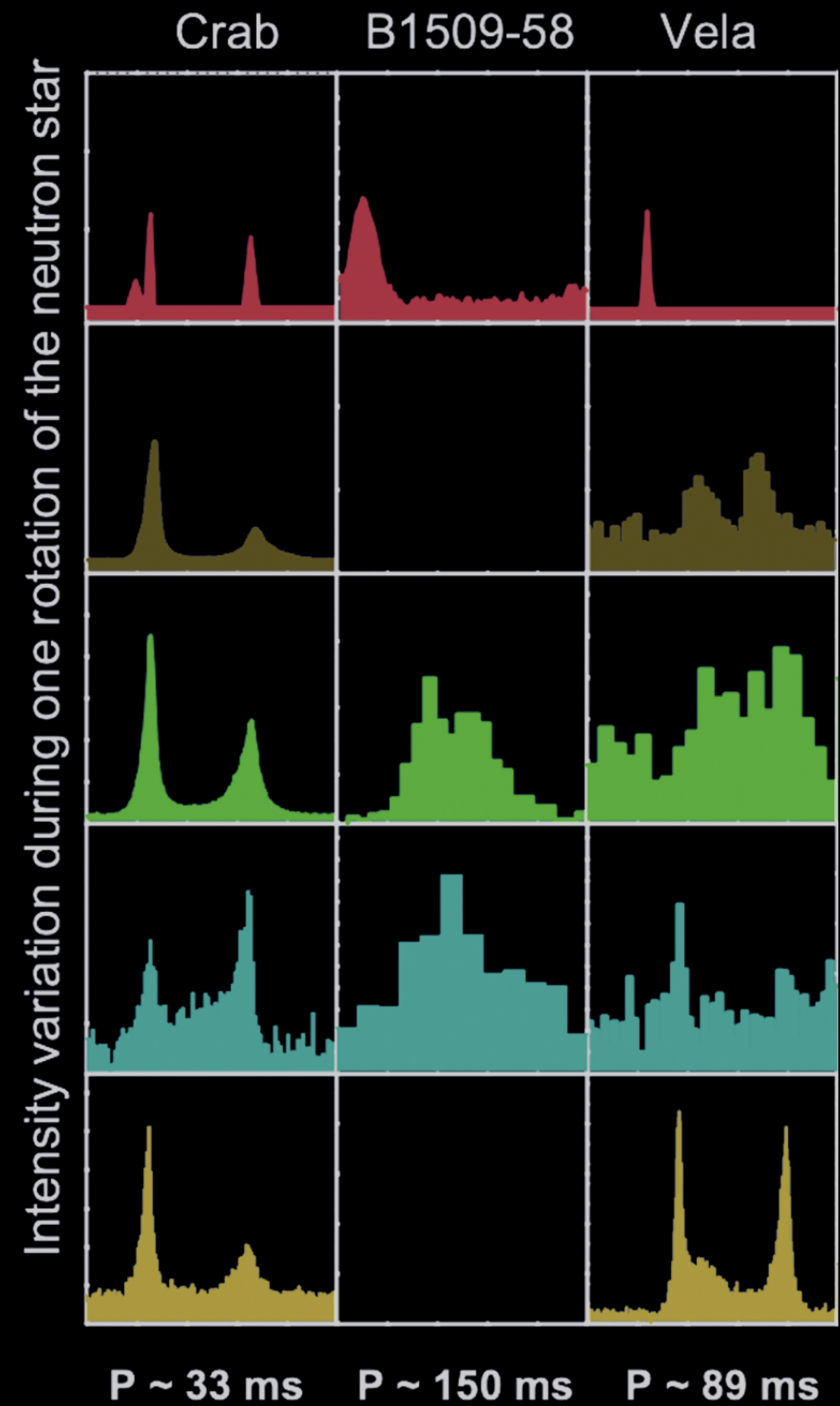
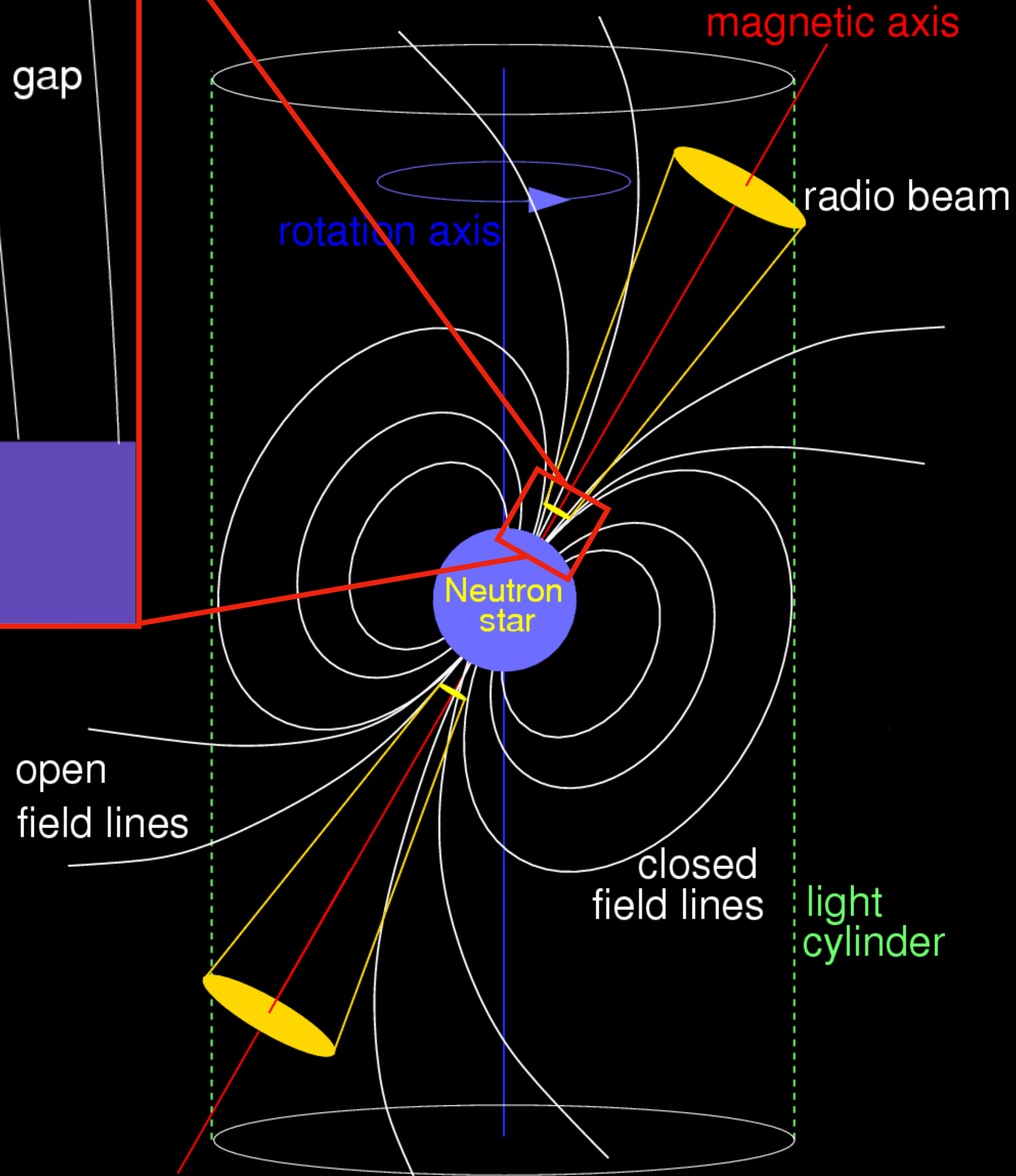
c) reconnecting current sheet (2D PIC)



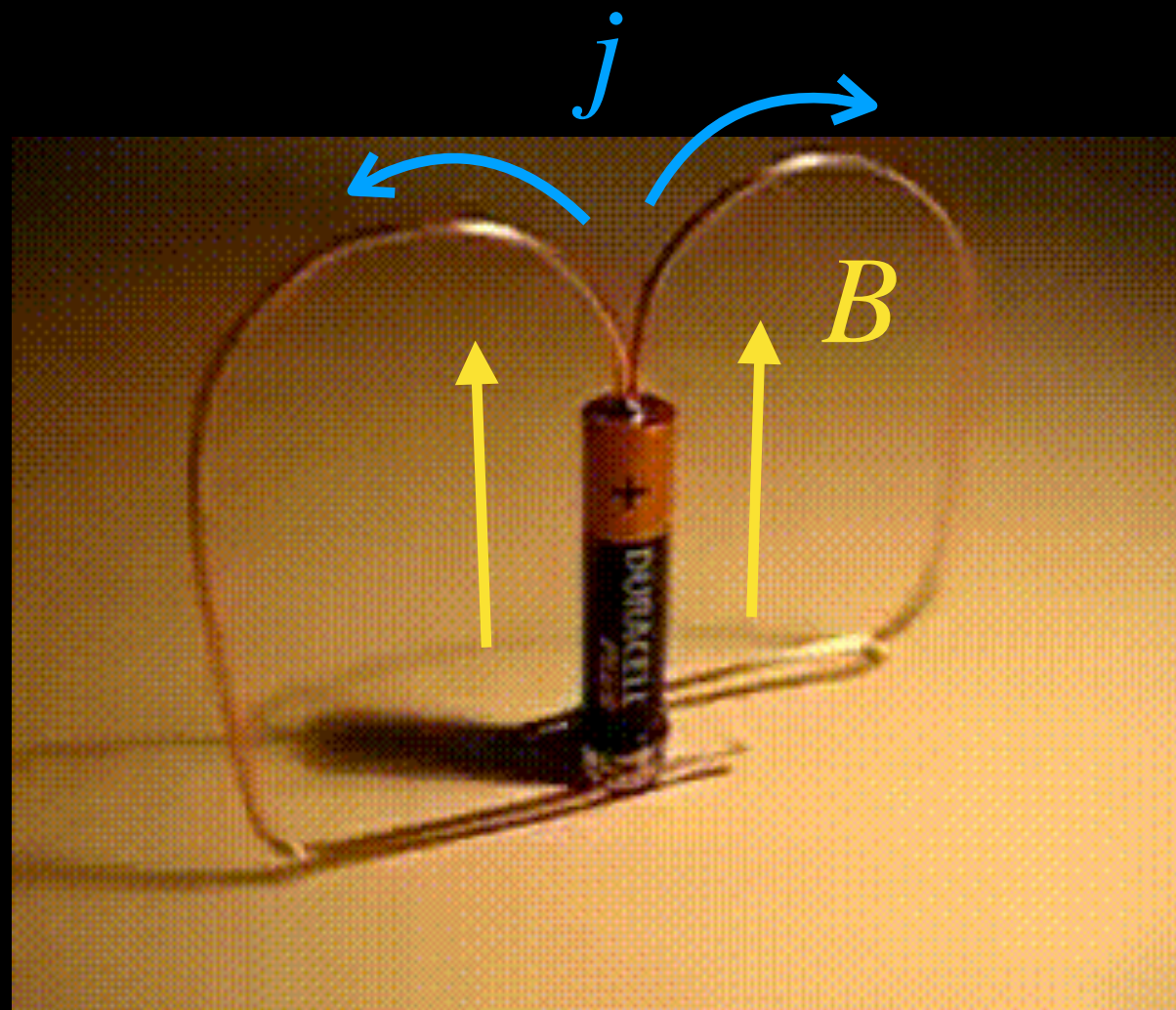
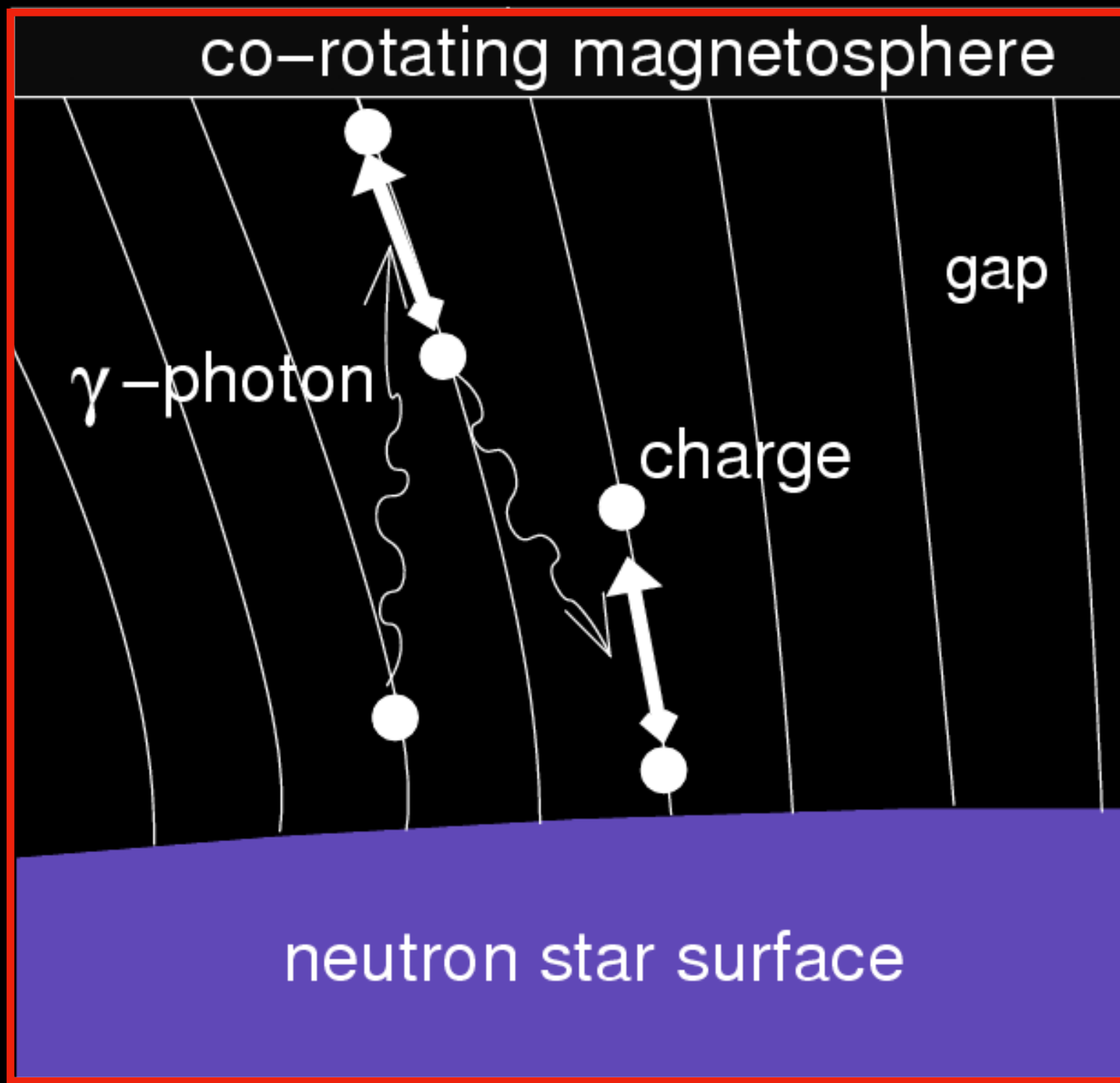
What is a pulsar?



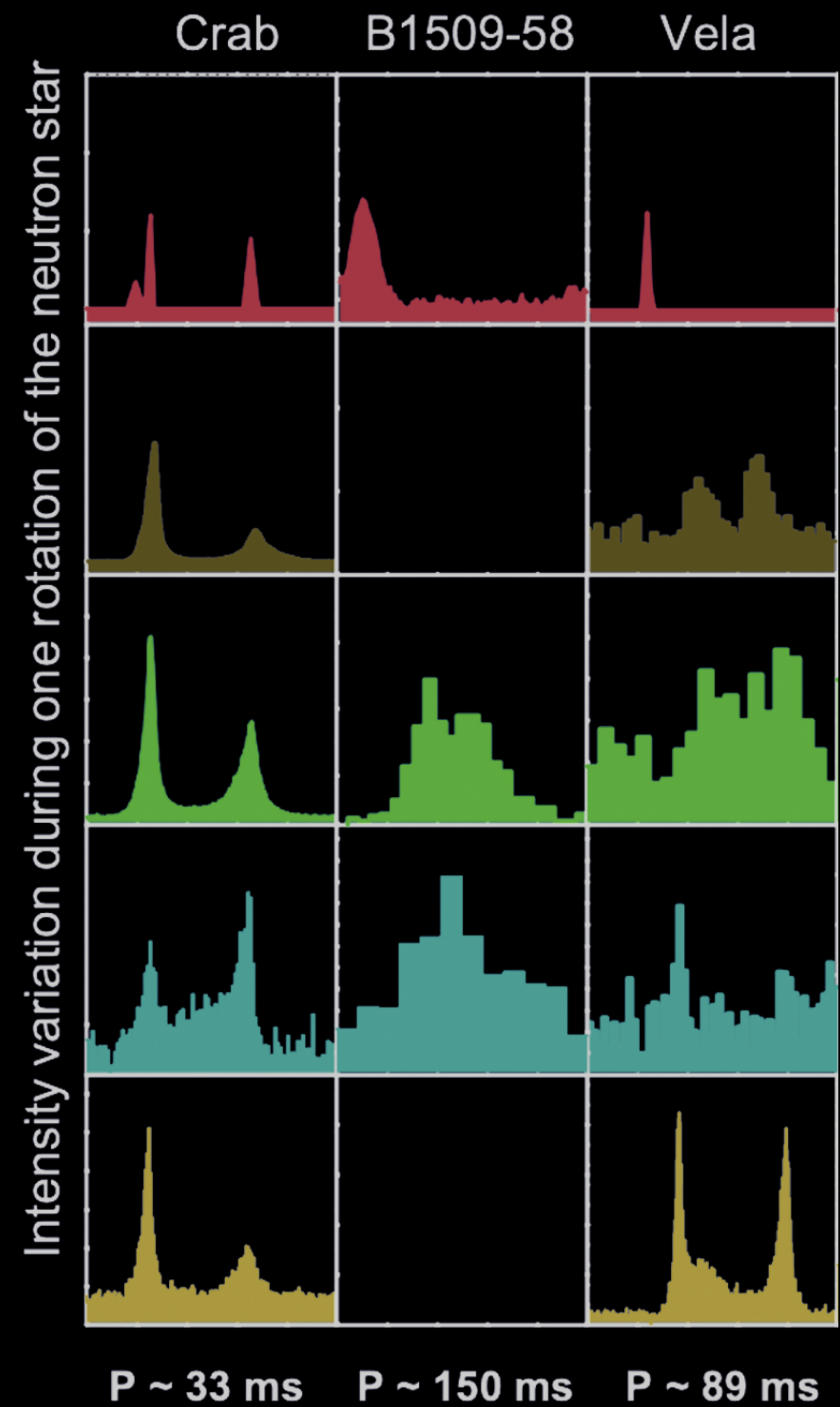
Unipolar induction



What is a pulsar?



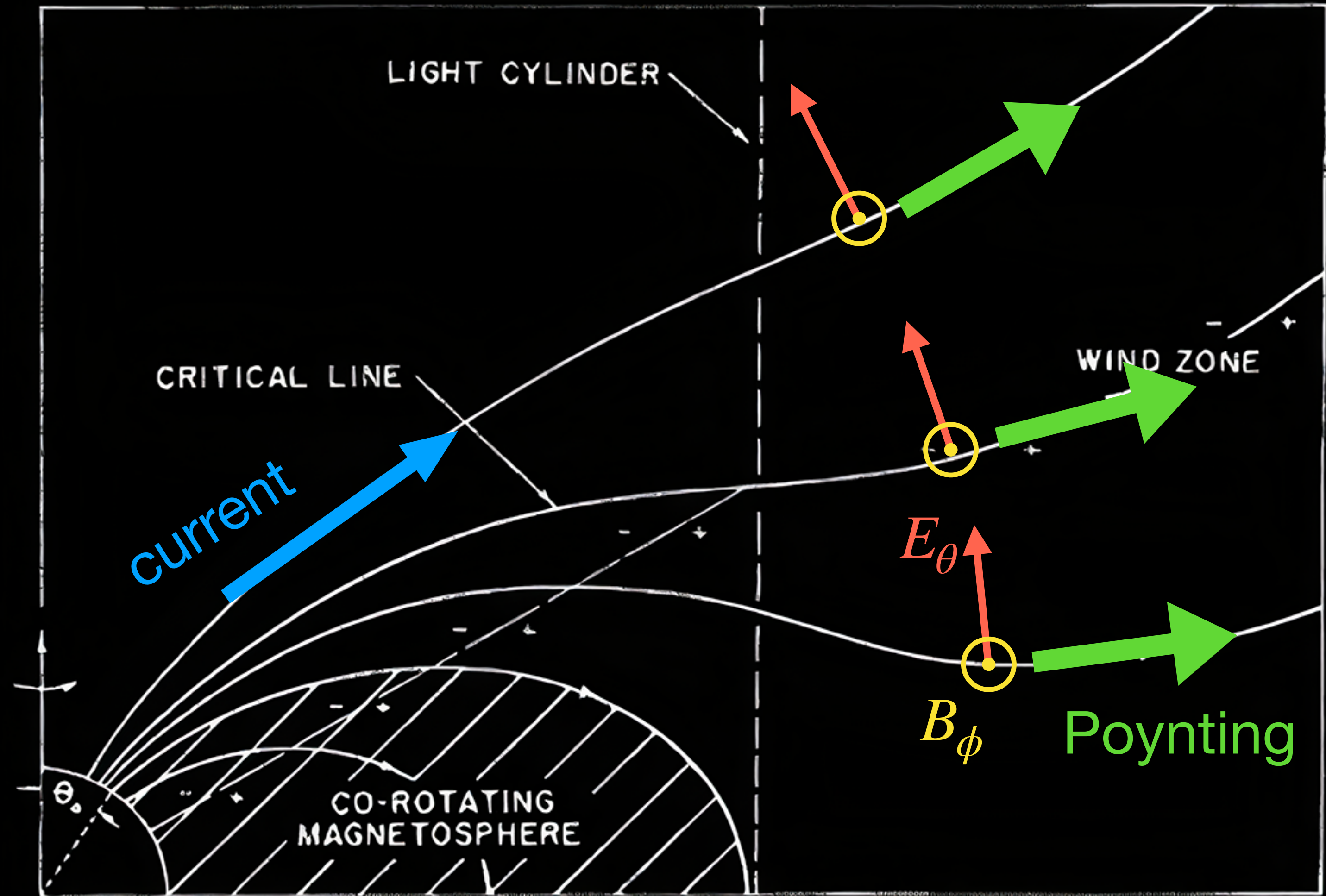
Unipolar induction



THEORETICAL CARTOON: GJ MODEL

$$\sigma \equiv \frac{B^2/4\pi}{\rho_{\pm}c^2} \gg 1$$

- Corotation **electric field**
- Sweepback of **B-field** due to poloidal current
- **Poynting flux** \Rightarrow electromagnetic energy losses



THEORETICAL (AND NUMERICAL) APPROACHES

Force-free
electrodynamics

Magnetized plasma without inertia

✓ OK in highly magnetized regions

- breaks when the existence of plasma is not a given, and in reconnection
- typical apps: neutron star magnetospheres, jets

Magnetohydrodynamics

Plasma as an ideal collisional fluid

✓ e.g., no thermal conduction, pressure is same in all directions; OK as a first approximation for global dynamics

- does not describe non-thermal particles
- typical apps: accretion flows

Kinetics

First-principles description for collisionless plasmas

✓ includes non-ideal effects (e.g., pressure is different along and across magnetic field, heat flux), describes particle acceleration

- computationally expensive and usually allows limited dynamic range
- typical apps: plasma instabilities, magnetospheres

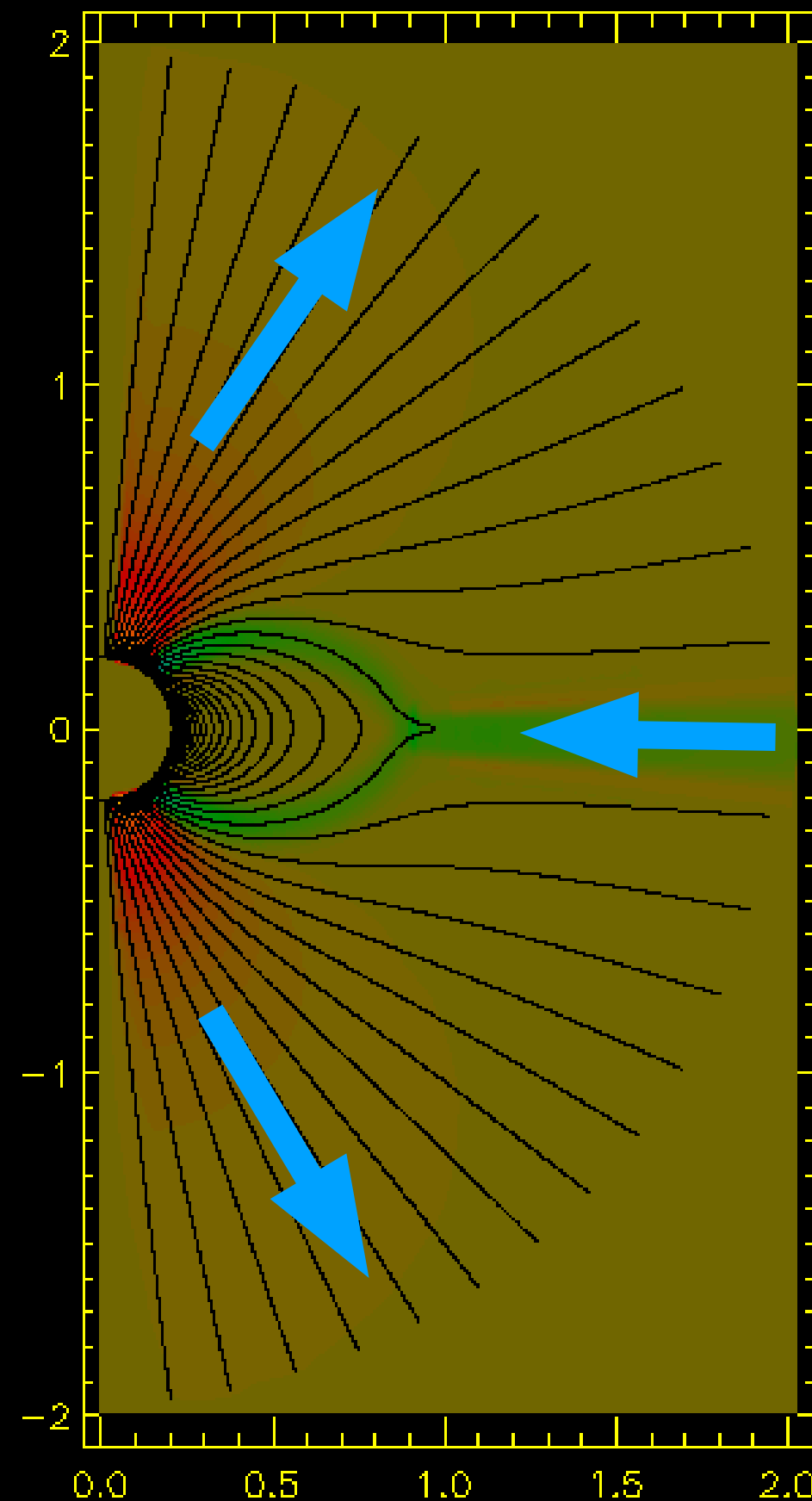
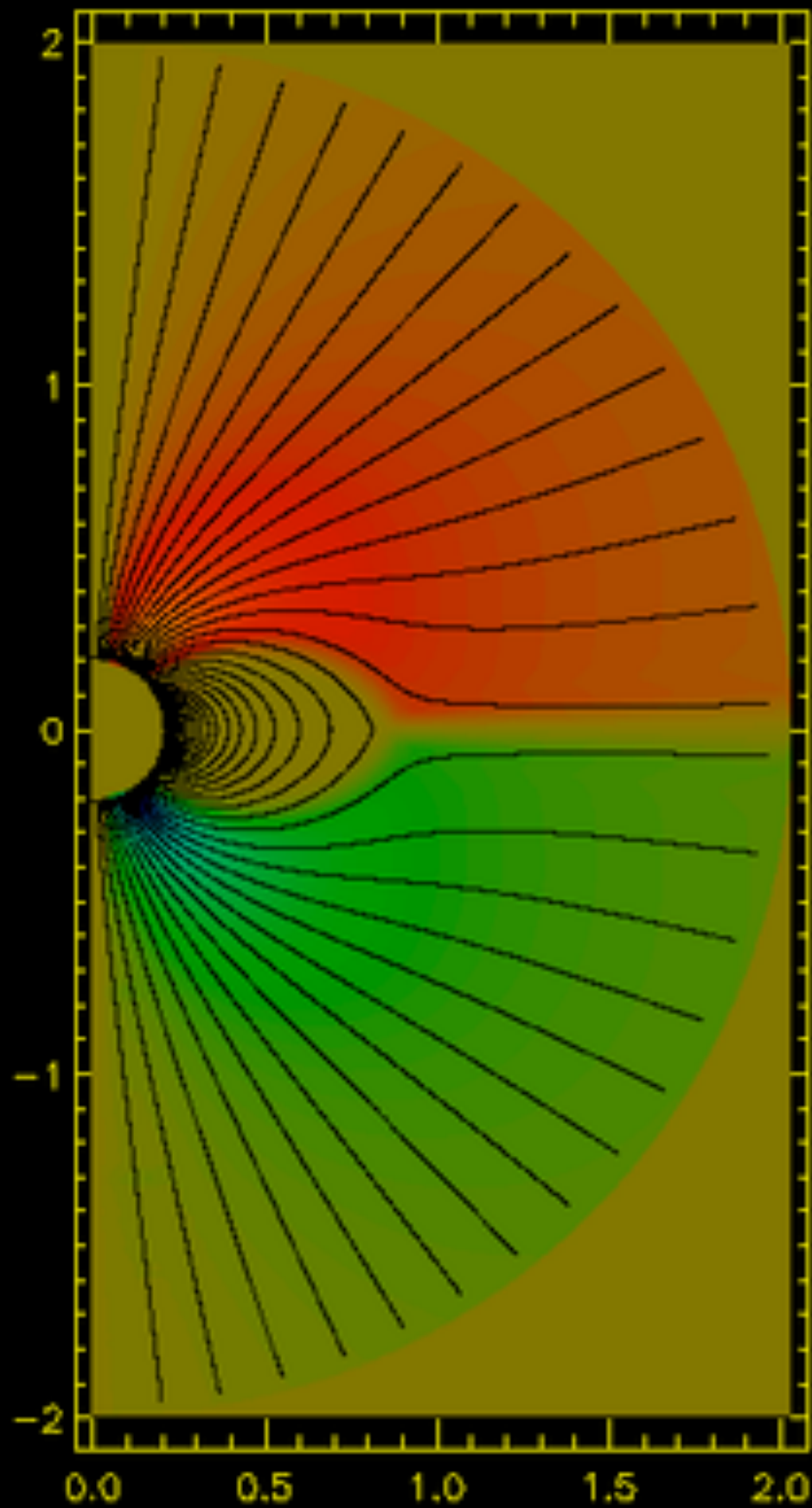
STANDARD PULSAR

Force-free paradigm

$$\rho_c \mathbf{E} + \mathbf{j} \times \mathbf{B} = \cancel{\frac{d\rho_n \mathbf{u}}{dt}} + \cancel{\text{pressure}}, \text{ and } \mathbf{E} \cdot \mathbf{B} = 0$$

+ Maxwell's equations

$$\Rightarrow \mathbf{j} = \mathbf{j}(\mathbf{E}, \nabla \times \mathbf{E}, \nabla \cdot \mathbf{E}, \mathbf{B}, \nabla \times \mathbf{B})$$



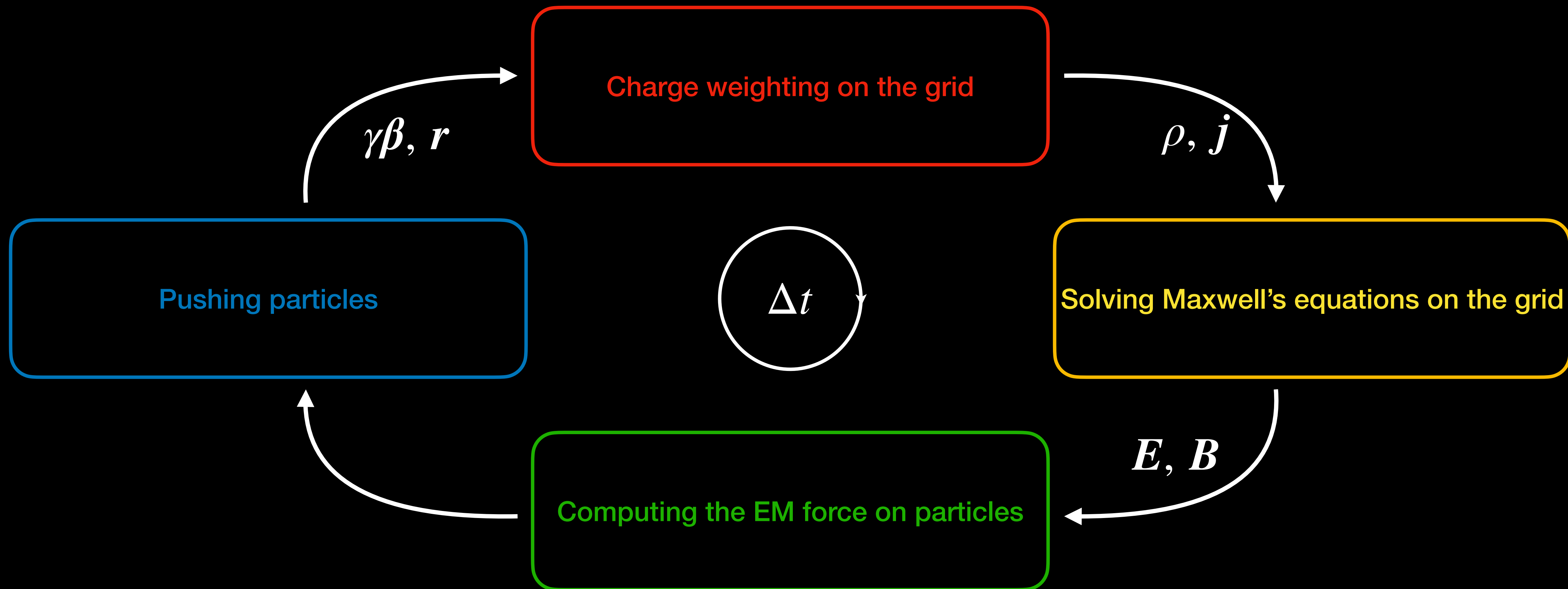
- closed-/open-field-line regions
- equatorial current sheet
- field lines are asymptotically radial
- predicts the spin-down law

$$L_{\text{PSR}} = k_1 \frac{\mu^2 \Omega^4}{c^3} (1 + k_2 \sin^2 \alpha)$$

- can not predict: particle acceleration, plasma supply, non-thermal radiation

Contopoulos+ (1999), Spitkovsky (2006), Kalapotharakos (2009), Petri (2012), Tchekhovskoy+ (2014) (MHD)

PLASMA PHYSICS ON A COMPUTER: (GR)(R)PIC



(GR) = general relativistic

(R) = radiation reaction force, photon emission, multiple pair production mechanisms

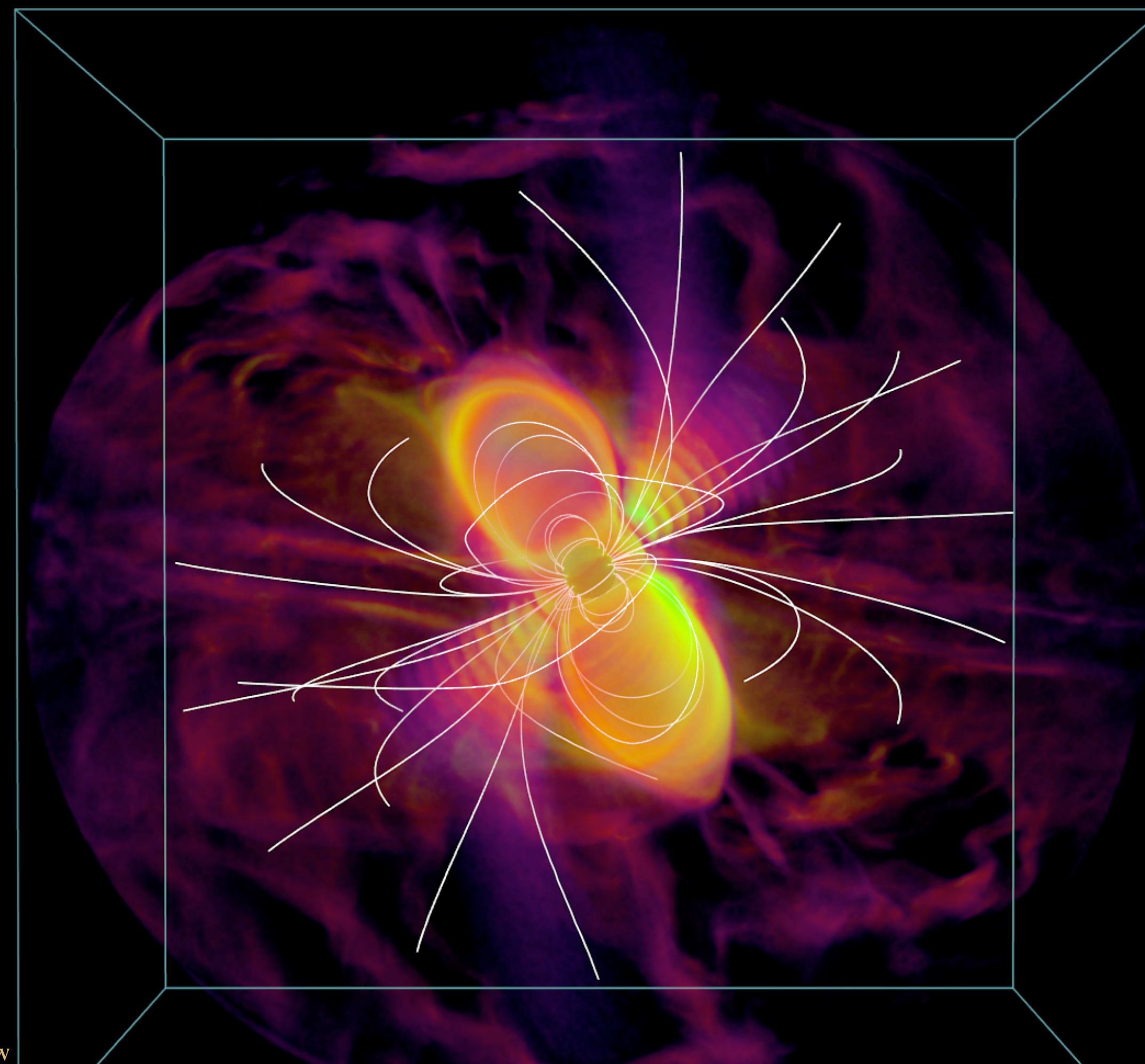
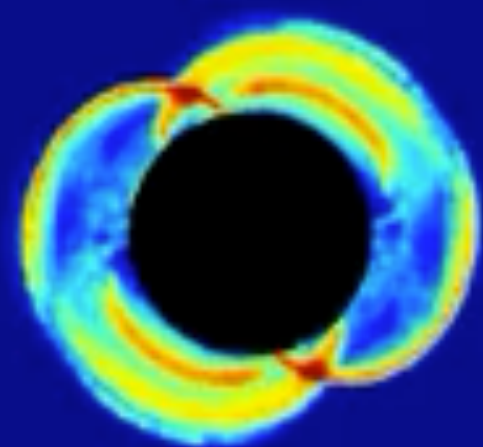
PIC = particle-in-cell

**MAGNETOSPHERIC STRUCTURE AND
HIGH-ENERGY EMISSION**

(GR) OBLIQUE ROTATOR WITH PAIR PRODUCTION

Philippov, Spitkovsky, 2018 (ApJ)

n_{\pm}

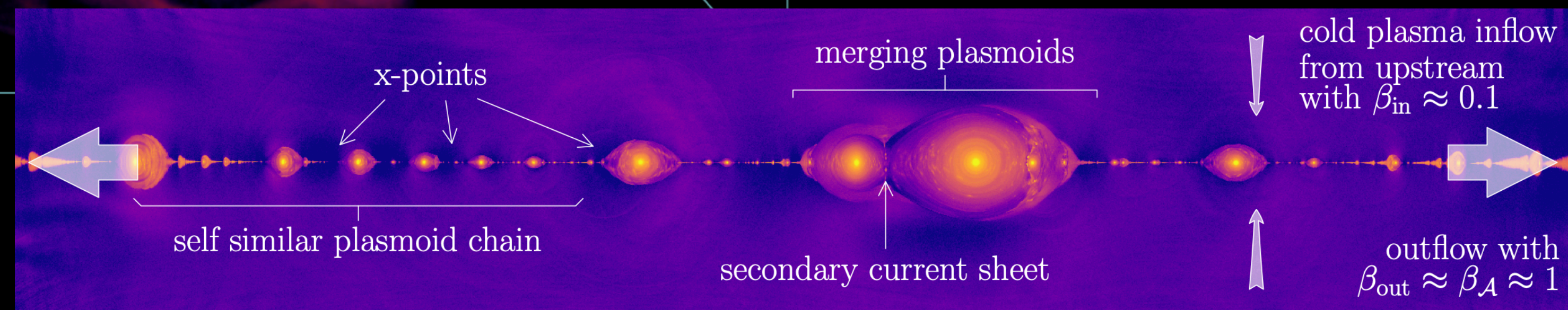
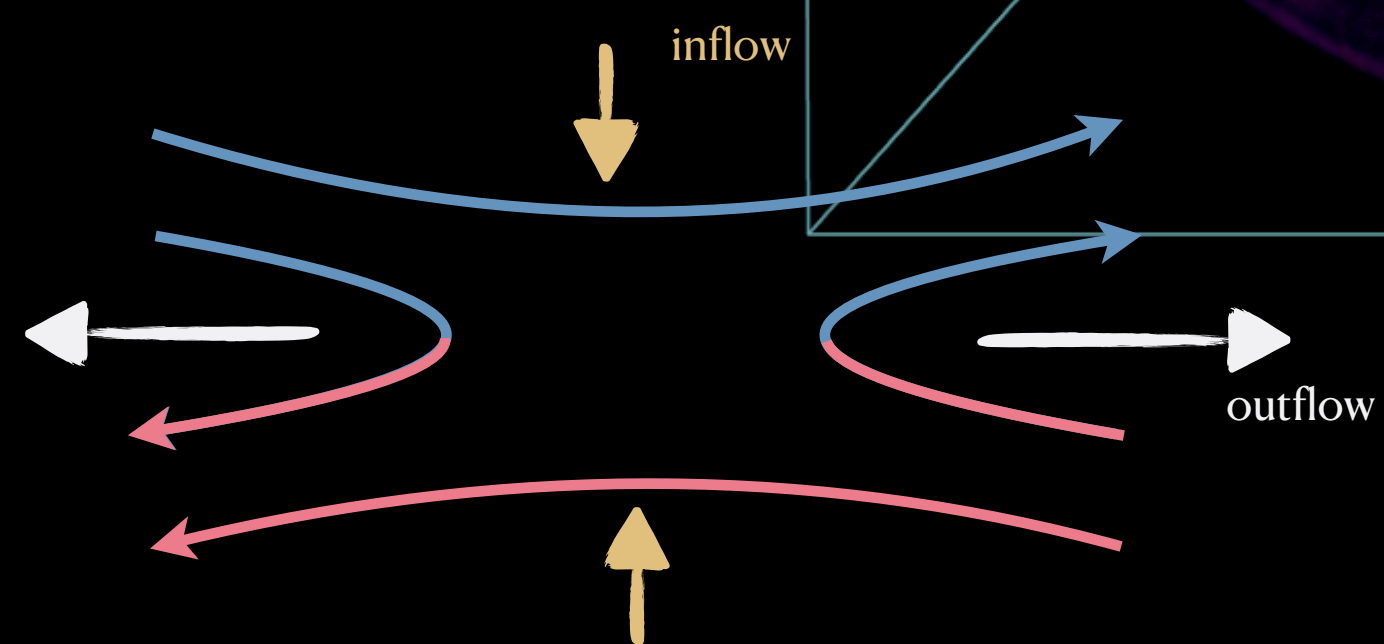


- Non-stationary discharge powers coherent radio emission.
- Relativistic magnetic reconnection in the current sheet powers high-energy emission. Its rate controls the radiated power.
- Current sheet is unstable to plasmoid (tearing) and drift-kink instabilities.

Hakobyan, Philippov et al., 2019 (ApJ)

(and many others after Loureiro et al., 2007)

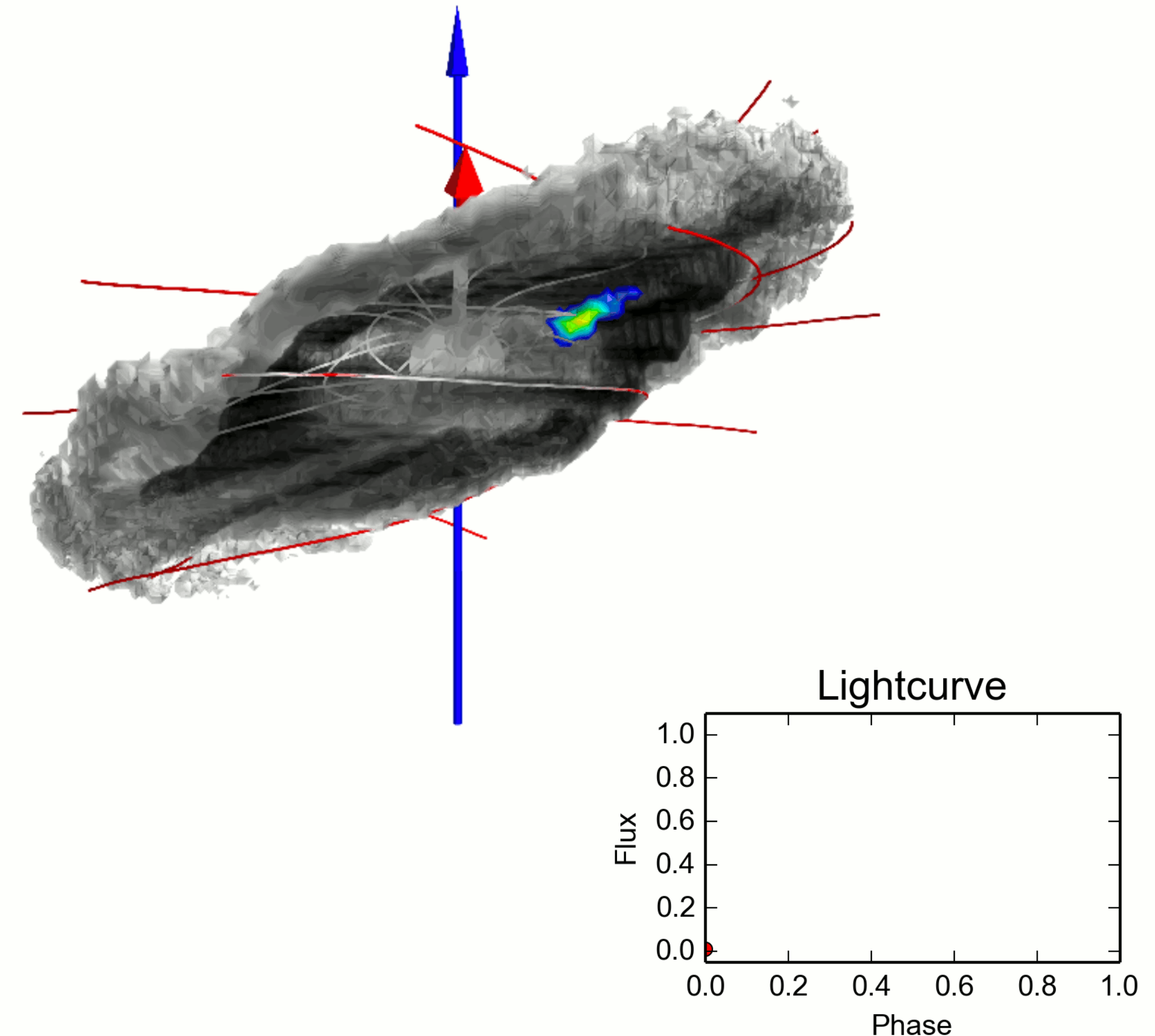
2D:



GAMMA-RAY EMISSION MODELING

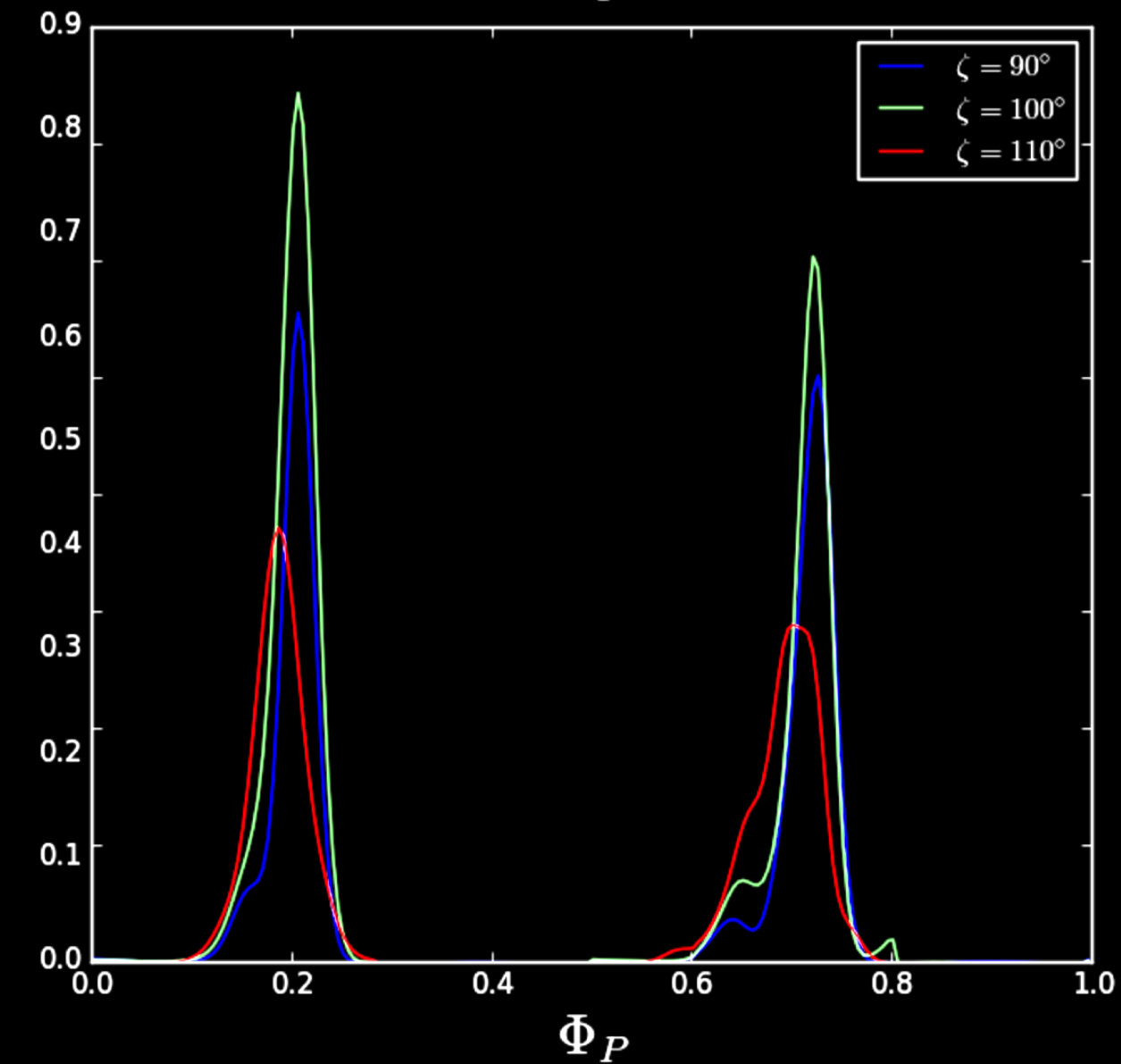
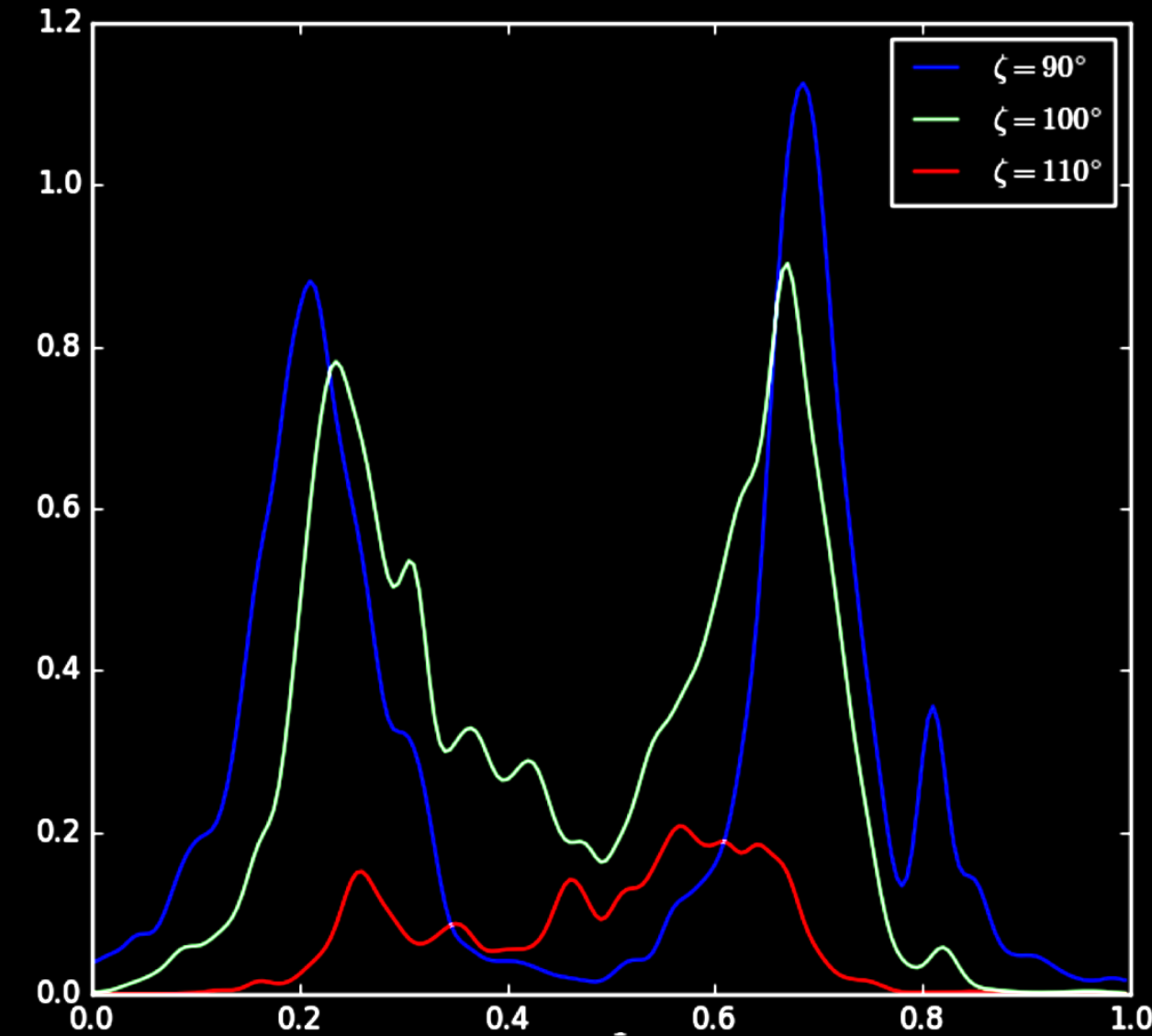
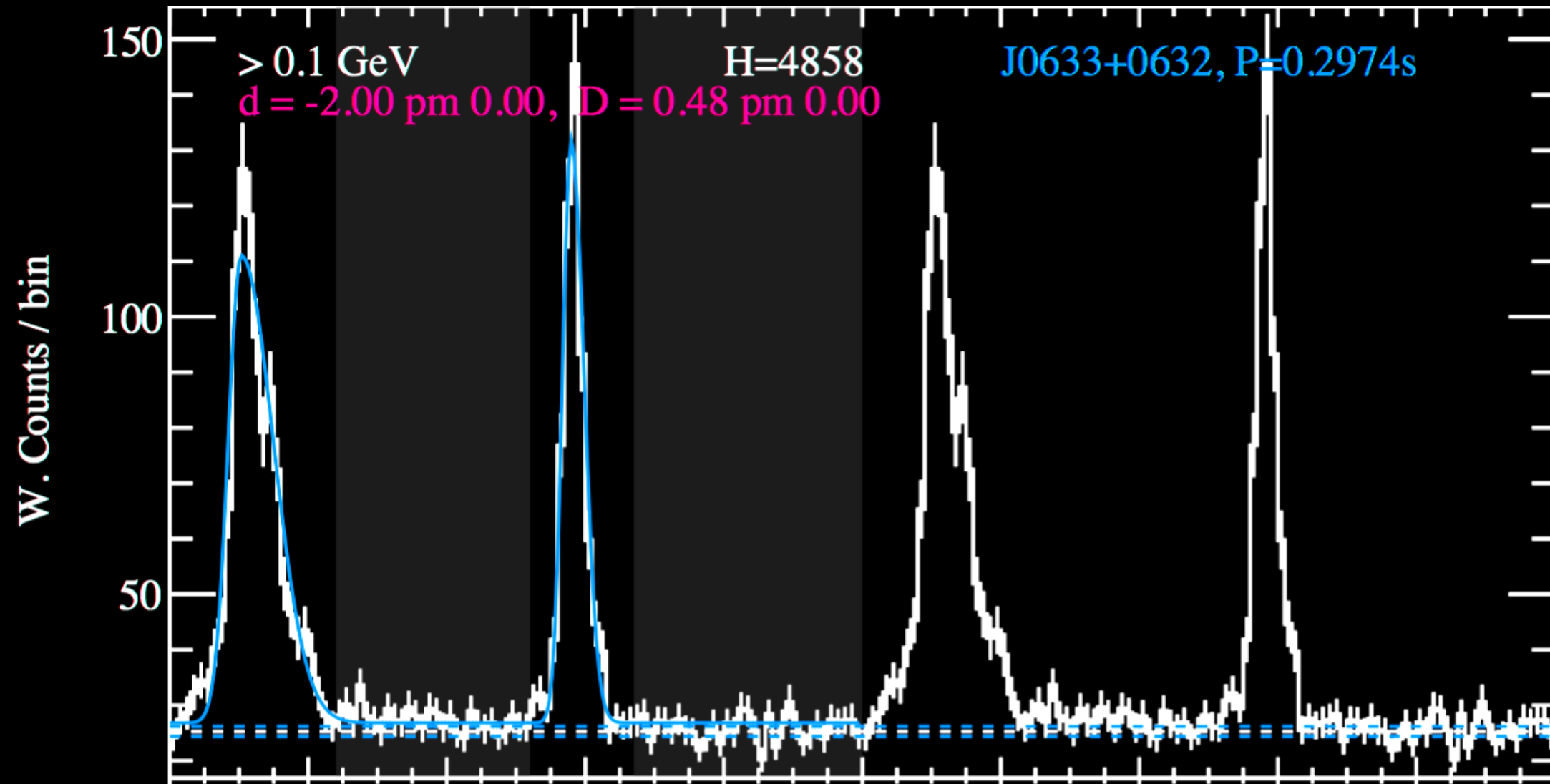
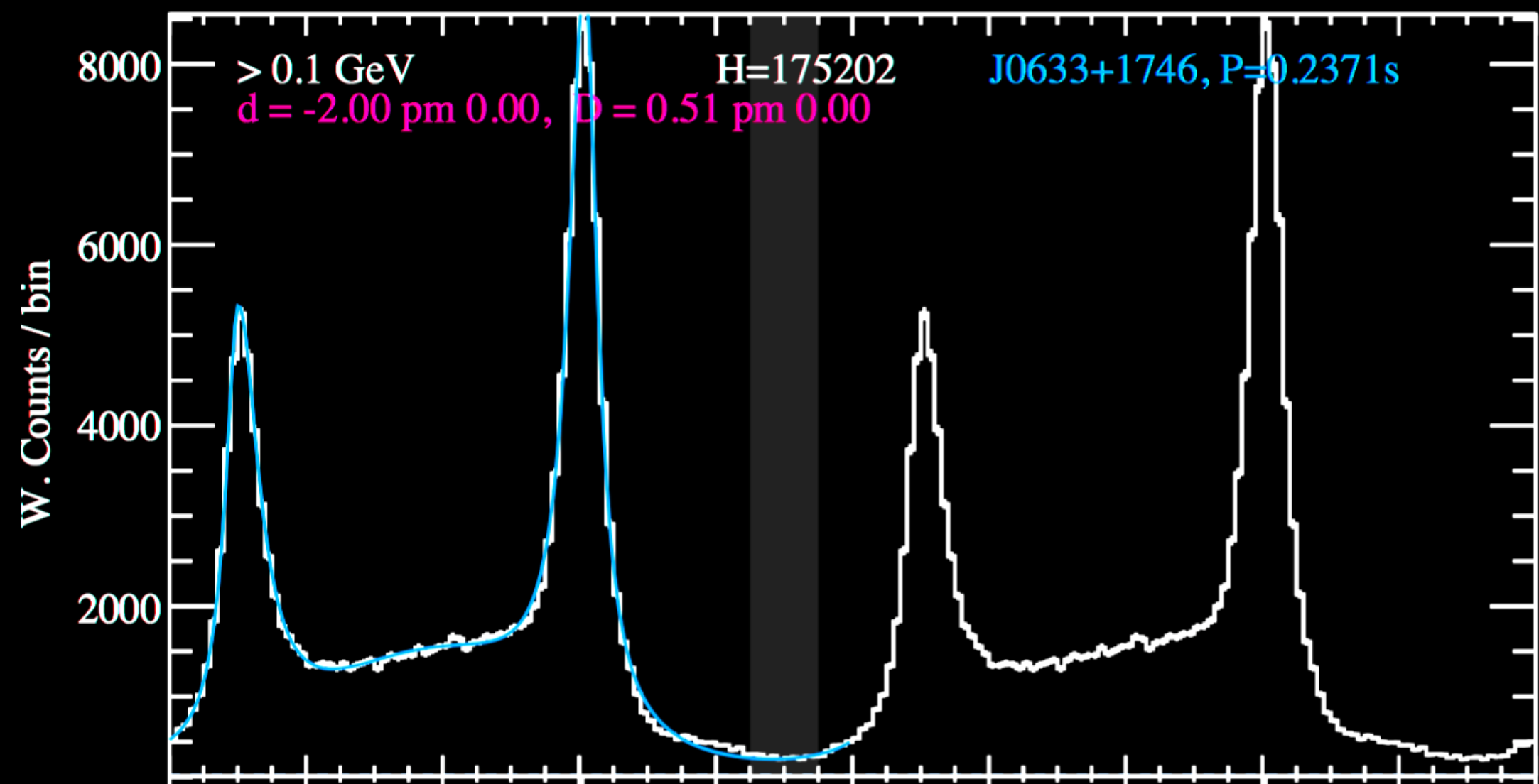
- Simulations prefer current sheet as a particle accelerator
- Particles radiate synchrotron emission
- Observe caustic emission
- Predict gamma-ray efficiencies 1-20% depending on the inclination angle and pair production efficiency in the sheet
- Higher inclinations are less dissipative

$i=30$ - Phase=0.00 - Positrons -



LIGHTCURVES

FERMI



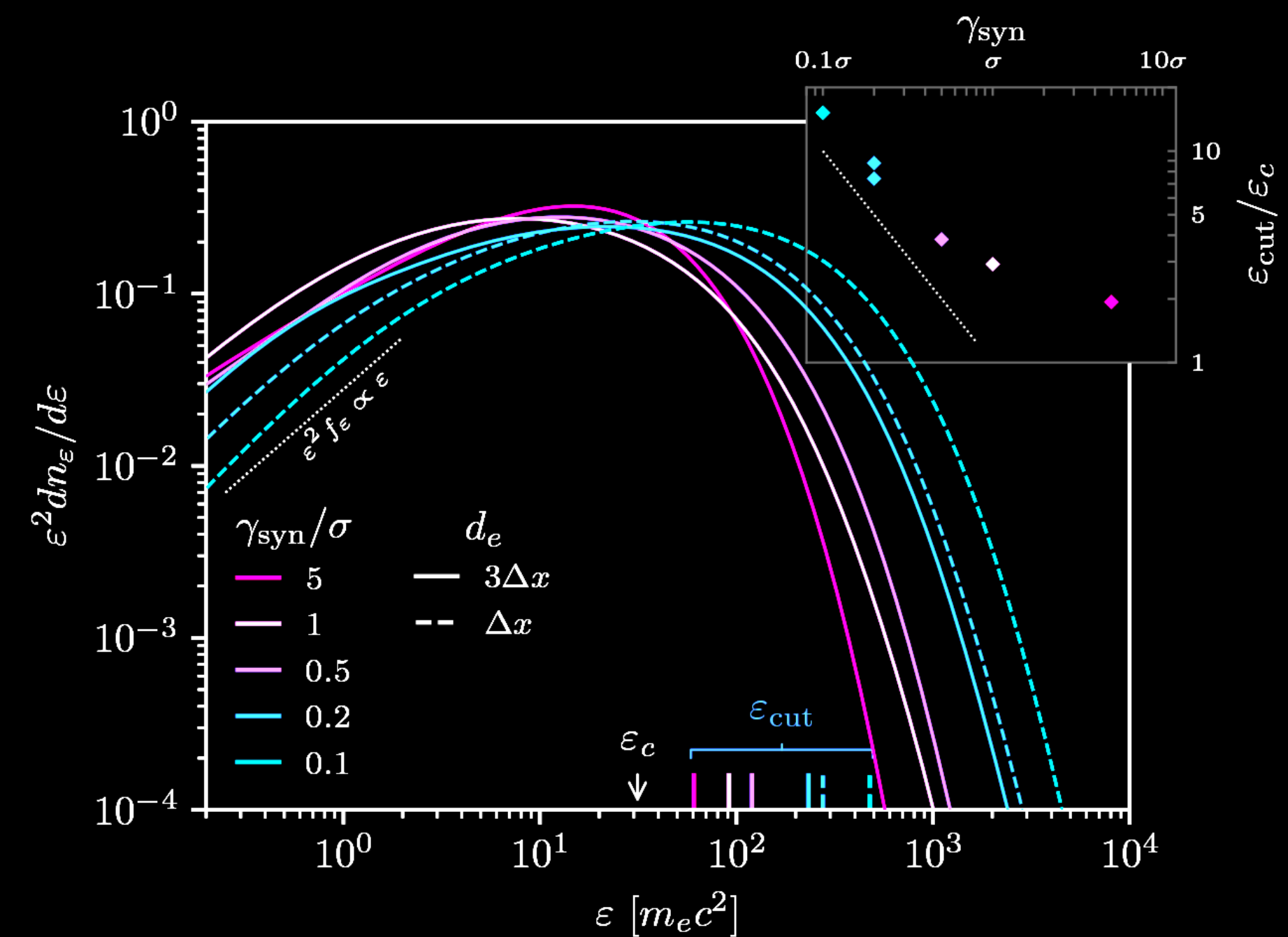
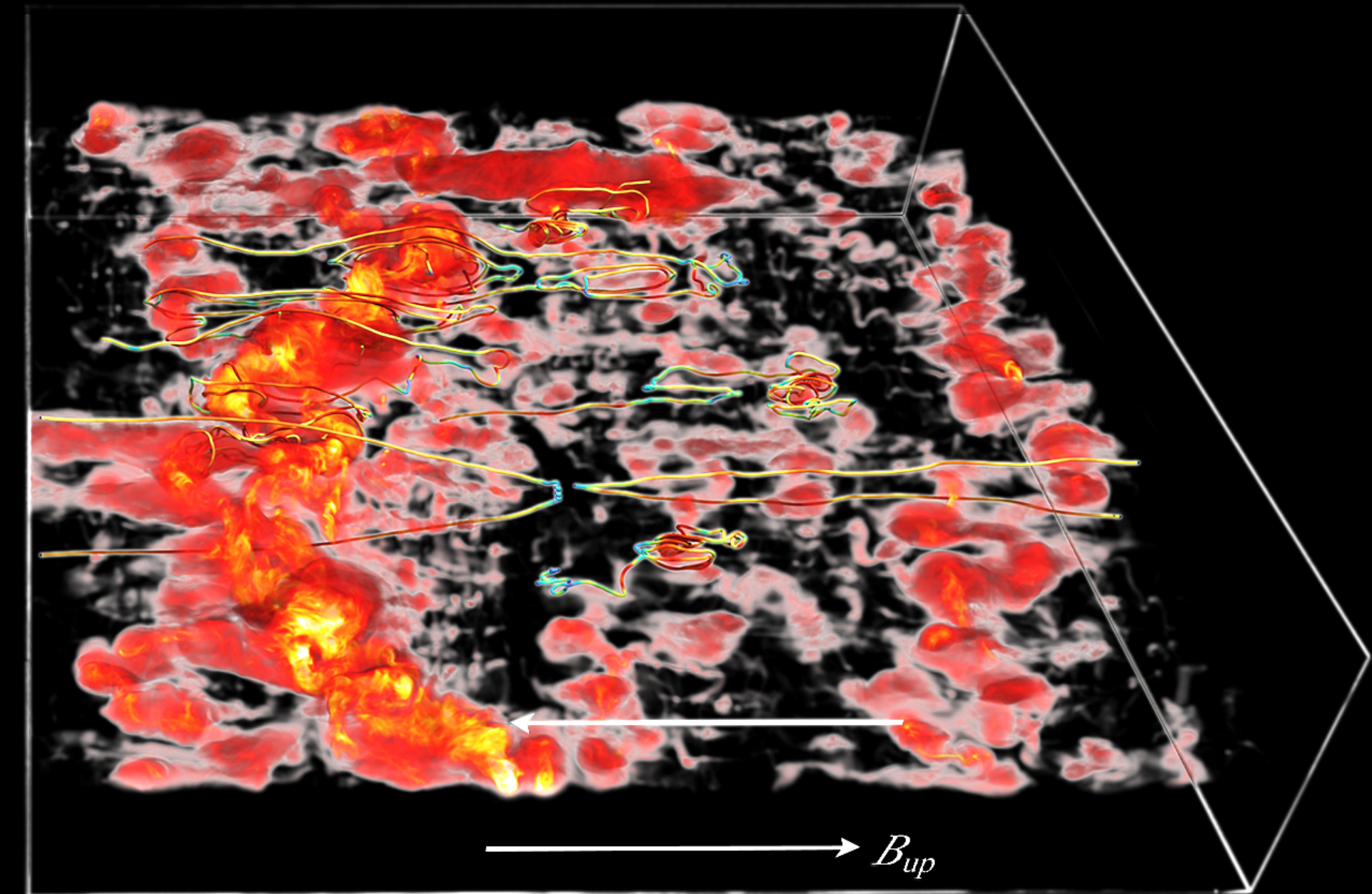
Double-peaked lightcurves are generic

Philippov, Spitkovsky, 2018 (ApJ)

RECONNECTION IN PULSAR MAGNETOSPHERES

- $B \sim 10^5 \text{ G}$, $\sigma = B^2 / (4\pi\rho_m c^2) \gg 1$
- Reconnection electric field accelerates particles, synchrotron cooling is important on the same timescale
- Pairs accelerate beyond the radiation reaction limit, up to $\gamma \sim \text{few} \times \sigma$
- Highest energy photons are beamed along the upstream magnetic field, consistent with the beaming of GeV lightcurves

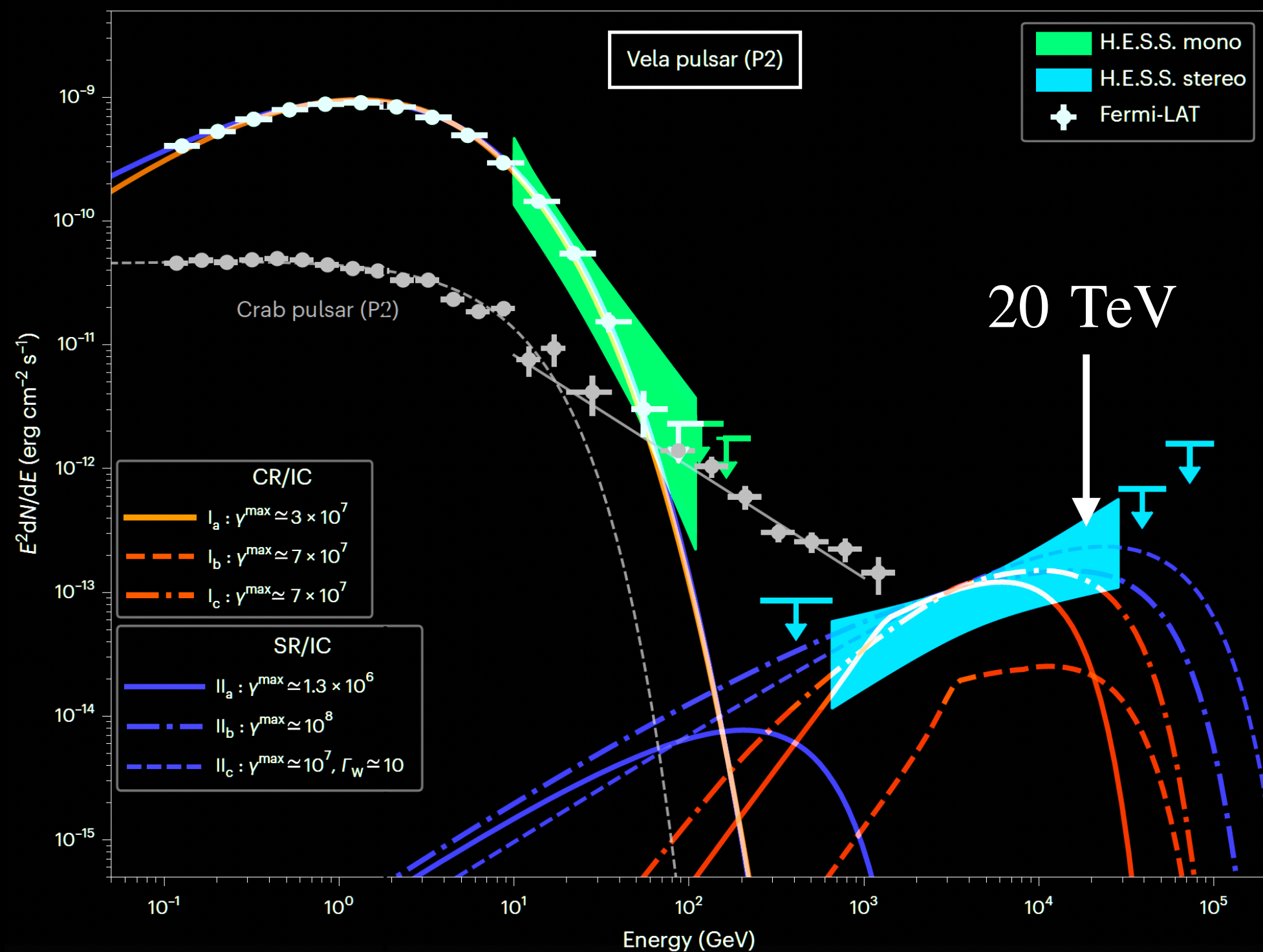
$$h\nu_{\text{max}} \approx 16 \text{ MeV} \cdot \left(\sigma / \gamma_{\text{syn}} \right)$$



NEW FRONTIER: MULTI-TeV FROM VELA PULSAR [IN PREP]



The H.E.S.S. Collaboration, Nature (2023)



Bransgrove et al, 2023 (ApJL)

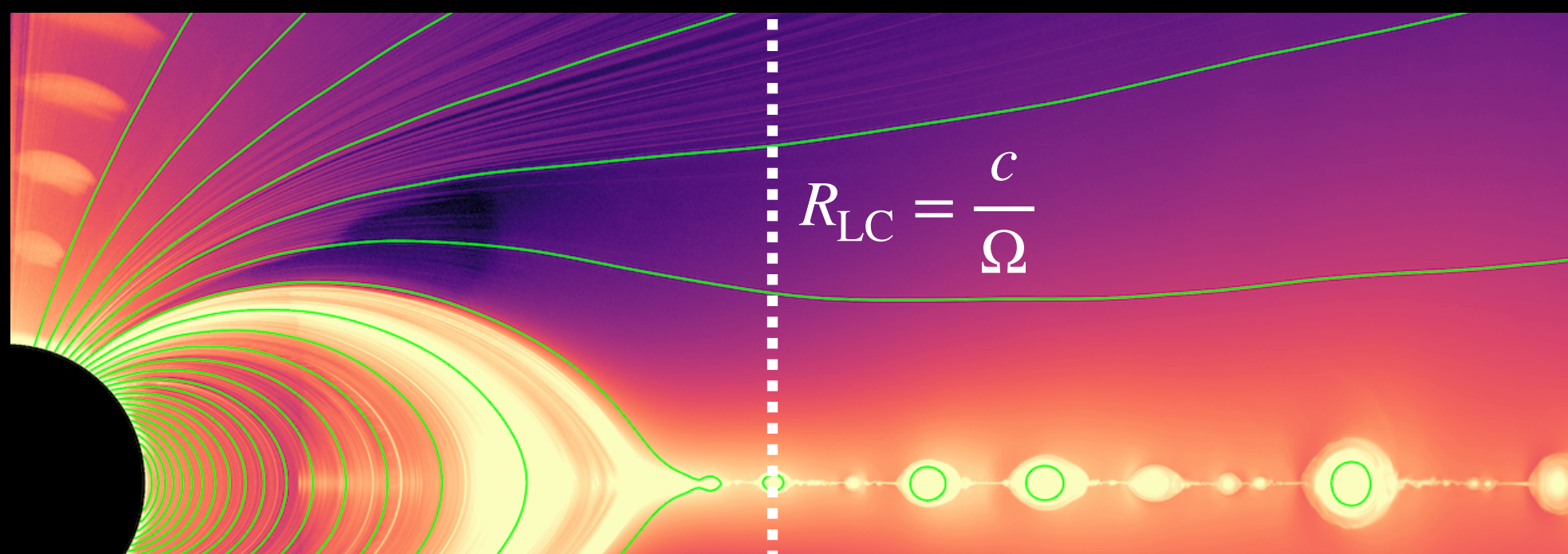
$$\gamma_{\text{syn}} \approx 10^5 \Rightarrow \sigma \approx \text{few} \times 10^7$$



$$\epsilon_{\text{ph}} = 16 \text{ MeV} \cdot \left(\sigma / \gamma_{\text{syn}} \right)$$

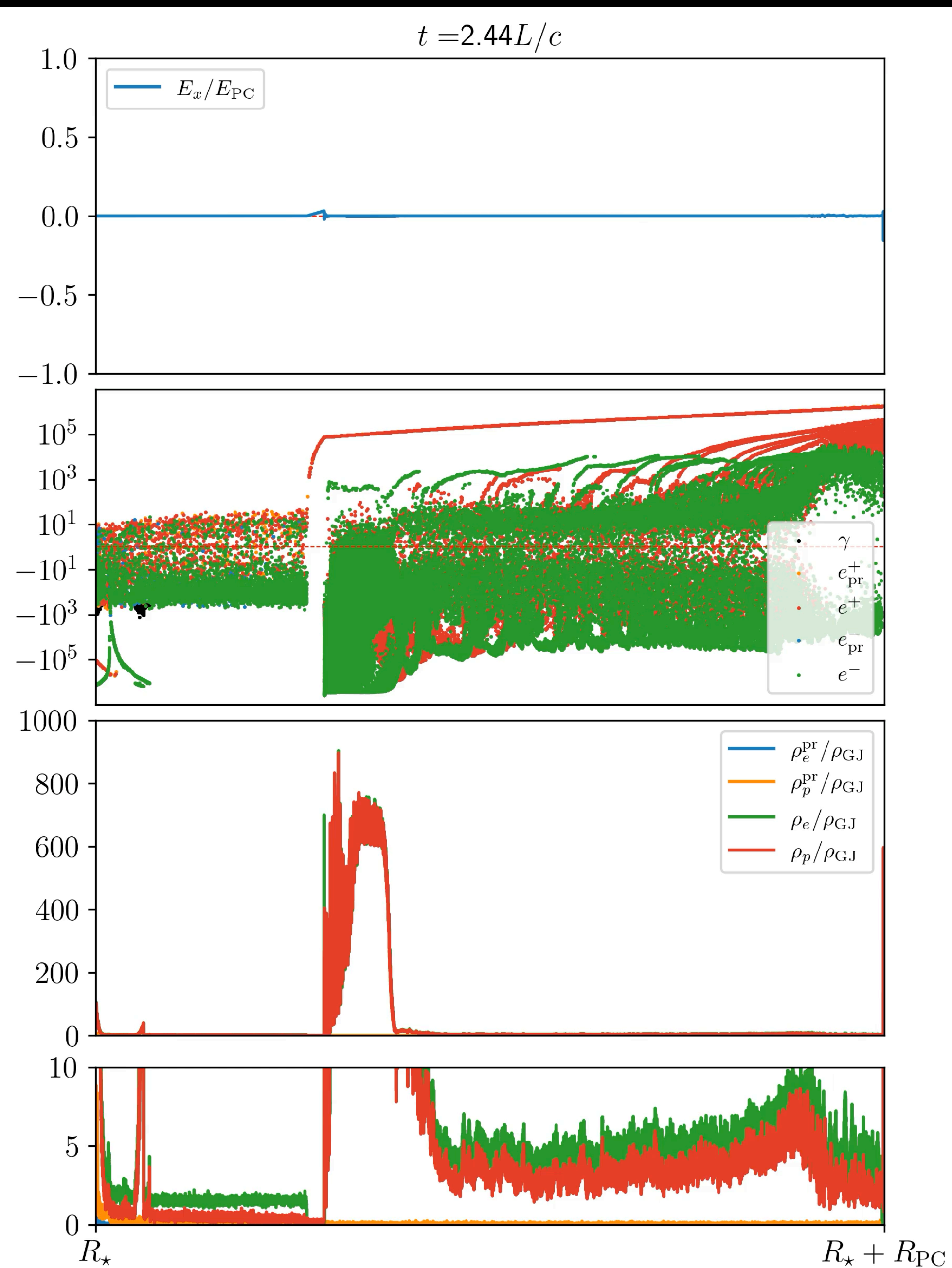
$$m_e c^2 \gamma_{\text{max}} = m_e c^2 \sigma \sim 10 \text{ TeV}$$

- Pair density is low because "return"-current discharge sends most of the plasma into the star
- Most of the plasma is produced in the current sheet



Prediction: CTA will see moderately energetic γ -ray pulsars as multi-TeV sources

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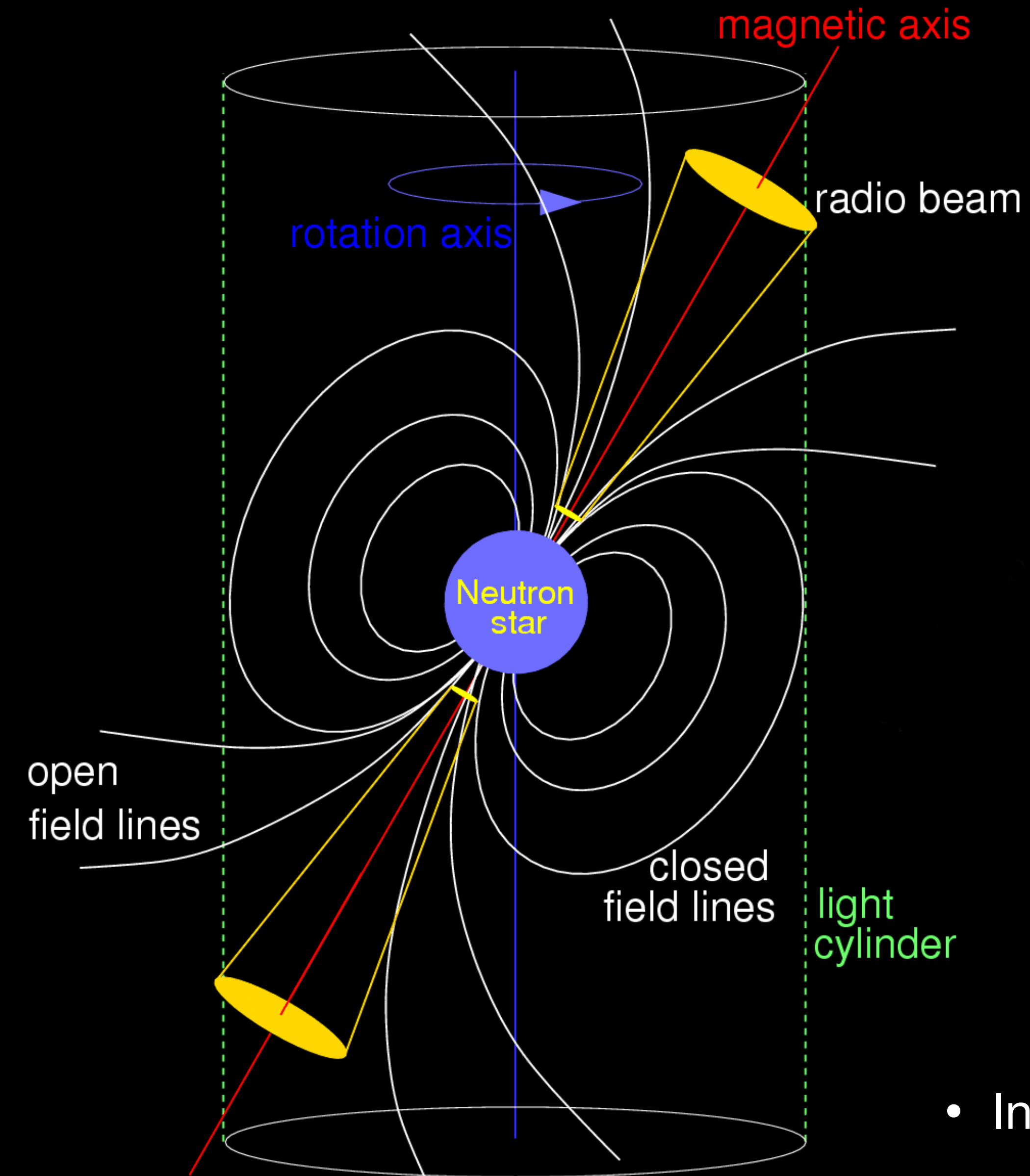
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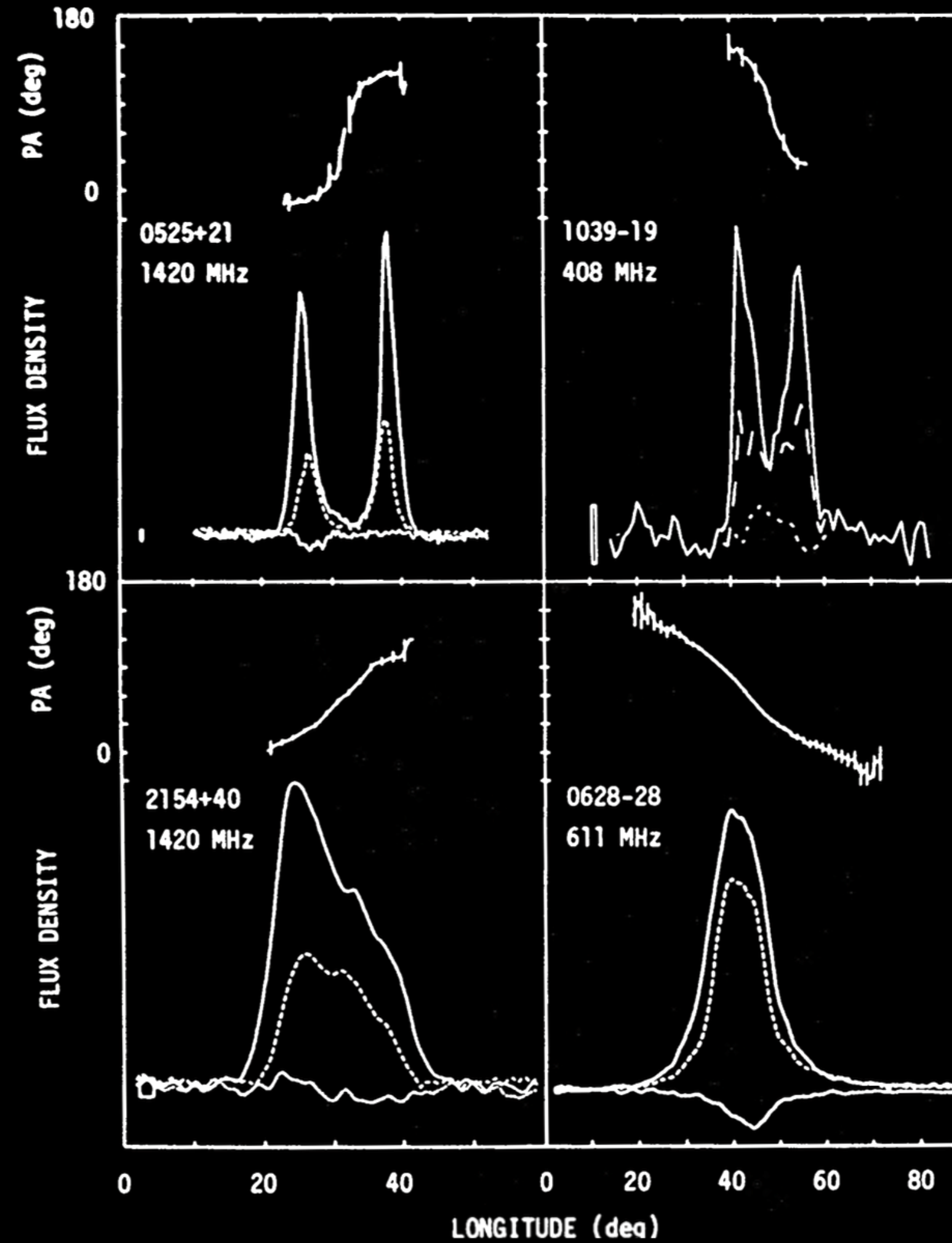
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RADIO EMISSION

POLAR RADIO EMISSION



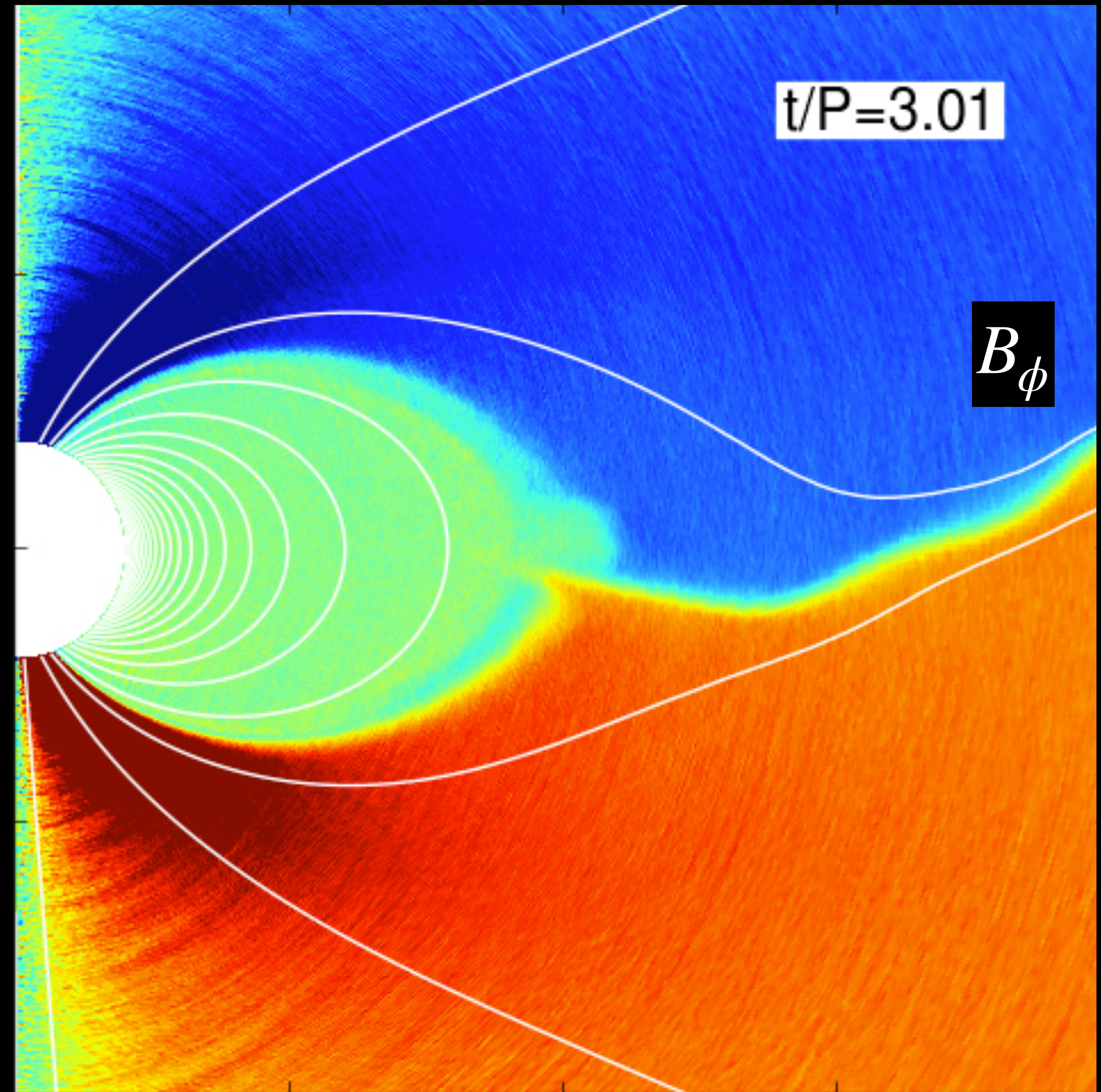
Lyne, Manchester (1988)



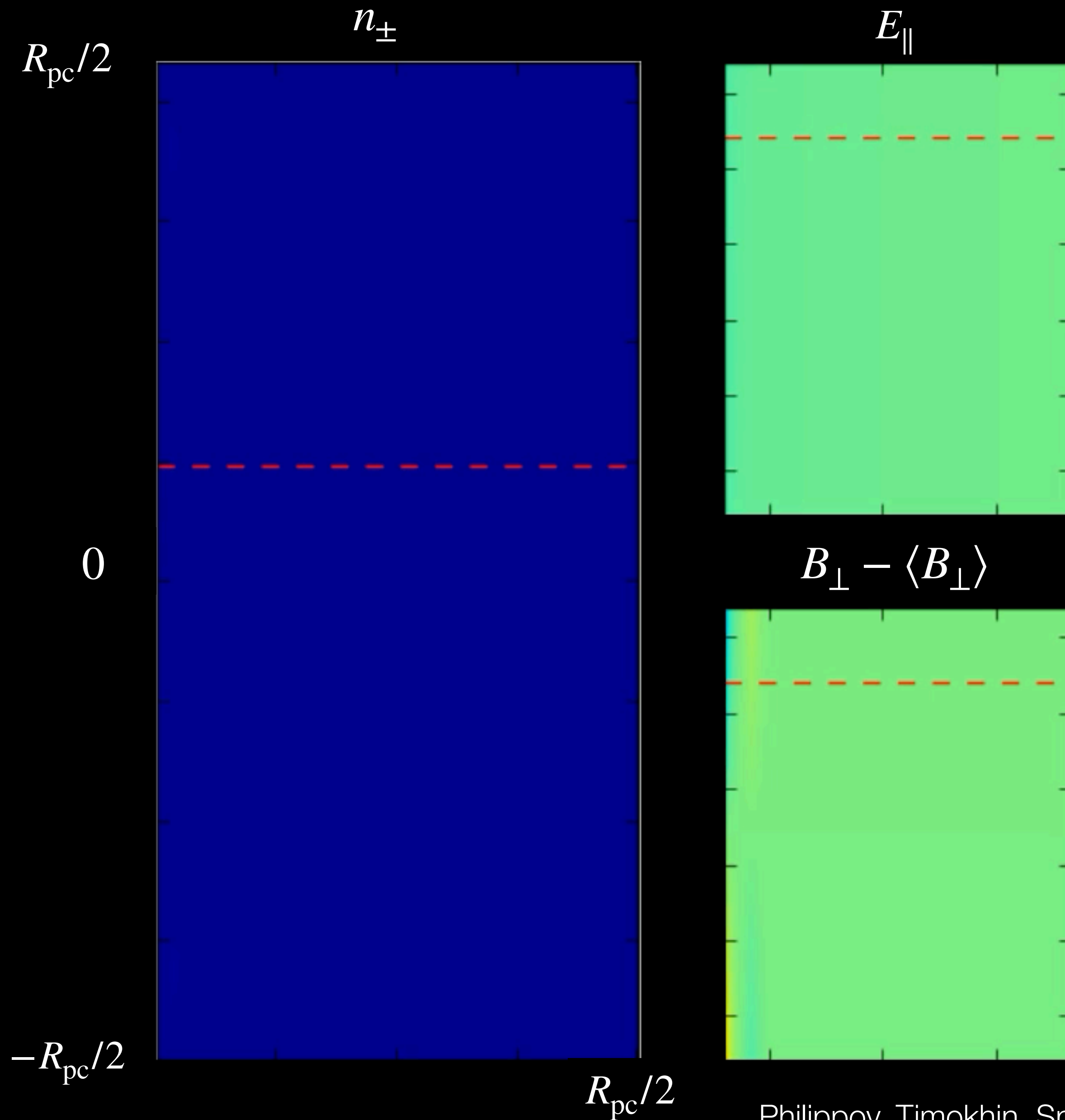
- In most cases we see one short pulse per period
- Beam width is related to the polar cap size

HINTS FROM GLOBAL SIMULATIONS

- ▶ Non-stationary discharge drives waves in the polar region.
- ▶ Waves are generated during the process of electric field screening. They are driven by collective plasma motions, thus, coherent (see also Beloborodov 2008; Timokhin, Arons 2013).



LOCAL SIMULATION OF 2D DISCHARGE

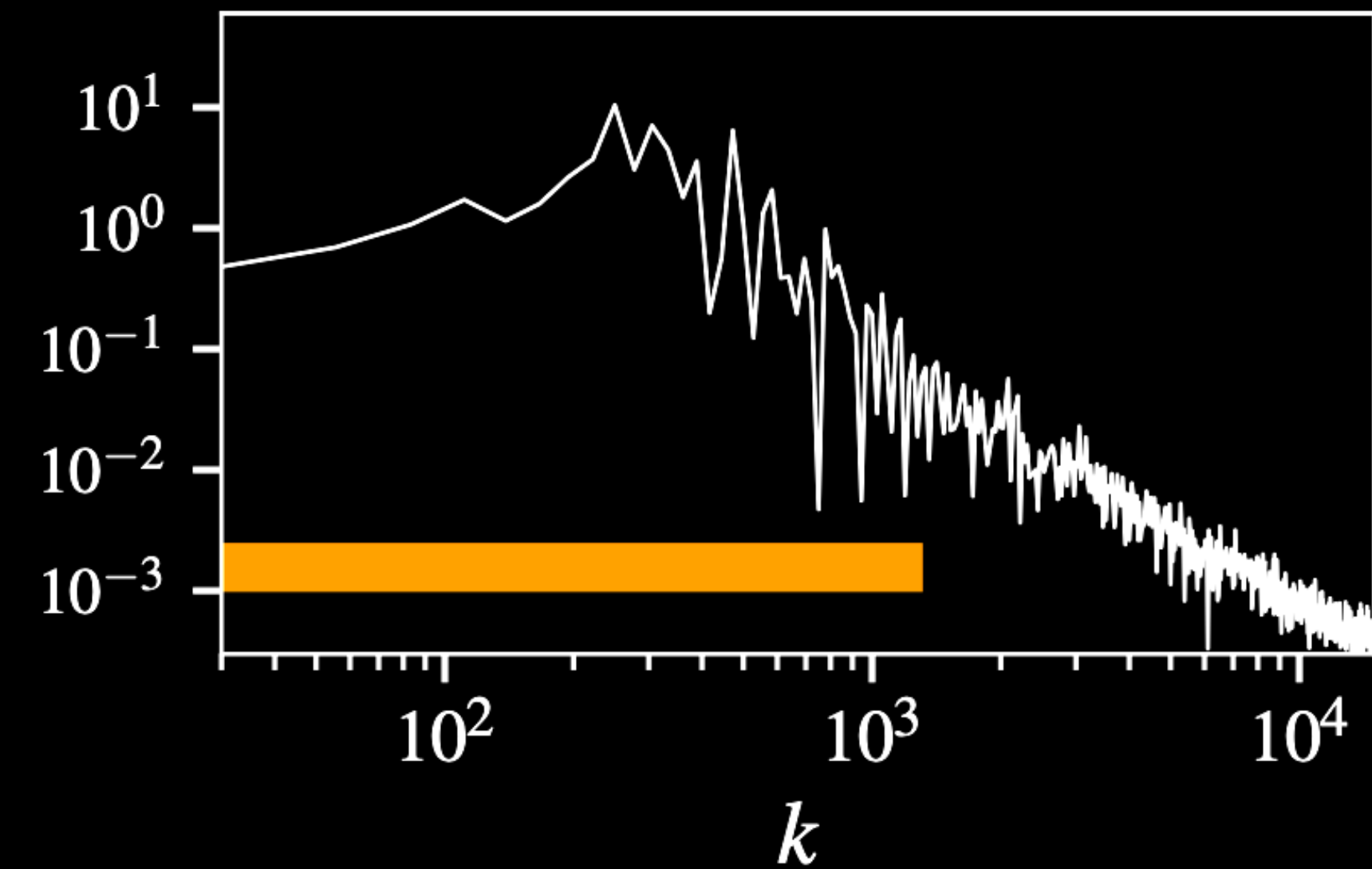
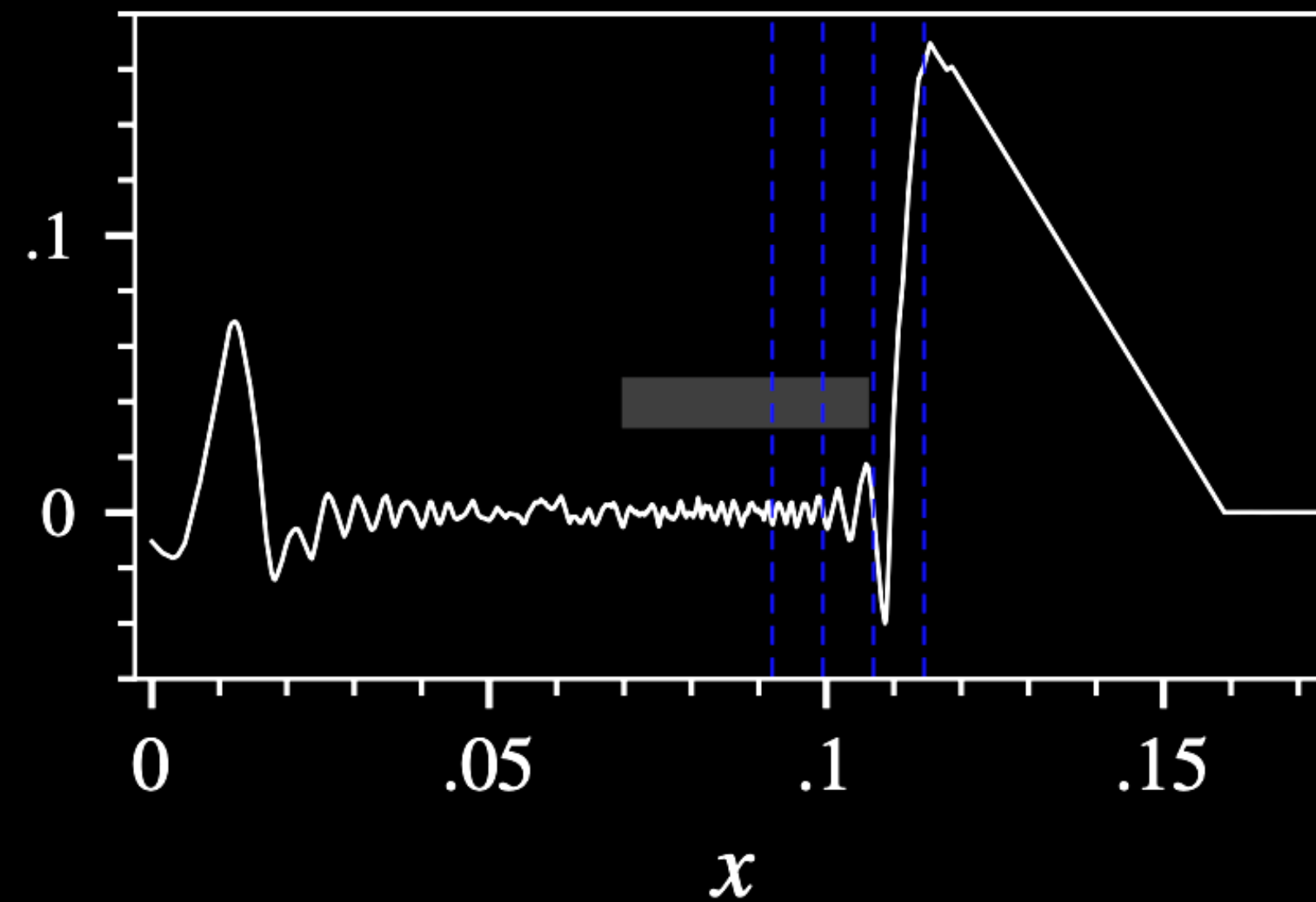
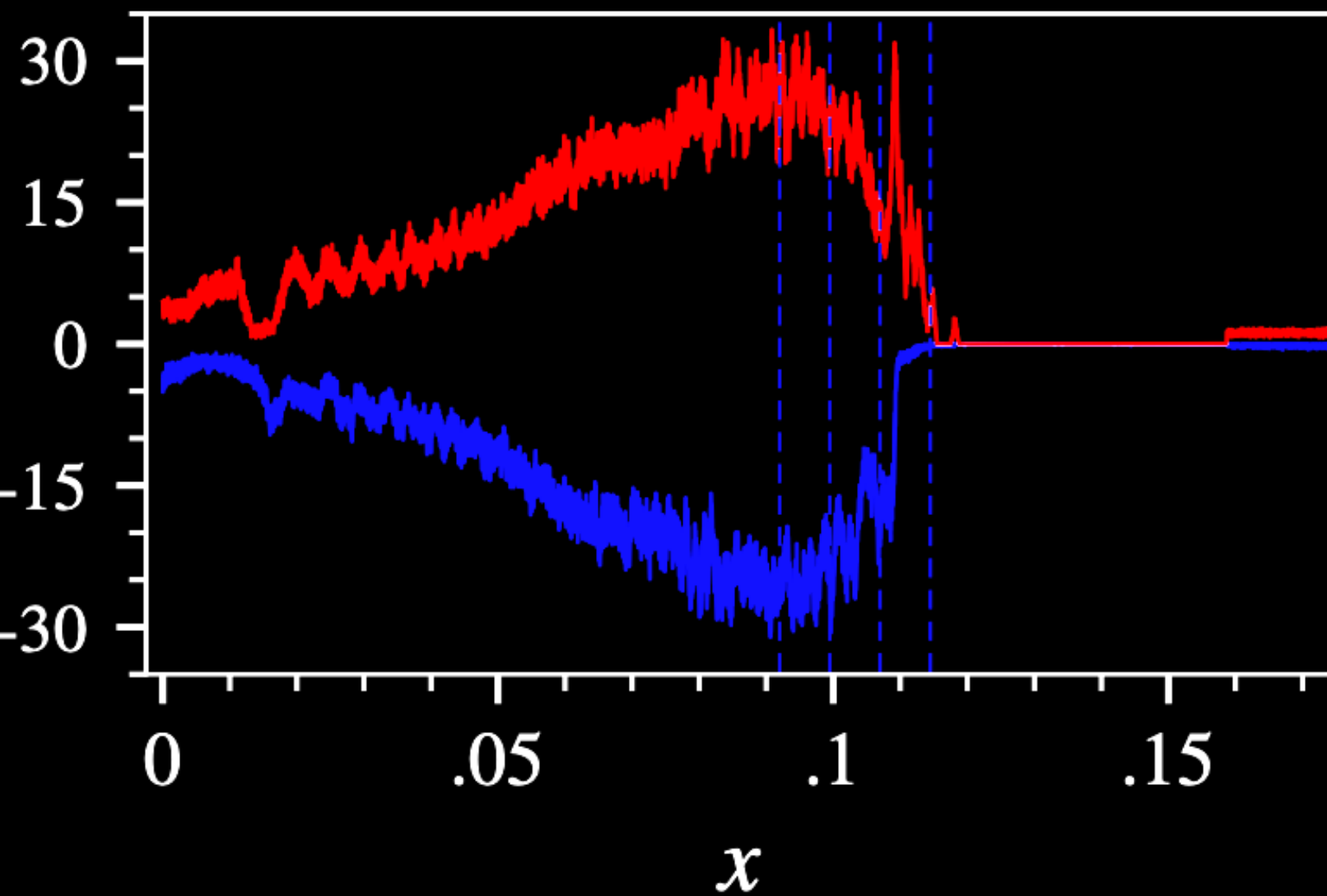


- Intermittency of the discharge results in production of coherent currents that are "screening" the electric field
- Oblique "screening" waves are electromagnetic and superluminal; thus, can escape from the magnetosphere
- The power is fixed at $\sim 10^{-4} L_{sd}$

SPECTRUM OF A 1D DISCHARGE



- Power cascades to a maximum plasma frequency in the cloud
- Clearly a very broad-band mechanism



$$\nu \simeq \sqrt{4\pi e^2 \kappa n_{\text{GJ}} / \langle \gamma^3 \rangle m_e^3} / 2\pi \sim 26 \sqrt{\kappa_5 B_{12} / r^3 P_{0.1} \gamma_{10}^3} \text{ GHz}$$

Prediction: close-by young pulsars should be ALMA sources

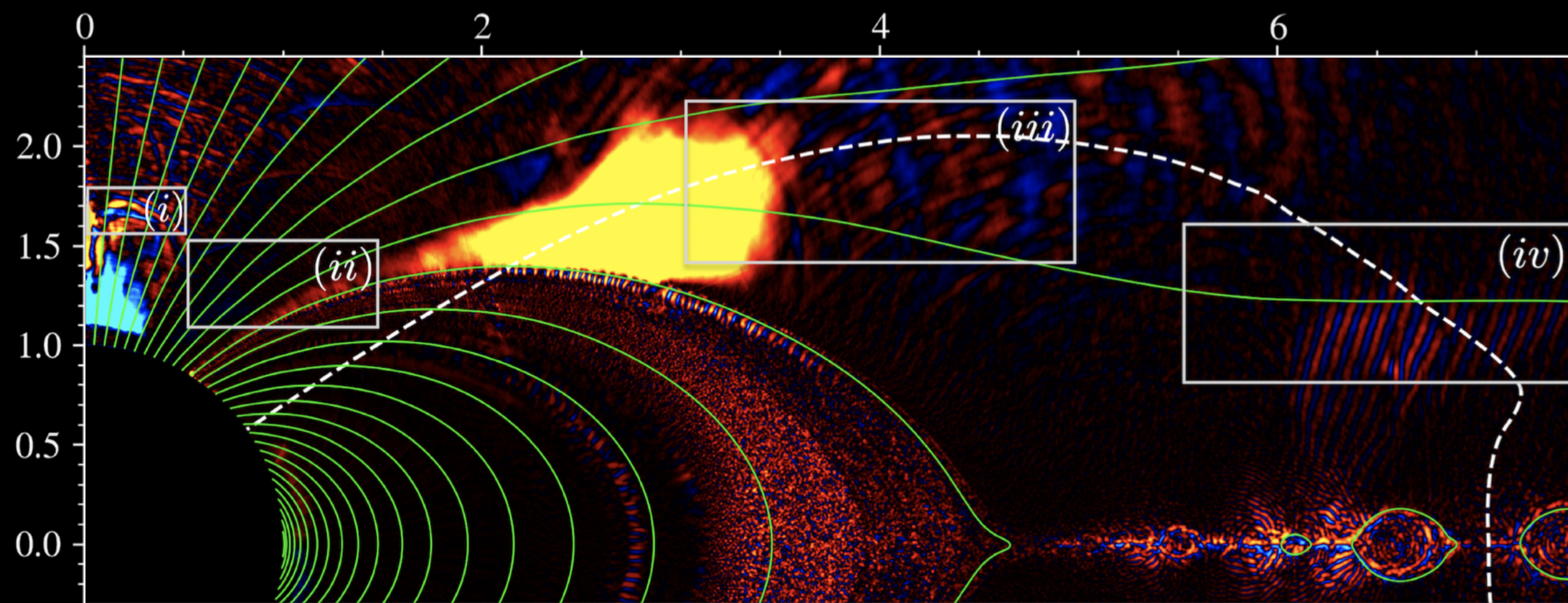
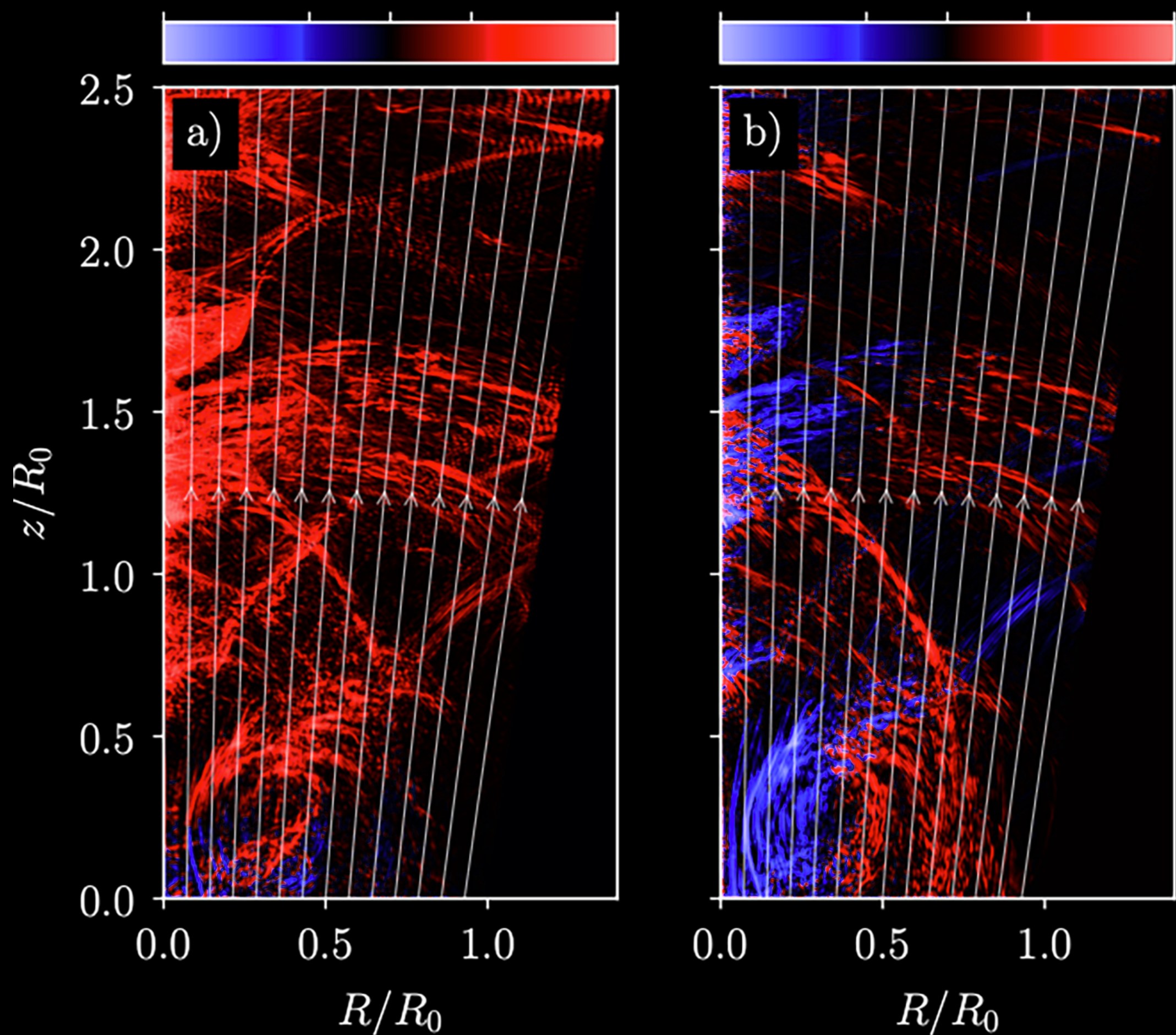
CONFIRMATION WITH DIFFERENT NUMERICAL CODES

$$tc/R_0 = 5.40$$

$$4\pi S_z / ce^2 n_{GJ}^2 R_0^2$$

$$4\pi S_R / ce^2 n_{GJ}^2 R_0^2$$

-10^{-3} -10^{-5} 0 10^{-5} 10^{-3} -10^{-3} -10^{-5} 0 10^{-5} 10^{-3}

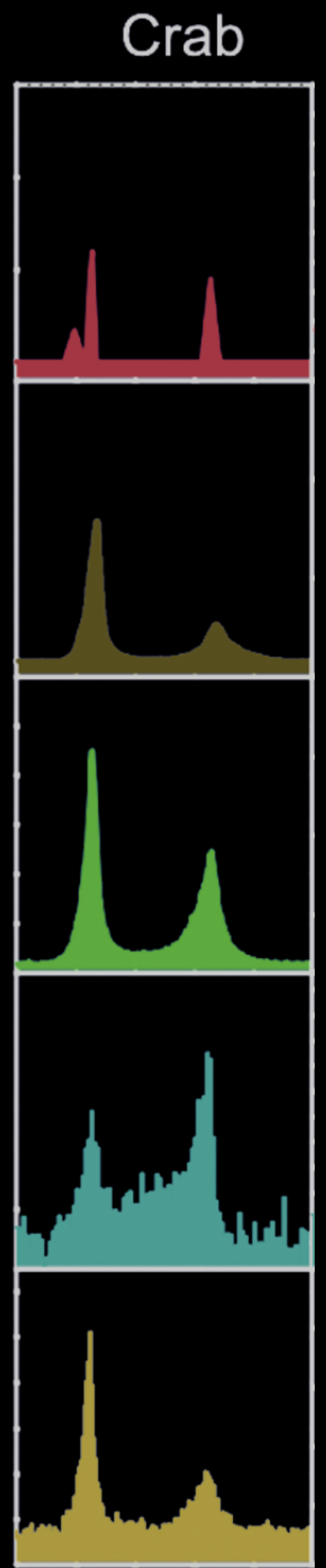
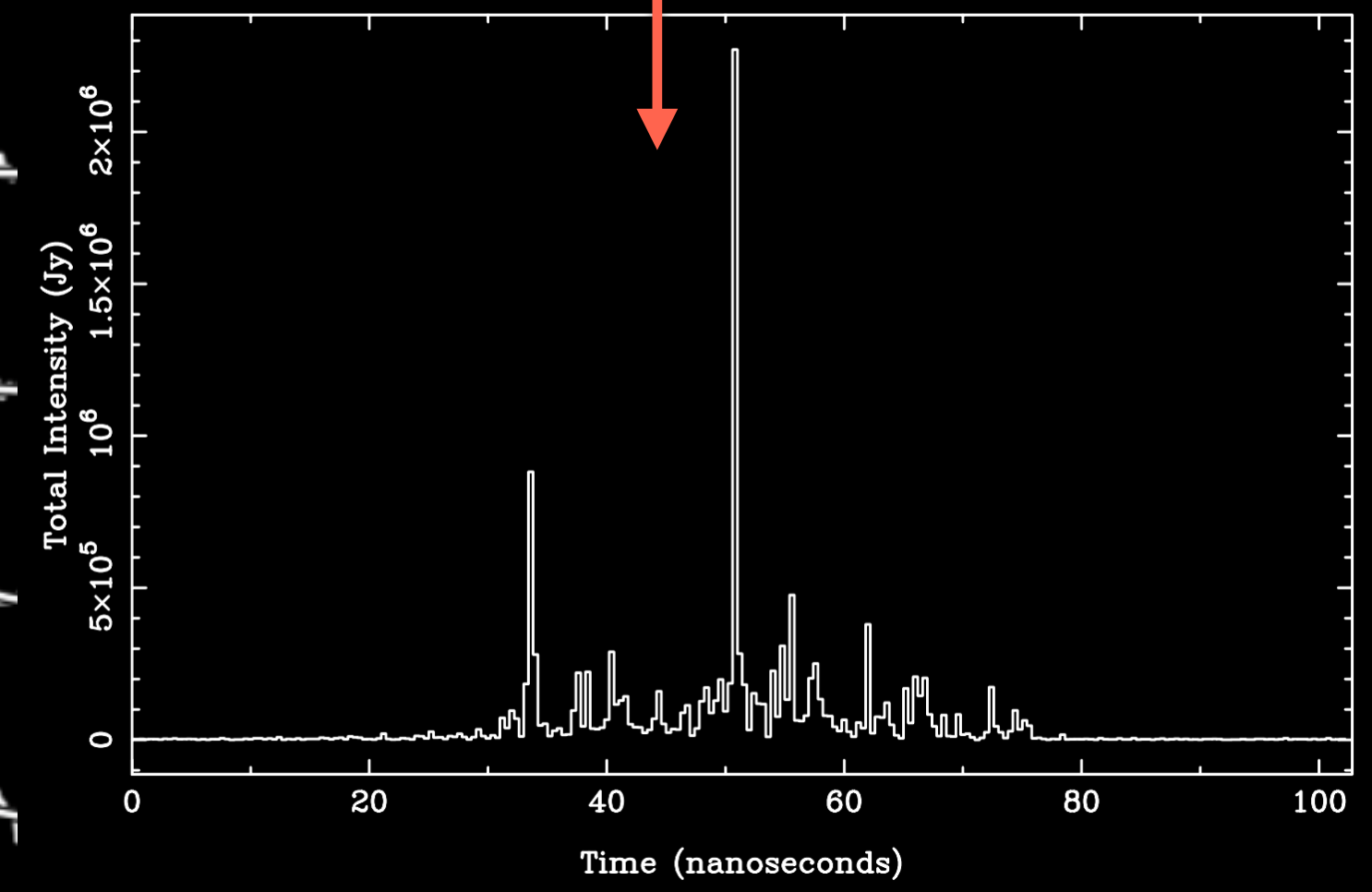
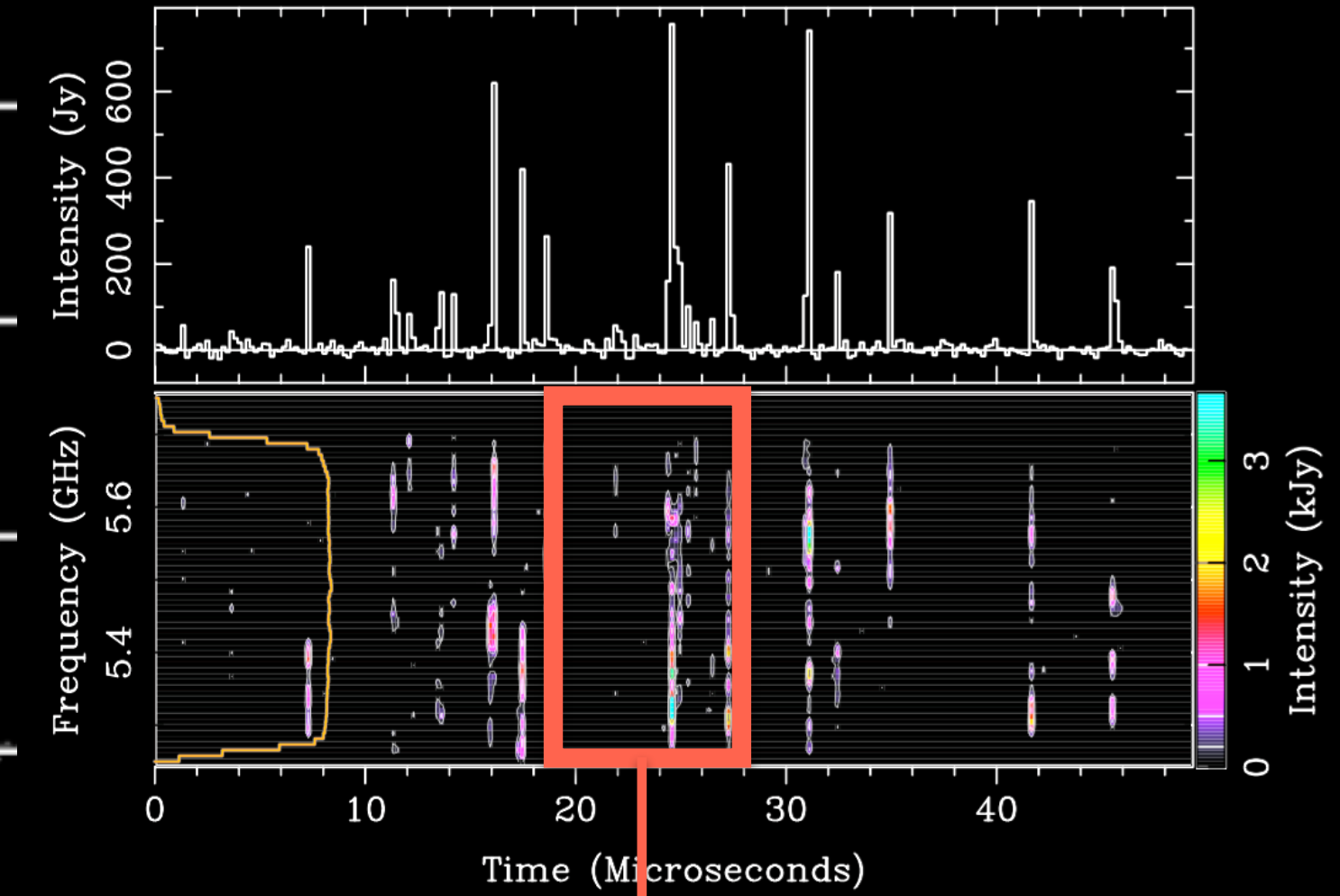
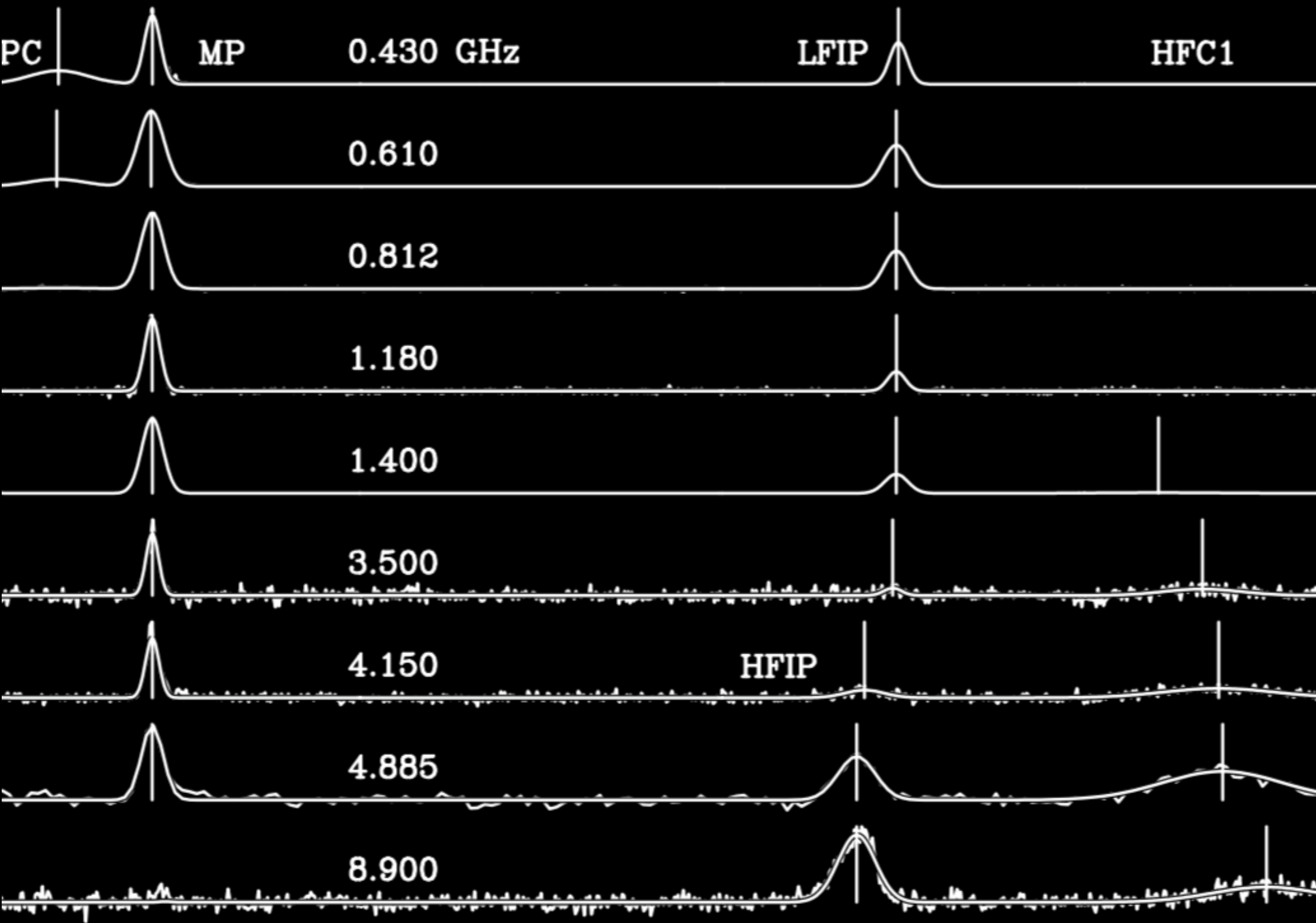


Bransgrove et. al., 2023, ApJL

Confirms order-of-magnitude luminosity
Core-cone geometry of the emission beam

PIC simulations with Osiris & Pigeon

CRAB RADIO EMISSION



MP and IP have high-energy emission counter-parts \Rightarrow definitely originate from the outer magnetosphere

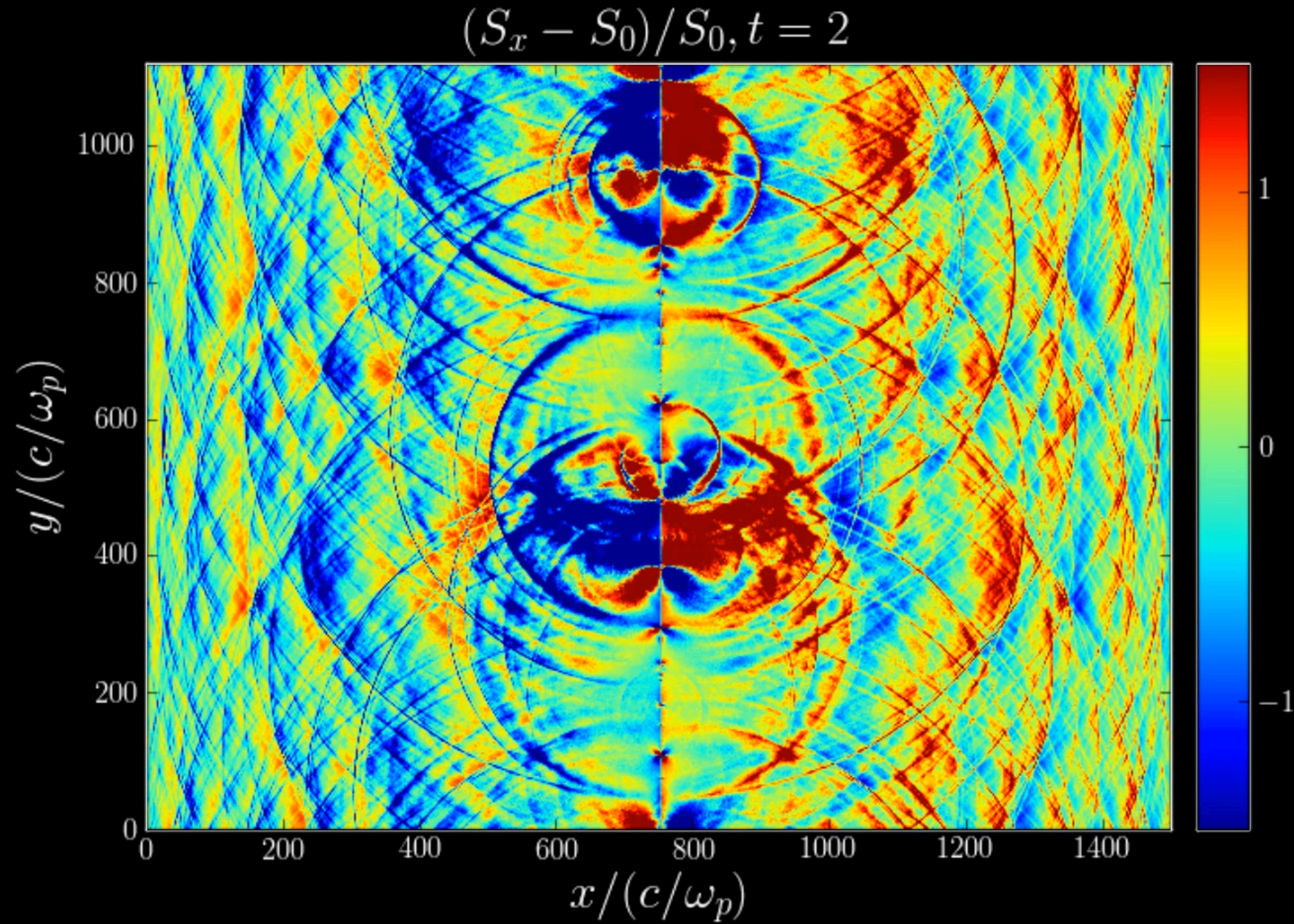
$$\nu_{\text{obs}} \delta t \sim \mathcal{O}(10)$$

Hankins, Eilek, 2016 (JPP)

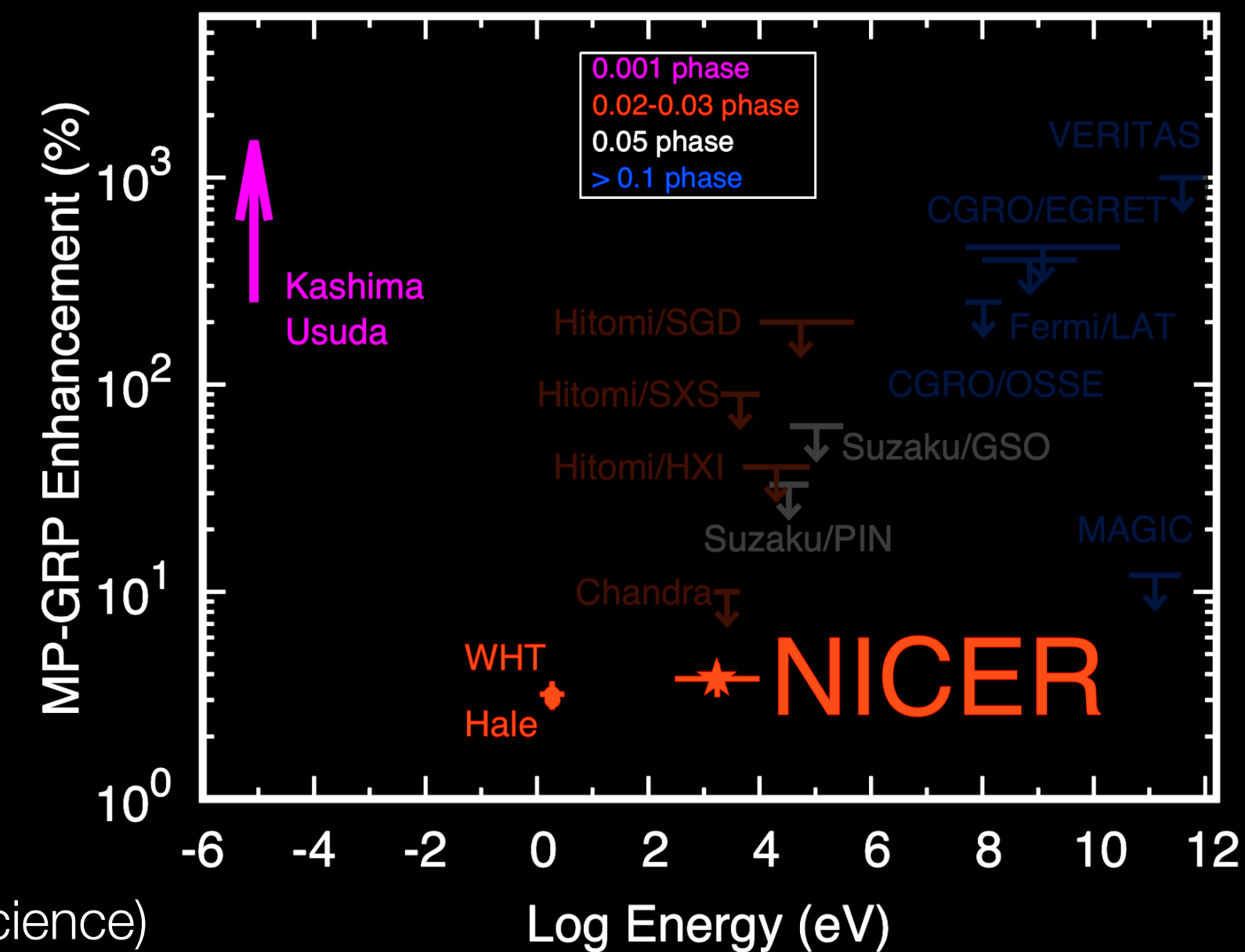
P ~ 33 ms

GIANT PULSES FROM RECONNECTION

- Current sheet breaks into plasmoids, plasmoids merge, EM waves are emitted.
- Amount of magnetic energy stored in a single plasmoid controls the brightness temperature. Can explain $T_B \sim 10^{38} \text{K}$!
- Plasmoid sizes set the frequency. Size is controlled by the strength of the radiative cooling, resulting in $\nu \sim c\Gamma/l \sim 1 \text{ GHz} \cdot B_6^{3/2}$. Requires MGs B-field strength at the light cylinder. Mergers of big plasmoids produce pulses with duration $\tau \approx 10/\nu$.
- Prediction: some correlation with the X-rays, also produced by reconnection.
- Similar waves exist in 3D, work in progress.



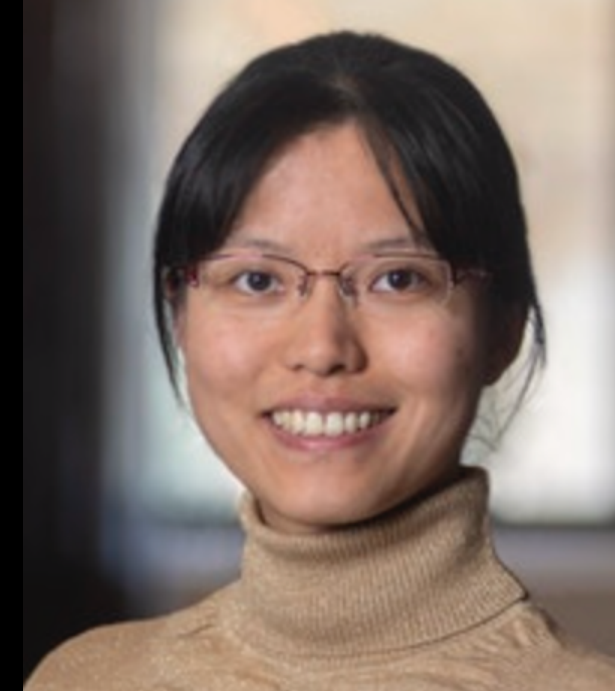
Philippov, Uzdensky, Spitkovsky, Cerutti, 2019 (ApJL)



Enoto et. al., 2021 (Science)

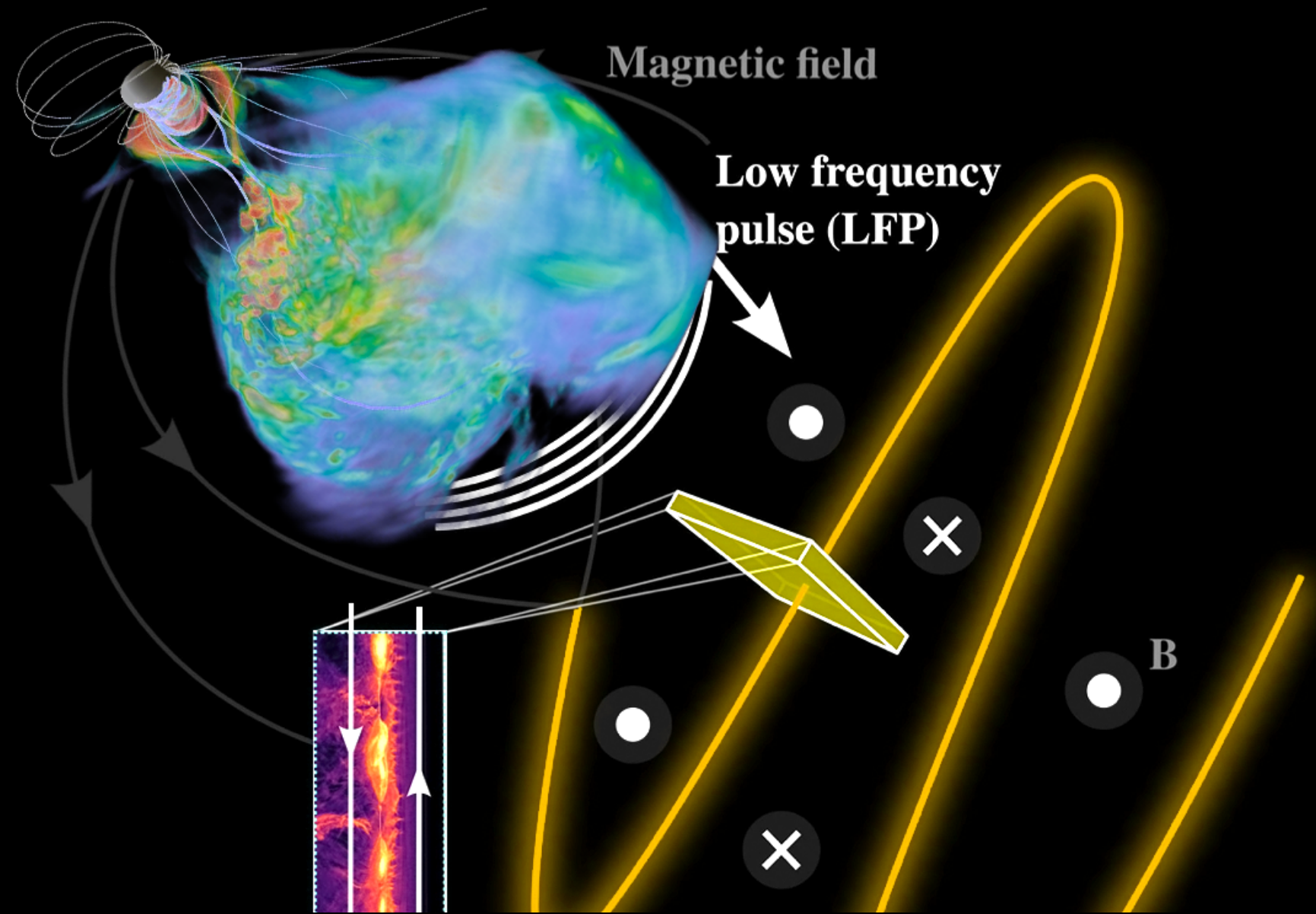
BEYOND PULSARS

APPLICATION TO FRBs: MAGNETIC EXPLOSIONS

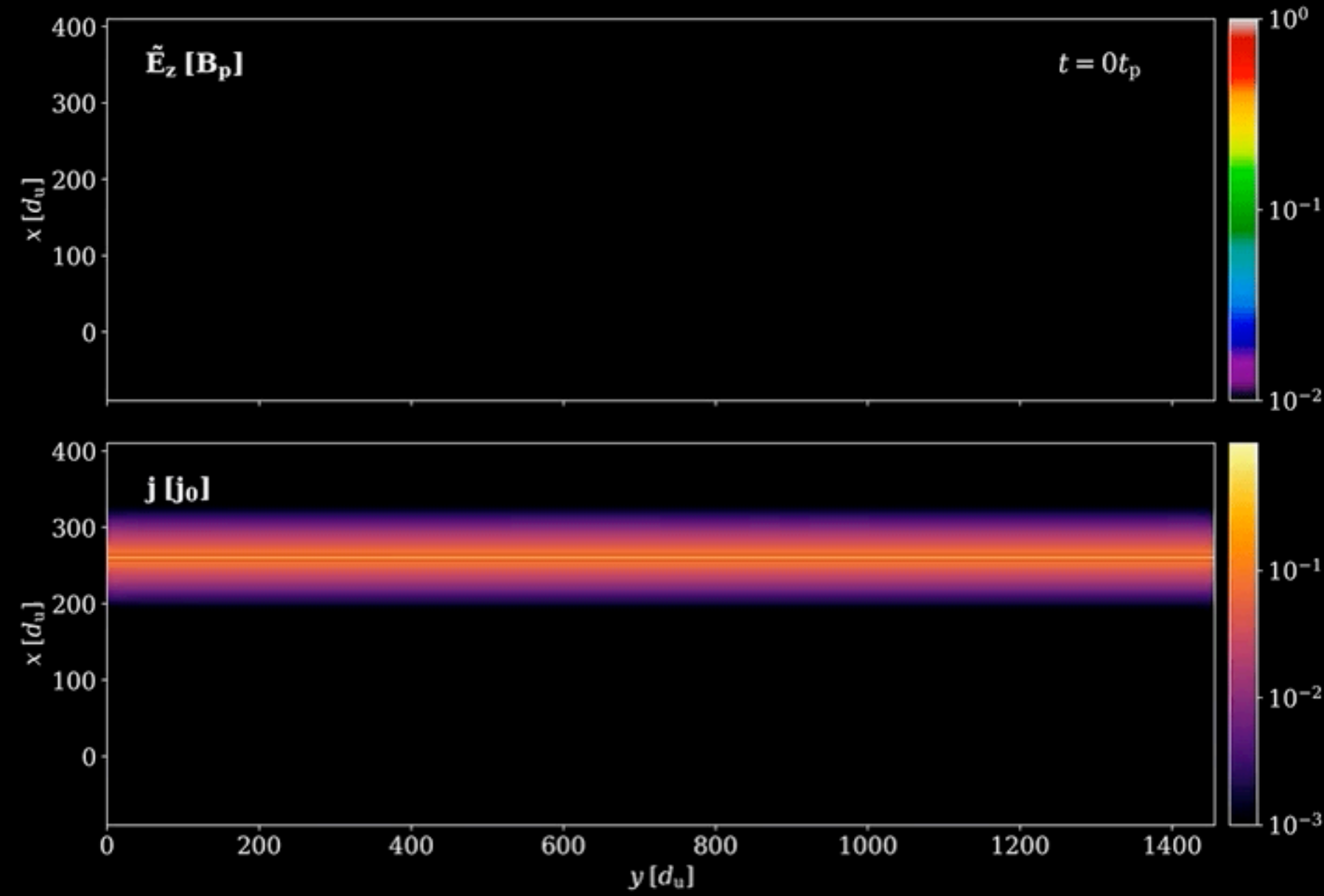
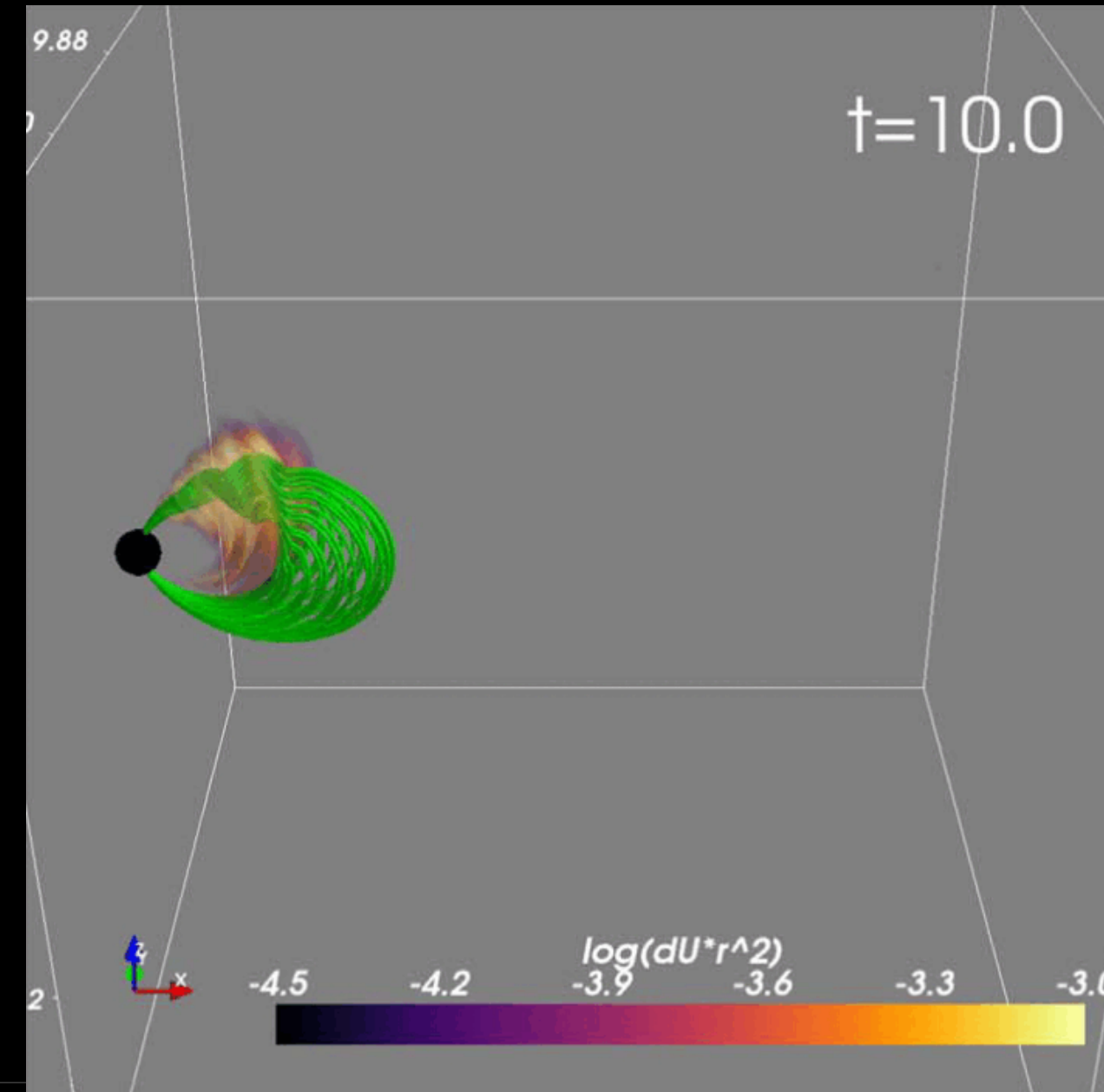


Potentially applicable to X-ray and FRB from galactic magnetar

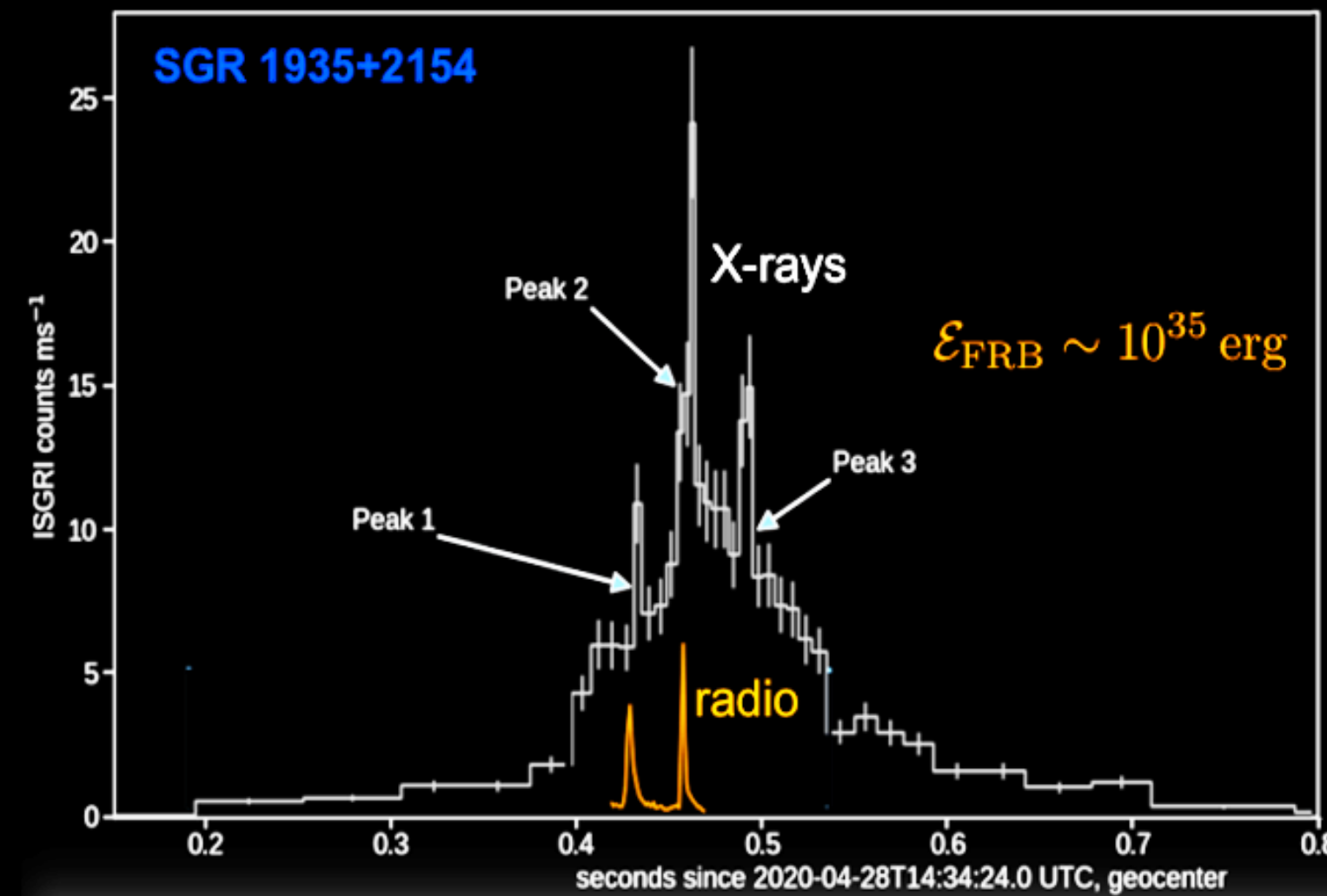
Large eruptions



Mild eruptions



April 2020: Fast Radio Burst from a local magnetar

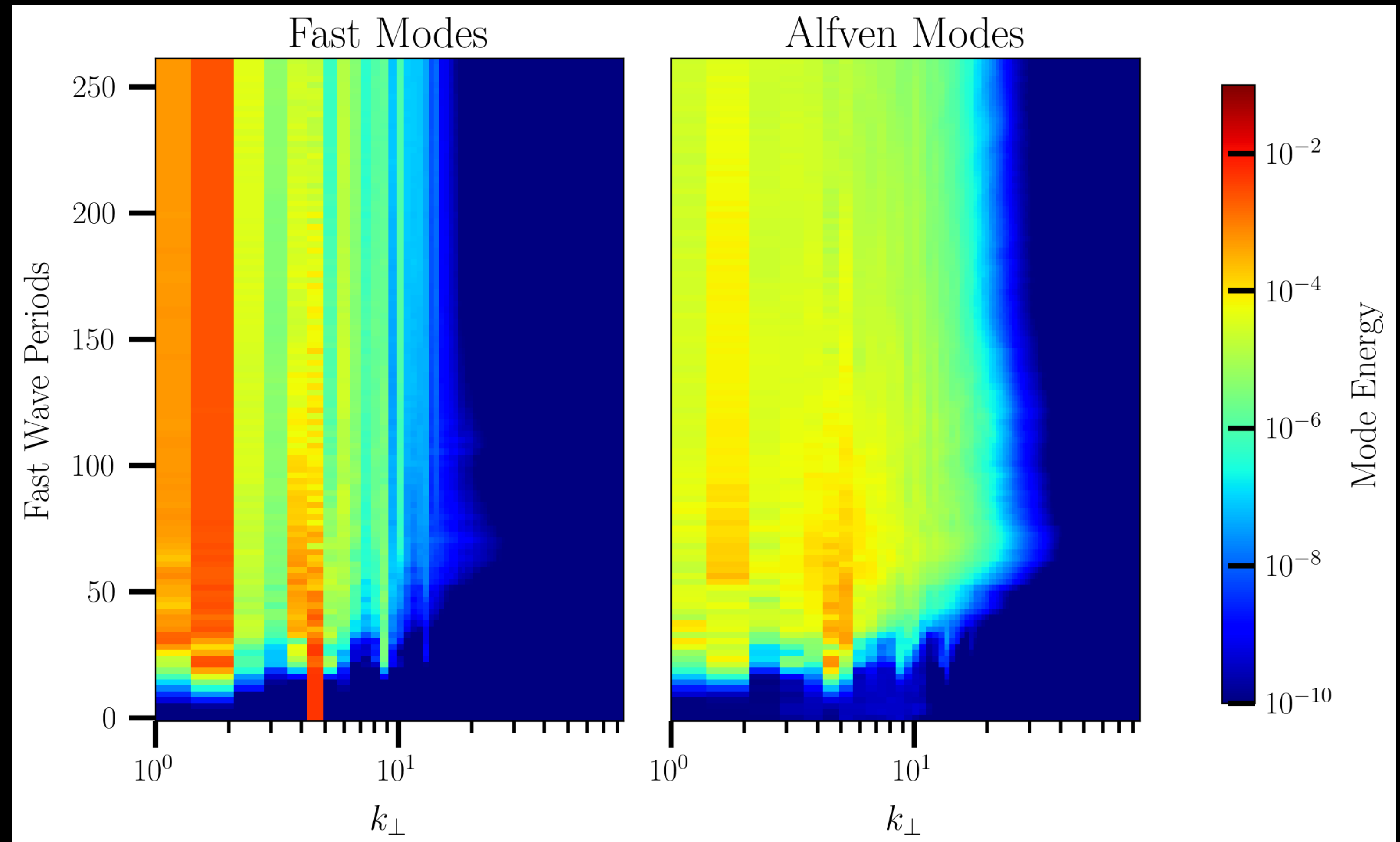


Lyubarsky, 2020 (ApJ)
 Mahlmann, Philippov et. al., 2022 (ApJL)
 Yuan et. al. (including Philippov), 2022 (ApJ)

FRB / GP PROPAGATION

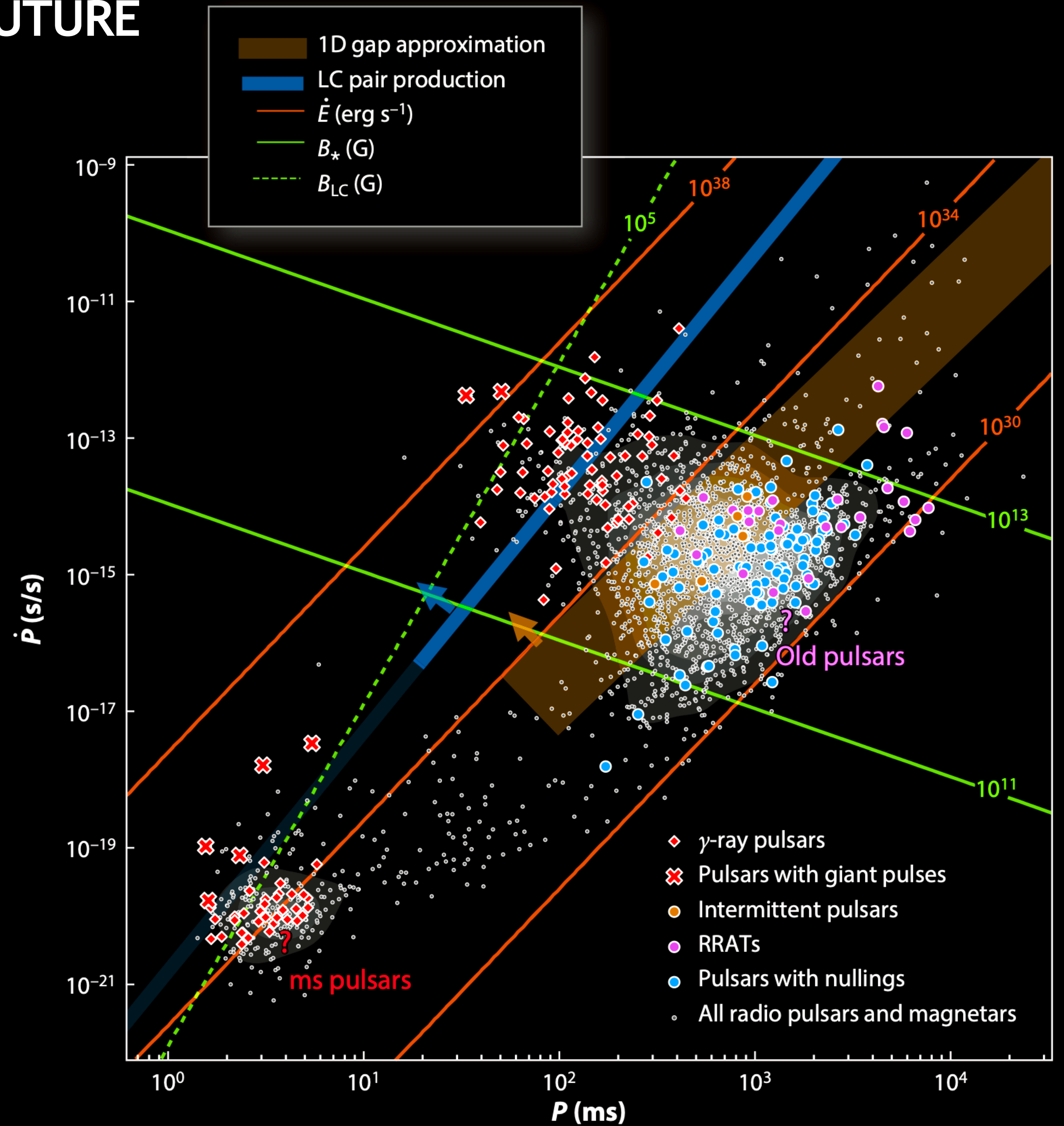


- High-amplitude FMS radiation is not free to leave.
- Quiescent magnetosphere: formation of shocks (Beloborodov 2023).
- More likely: surfing electromagnetic explosions.
- Non-linear wave interactions:
 - $F \leftrightarrow A+A$
 - $F \leftrightarrow A+F$



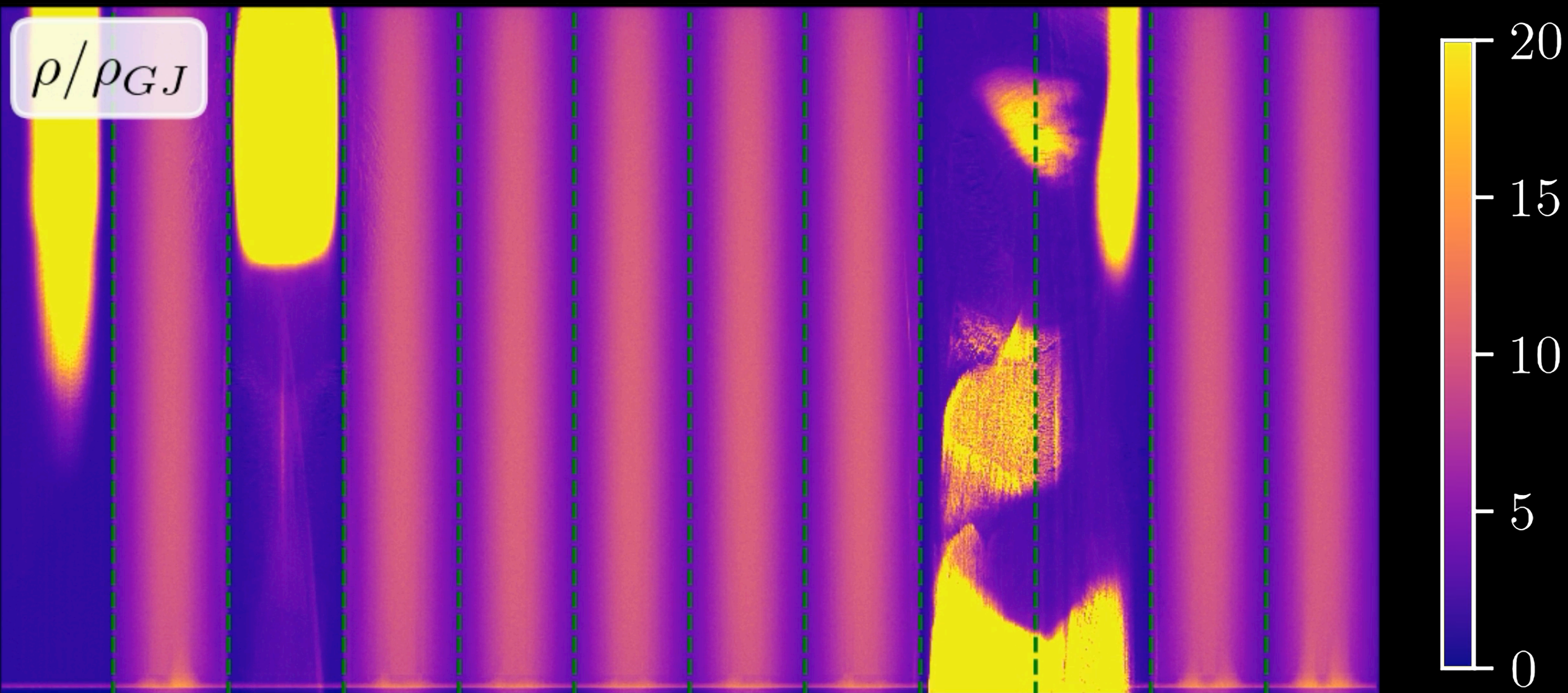
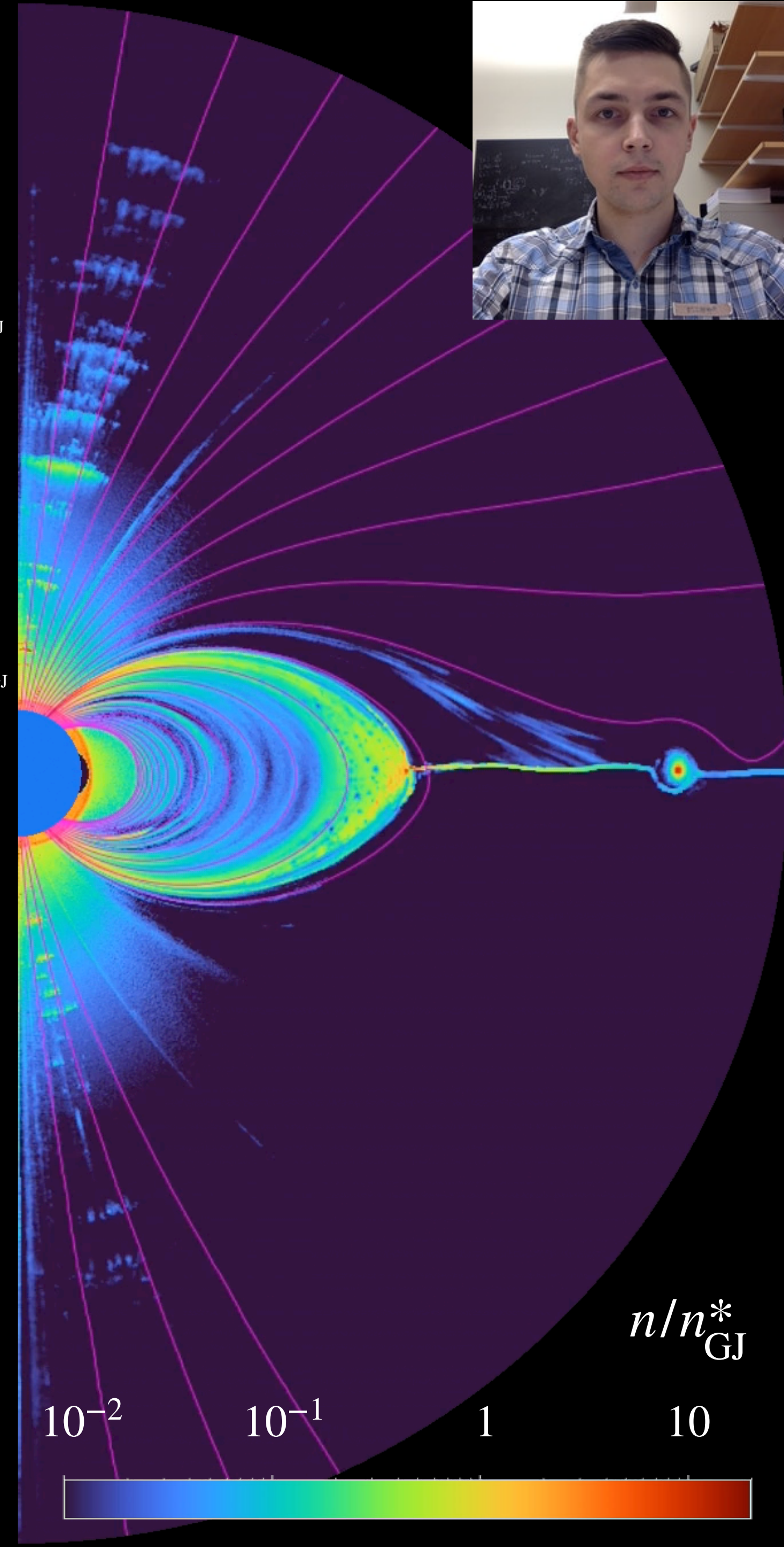
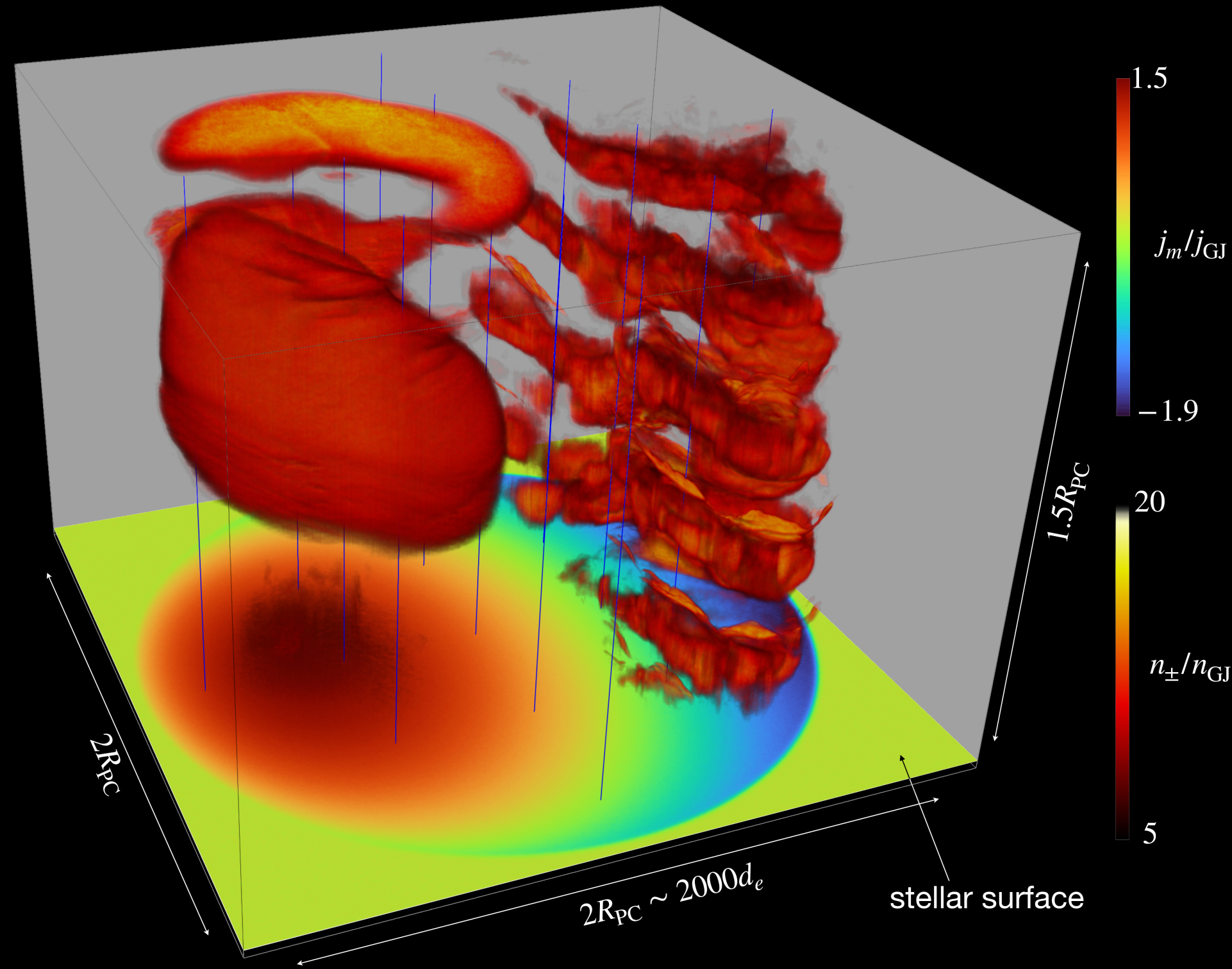
FUTURE

- Pair production in 3D
- Old and Millisecond Pulsars
- Variability: nulling, thunderstorms, raindrops, drifting subpulses
- Other bands: optical, X-ray, etc.

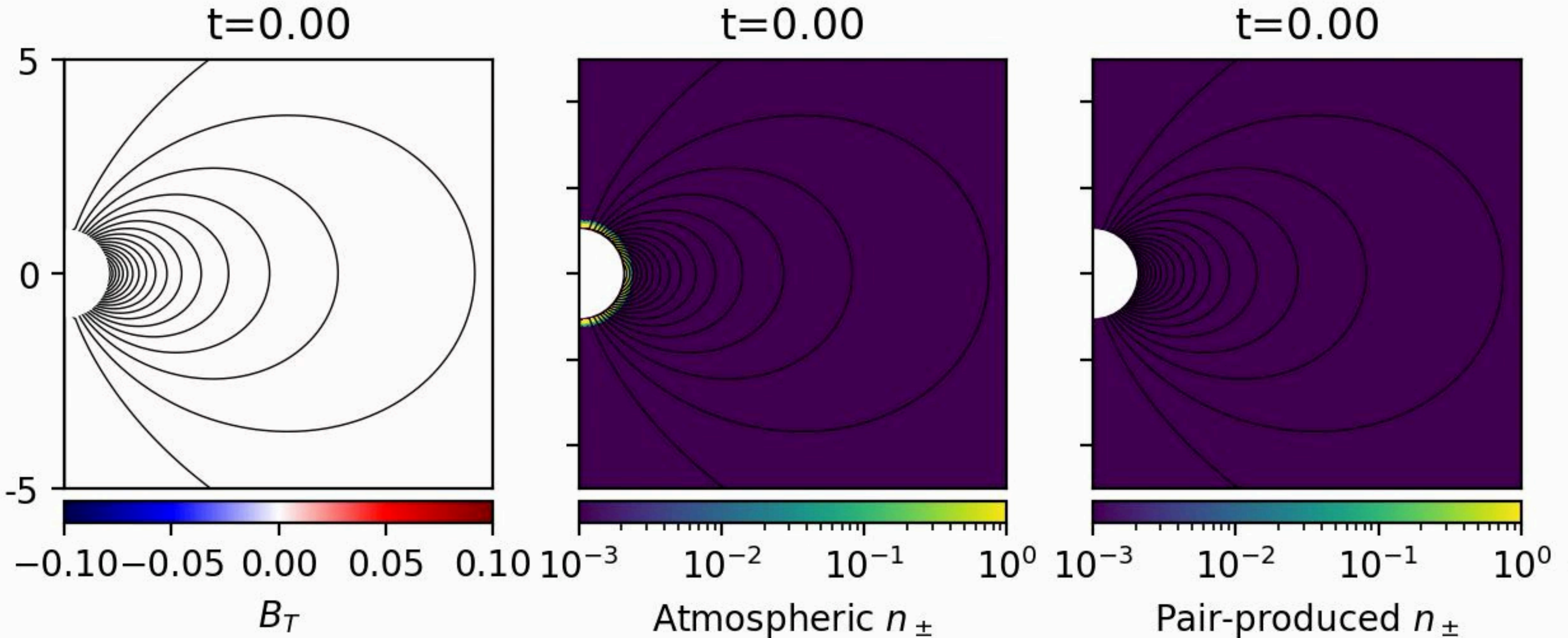


FUTURE

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MAGNETAR MAGNETOSPHERES



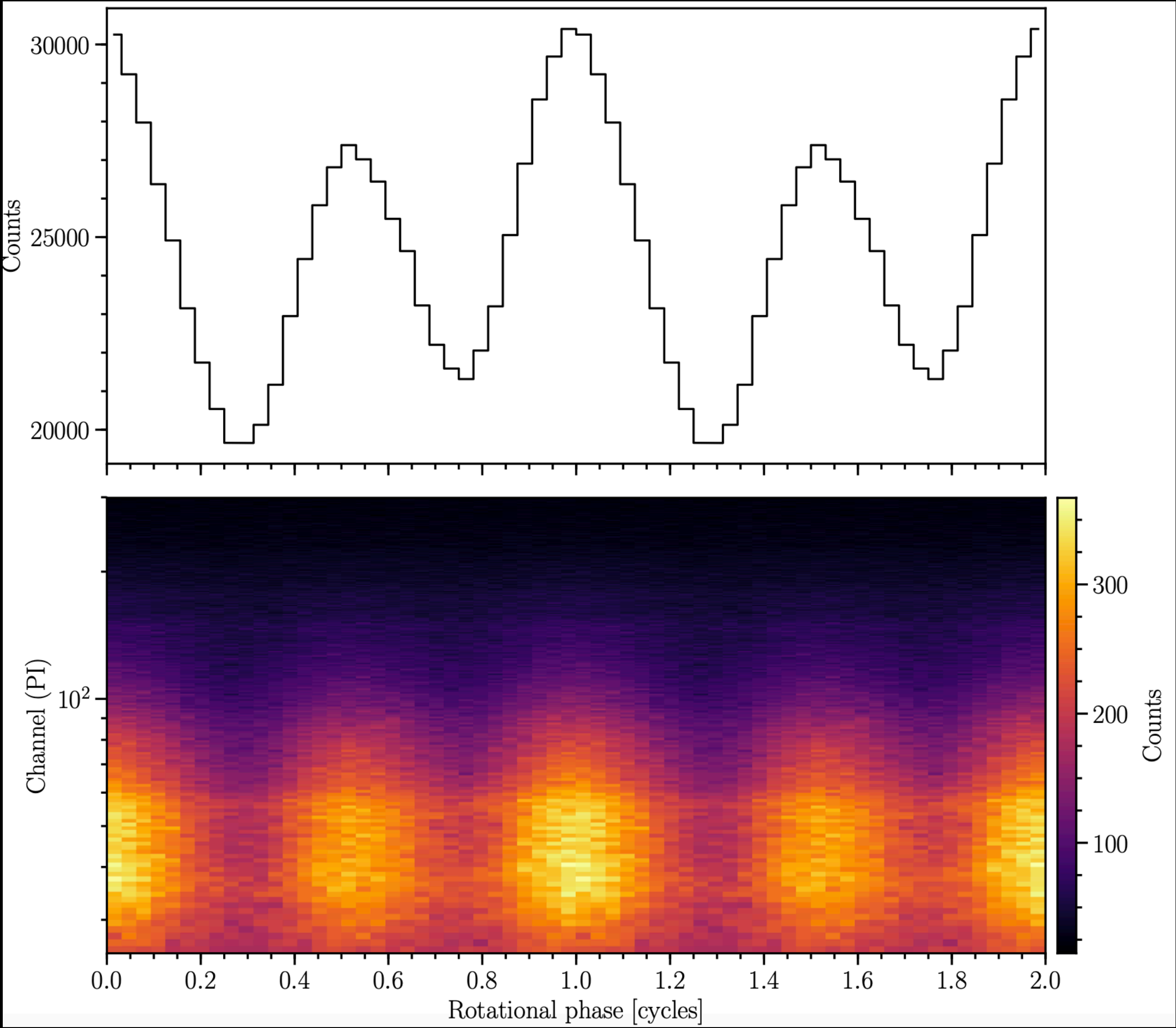
- Ongoing: RICS pair-production
- Near-term: inclusion of QED processes for strong fields

Conclusions and outlook

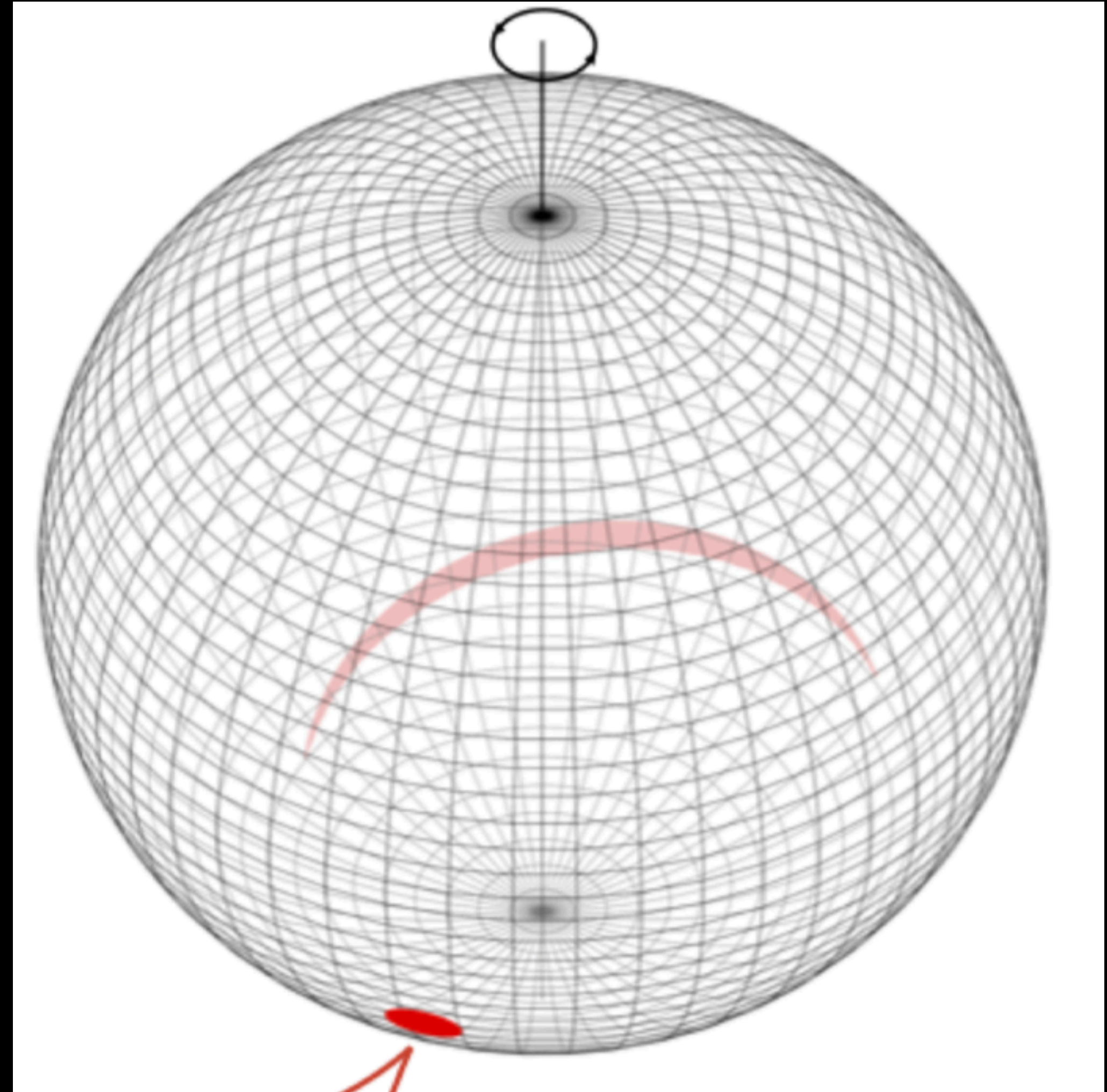
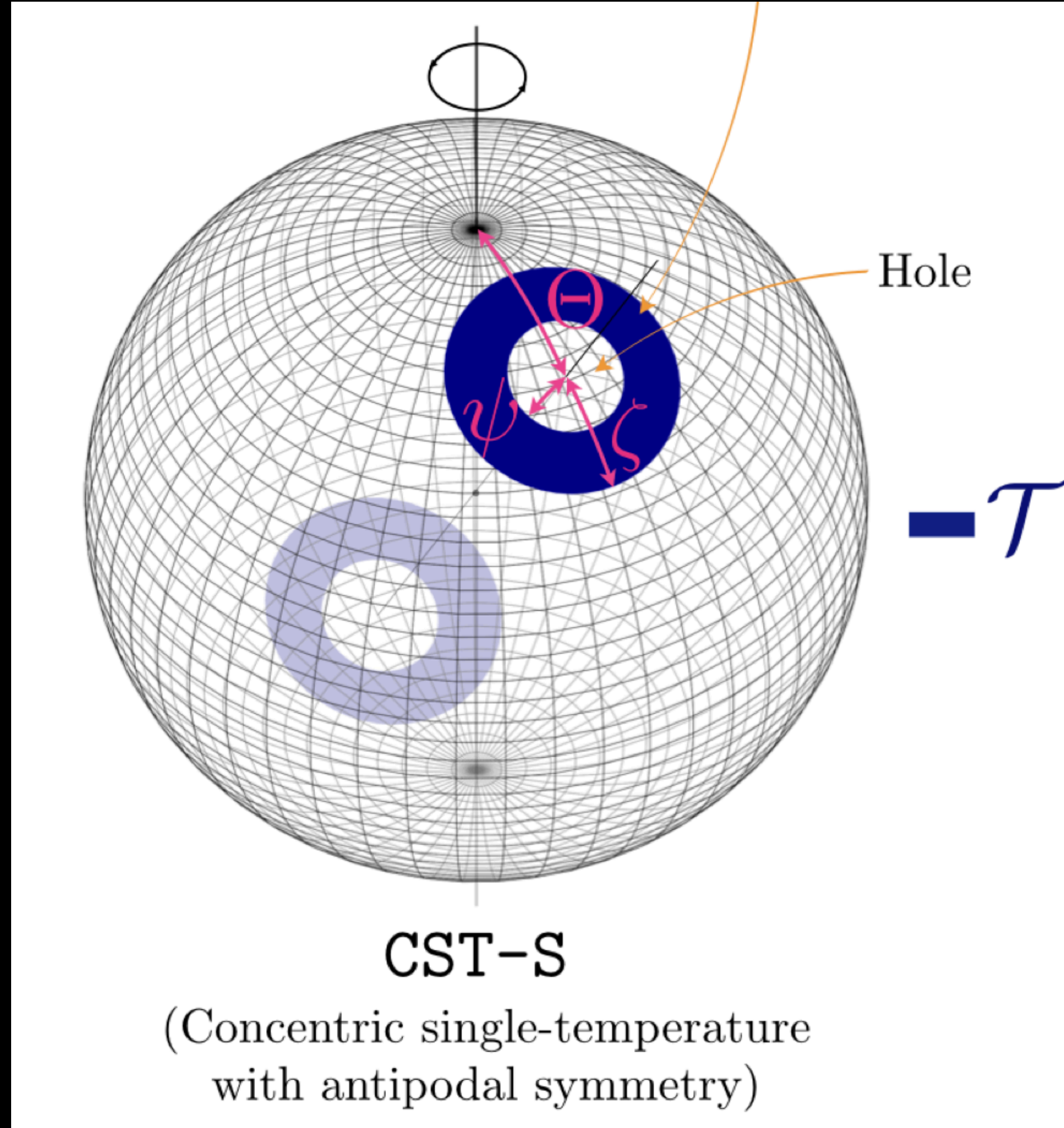
1. Origin of pulsar emission has been a puzzle since 1967 - kinetic plasma simulations are finally addressing this from first principles.
2. Current sheet is an effective particle accelerator. Particles in the sheet emit powerful gamma-ray mainly via synchrotron mechanism. Highest energy TeV photons can be produced in the current sheet as well.
3. Low altitude radio emission is produced during non-stationary discharge at the polar cap, not a plasma instability in the uniform plasma flow. Giant pulses and nanoshots are powered by plasmoid mergers in the current sheet beyond the light cylinder.

PULSARS AND NICER

J0030+0451



PULSARS AND NICER



of parameters to describe. The inferred mass M and equatorial radius R_{eq} are, respectively, $1.34_{-0.16}^{+0.15} M_{\odot}$ and $12.71_{-1.19}^{+1.14}$ km, while the compactness $GM/R_{\text{eq}}c^2 = 0.156_{-0.010}^{+0.008}$ is more tightly constrained; the credible interval bounds reported here are approximately the 16% and 84% quantiles in marginal posterior mass.